

# ESCAPING NEAR-RELATIVISTIC ELECTRON BEAMS FROM THE SOLAR CORONA

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## ABSTRACT

Processes in the solar corona are prodigious accelerators of energetic ions, and electrons. The angular distribution, composition, and spectra of energetic particles observed near Earth gives information on the acceleration mechanisms. A class of energetic particle observations particularly useful in understanding the solar acceleration is the near-relativistic impulsive beam-like electron events. During five years of operation the Advanced Composition Explorer (ACE) has measured well over 400 electron events. Approximately 25% of these electron events are impulsive beam-like events that are released onto interplanetary field lines predominantly from western solar longitudes. We extend our initial  $\sim 3$  year study during the rise to solar maximum (Haggerty and Roelof, 2002; Simnett et al., 2002) to a five year statistical analysis of these beam-like energetic electron events in relationship to optical flares, microwave emission, soft X-ray emission, metric and decametric type-III radio bursts, and coronal mass ejections.

## INTRODUCTION

Large solar flares and their coronal mass ejections (CMEs) are associated with the acceleration and release of large fluences of energetic ions and electrons that are observed in the inner heliosphere at 1 AU. The mechanism for the prompt acceleration and release of the solar energetic particles (SEP) has remained elusive. Simnett (1974) showed that the onsets of fully relativistic electron bursts observed near Earth were delayed relative to solar flare EM emissions by  $> 10$  minutes. Cliver et al. (1982) and Kahler (1994) showed that the onsets for energetic protons  $> 10$  MeV and the onsets for ground level events (GLE) were also related to CME driven shocks in the high corona with a typical release altitude of  $> 2$  solar radii ( $R_{\odot}$ ). Recent studies have shown that the near-relativistic beam-like electron events observed at the L1 Lagrangian orbit near 1 AU are released onto open field lines  $\sim 10$  minutes after metric and decametric type III radio bursts (Krucker et al., 1999; Haggerty et al., 2001; Haggerty and Roelof, 2002) as well as microwave, optical flare, and soft X-ray (SXR) emission (Haggerty and Roelof, 2002). These scatter-free electron events were found to be directly related to west limb CMEs with the intensity and spectral hardness of the electron events correlated with the speed of the CME (Simnett et al., 2002). For this study we have examined five years of ACE/EPAM (Electron, Proton, and Alpha monitor) near-relativistic electron observations, microwave, metric, and decametric radio emission,  $H\alpha$  optical flare emission, and SOHO/LASCO CME images. We report on timing delays between the injection of the SEP electron beams and various EM emissions. We also report on the correlations between the launch time and speed of the CMEs and the injection of the electron events.

## EVENT CRITERIA

The EPAM instrument (Gold et al., 1998) measures near-relativistic electrons ( $38 < E < 315$ ) keV in three different detector telescopes at  $30^\circ$ ,  $60^\circ$ , and  $150^\circ$  from the spacecraft spin axis. These detectors sample much of the  $4\pi$  solid angle at a resolution of up to 1.5 seconds/sector and supply the required angular distributions for the study. From the launch of ACE (Aug 25, 1997) through Mar 31, 2002, EPAM observed 406 near-relativistic

electron events where the increase in the 38-53 keV electron intensity  $> 10^3 \text{ (cm}^2\text{.sr.s.MeV)}^{-1}$ . A wide variety of intensity variations, angular distributions, and spectral evolution are observed in the events, so a precise method of selecting only well connected, scatter-free, beam-like electron events needs to be established. The method we have employed consists of three clearly defined observations:

1. Unambiguous onset time. The event onset needs to have a large enough signal to noise ratio that the onset time can be firmly established.
2. A scatter-free onset. The observation of velocity dispersion is required, so this also demands that the event is observed in at least two channels between 35 and 315 keV.
3. A field aligned angular distribution. The electron pitch angle distribution (PAD) must be beam-like along the IMF.

Based on these criteria, 113 near-relativistic, scatter-free, beam-like electron events (28% of 406) were selected for this study. Two different methods for determining the precise electron event injection time at the Sun have been reported. Krucker et al. (1999) identified injection times by plotting channel onset times versus inverse channel velocity. A straight-line fit to the points yields both the path length traveled and the release time at the sun. The method used for this study, Haggerty and Roelof (2002), uses the onset in the peak sector of extremely collimated electron PADs from the highest energy channel in which they were measured. The injection time is shifted back from the onset time by a transit time of 1.2 AU divided by the velocity in that highest channel. A comparison of the injection times from the two methods gives an approximately normal distribution with zero mean and a standard deviation of 2.5 minutes (S. Krucker 2002, private communication). We therefore conclude that the method employed by Krucker et al. (1999) and the method used in this study (Haggerty and Roelof, 2002) yield effectively the same results.

The CME coronagraph images were obtained from the Large Angle and Spectrometric Coronagraph Experiment (LASCO) C2 on the Solar and Heliospheric Observatory (SOHO). CME launch times were determined via the standard technique of mapping the height-time profiles back to  $1 R_{\odot}$ , while the speed to the CME is directly determined from the slope of the height-time profile. See Simnett et al. (2002) for a thorough discussion on the CME identification and analysis. Owing to several factors LASCO was not observing during 30 of the 113 events. In the remaining 83 electron events, 70 (84%) well correlated CMEs were observed.

The decametric radio burst observations are made using high resolution (16 second) WIND/WAVES (Bougeret et al., 1995) data. The timing of the type III radio bursts observed at 14 MHz ( $\sim 2R_{\odot}$  based on the density model of Leblanc et al., 1998) was obtained by taking a frequency cut of the WAVES channel associated with 14 MHz emission. The timing of the peak of the 14 MHz frequency cut was used as the timing of the decametric type III emission. A one-to-one correlation was obtained for all 111 of the electron events for which WIND/WAVES data were available. The timing of the decametric type III emission (observed near 1 AU) and all of the EM emissions used in this study are taken back to the Sun by 500 seconds so we may refer to them as an emission time, and compare them directly with the electron injection times.

The time of the remaining EM emissions used in this study, microwave and metric radio bursts, H $\alpha$  optical flares and SXR emission were obtained from the NOAA Solar Geophysical data (SGD). For some of the electron events, multiple metric type IIIs, H $\alpha$  flares, or SXR events fell within a  $\pm 1$  hour interval of possible association. To ensure we were not biasing the selection of EM events, we included all events within the interval. There were 20 electron events with no reported metric type III emission ( $\pm 1$  hour of electron injection) in the SGD. For the remaining 93 events, a total of 174 metric type III radio bursts fell within the timing intervals and have been included in the statistical analysis presented in this study. There were 26 electron events with no SXR emission reported while 7 electron events had multiple possible SXR associations. The total number of 87 (77%) of the beam like electron events have associated SXR emission. There were 39 electron events with no reported H $\alpha$  optical flare reported in the SGD. Of the remaining 74 events, 26 had at least two possible H $\alpha$  observations within the prescribed interval. Therefore the total number of H $\alpha$  observations presented in this study is 116 that are correlated with 79 near-relativistic electron events. There were 75 electron events with no reported microwave bursts ( $> 8800$  MHz) in the SGD. Of the remaining 38 (33%) electron events, 4 had multiple microwave bursts reported yielding a total of 43 microwave bursts used in this study.

## WELL CONNECTED ELECTRON EVENTS

While the event selection criteria result in a list of electron events that propagate scatter-free, that are field-aligned, and that are clean enough to use in a statistical timing study, they do not demand that the selected events be well connected to western longitudes. While we do not argue that these electron events are accelerated precisely at the flare site, we will use the longitude of the H $\alpha$  flares to estimate the approximate longitude of the electron injection. The left panel of Figure 1 shows that the longitude of the electron-event-correlated H $\alpha$  flares had a median longitude of 44°. Note that the majority of eastern flares shown in Figure 1 have at least one western flare that was associated with the same electron event, so the former are therefore probably not associated with the electrons that arrive at 1 AU. On the ordinate of the left panel of Figure 1 is the time delay between the start time of the optical flare and the injection of the near-relativistic electrons. We find no clear trend in these observations that would indicate that the electron injection delay is related to the longitude of the associated event.

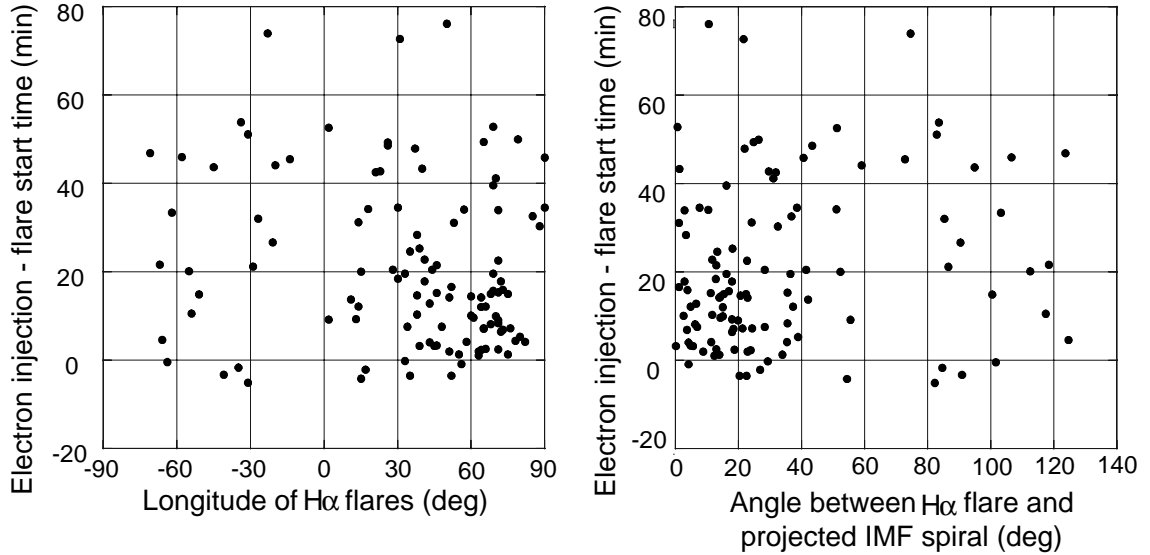


Fig. 1. Left panel. Longitude of H $\alpha$  flares versus the time delay between start time of the flare and the electron injection. Right panel. Angle between the longitude of the H $\alpha$  flare and the projected longitude of IMF connection based on the solar wind speed at the time of the electron onset versus the time delay between start time of the flare and the electron injection.

Assuming that the magnetic-field lines are frozen into the expanding plasma at a source surface in the corona, one may estimate the actual longitude of IMF connection based on the local plasma speed at the time of the electron onsets. The right panel in Figure 1 shows the difference between the angle of the H $\alpha$  longitude and the projected IMF spiral angle versus the time delay between the optical flare start time and the electron event injection. Although this gives another estimate of how well connected the optical flare site is to the L1 orbit, it does not seem to order the delays. Nor is there any clear pattern in this combination of observations that would indicate acceleration and release of electrons due to a source propagating along the solar surface.

## LOW CORONAL ELECTROMAGNETIC EMISSION

The traditional view of HXR emission from the low corona consists of a beam of non-thermal electrons being driven down the coronal field lines into the chromosphere where thick target bremsstrahlung results in the emission of the HXR. Timing between the escaping population of energetic electrons and these low coronal emissions can indicate whether or not the escaping population is accelerated through the same process. Microwave emission and HXR emissions are well correlated, and the emission time between them is known to differ by less than a few seconds. We therefore take the microwave emission as a proxy of the HXR event. Figure 2 shows the distribution of timing delays between the injection of the escaping near-relativistic electron beams and the microwave emission. While a few of the electron events in our list are associated with the start time of the microwave events ( $t=0$  indicates the start time of the microwave radio emission), the median delay of the injection of the escaping population is 14 minutes (average 13 minutes).

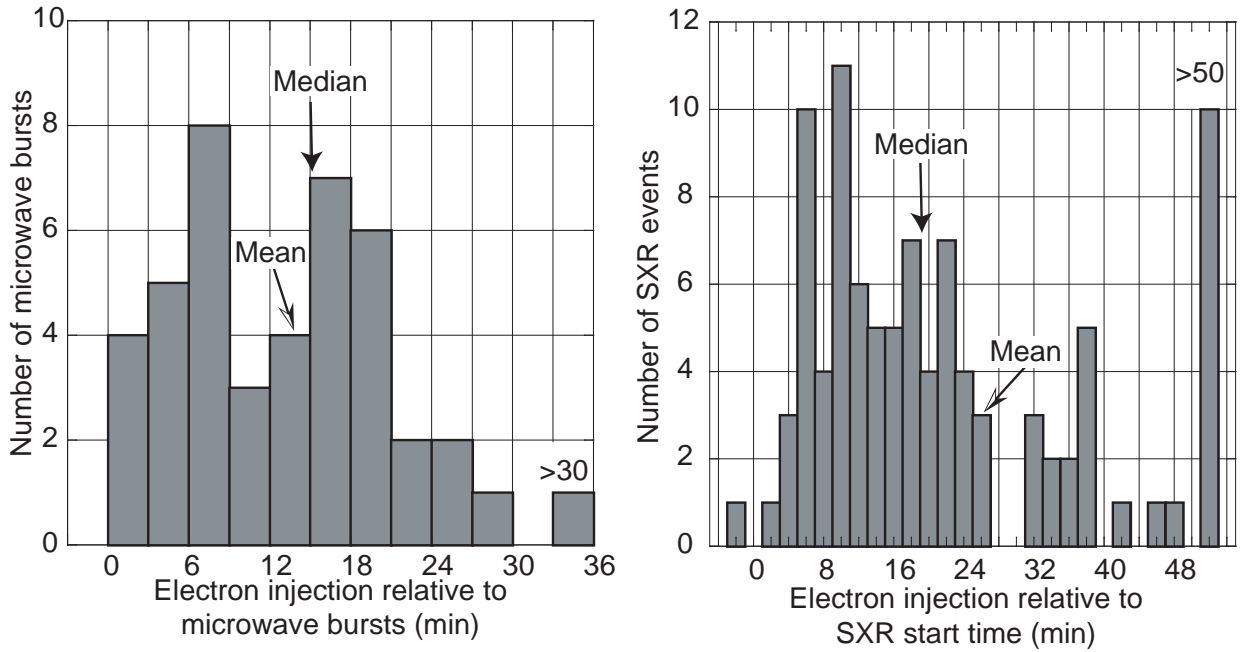


Fig. 2. Left panel. Distribution of delay times at the Sun between the injection of the escaping near-relativistic electron beams and the emission of the microwave radio bursts. Right panel. Distribution of delay times at the Sun between the escaping population and the SXR emission.

The right panel of Figure 2 shows the distribution of delays between the injection of the escaping population of near-relativistic electrons and the start time of the SXR event. The last histogram in the right panel of Figure 2 at  $t=52$  minutes give the integral of the distribution that extends beyond  $t>50$  minutes. The median delay between the escaping population and the SXR start time is 17 minutes. Comparisons to both the HXR (microwave) start time and the SXR start time show that the injection of the escaping near-relativistic electrons is delayed by  $\sim 10$ -20 minutes relative to the start of the solar flare process.

## TWO DISTINCT POPULATIONS OF ESCAPING ELECTRONS

Type III radio bursts are known to be produced by plasma-wave interactions from electron beams escaping into the high corona and interplanetary space. The energy of the electrons within the escaping beam are known to be non-thermal and contain a spectrum of electron energies. The speed at which the type III exciter front propagates corresponds to an electron energy of a few keV (Haggerty et al., 2001; Haggerty and Roelof, 2002). We have shown that the injection of the escaping population of near-relativistic electrons is delayed with respect to optical flares observed in H $\alpha$  as well as low coronal HXR and SXR emission. We now show that a similar delay exists between the injection of the escaping energetic electrons and the type III burst produced by lower energy electrons. Figure 3 shows the distribution of delays with respect to the metric type III (left) and the decametric type III radio bursts (right). In both figures  $t=0$  is the start time of the type III emission and the histograms detail the distribution of the energetic electron injection delays. The last histogram at  $t=50$  minutes in the left panel of Fig. 3 represents the integral of the distribution that extends beyond  $t>50$  minutes. The median delay between the start time of the metric type III radio burst and the injection of the near-relativistic electrons is  $\sim 13$  minutes while the median delay between the start time of the decametric type III radio burst and the injection of the high energy electrons is  $\sim 8$  minutes. These histograms give direct evidence that the injection of the near-relativistic electrons that are observed at 1 AU are significantly delayed with respect to the type III radio bursts. No evidence is observed in either of the histograms that near-relativistic ( $E>53$  keV) electrons are part of the electron population that produces the type III emission, at least for those electron events whose onsets have significant SNR at EPAM (see Haggerty and Roelof, 2002 for an analysis of the SNR and quality of the observed onsets).

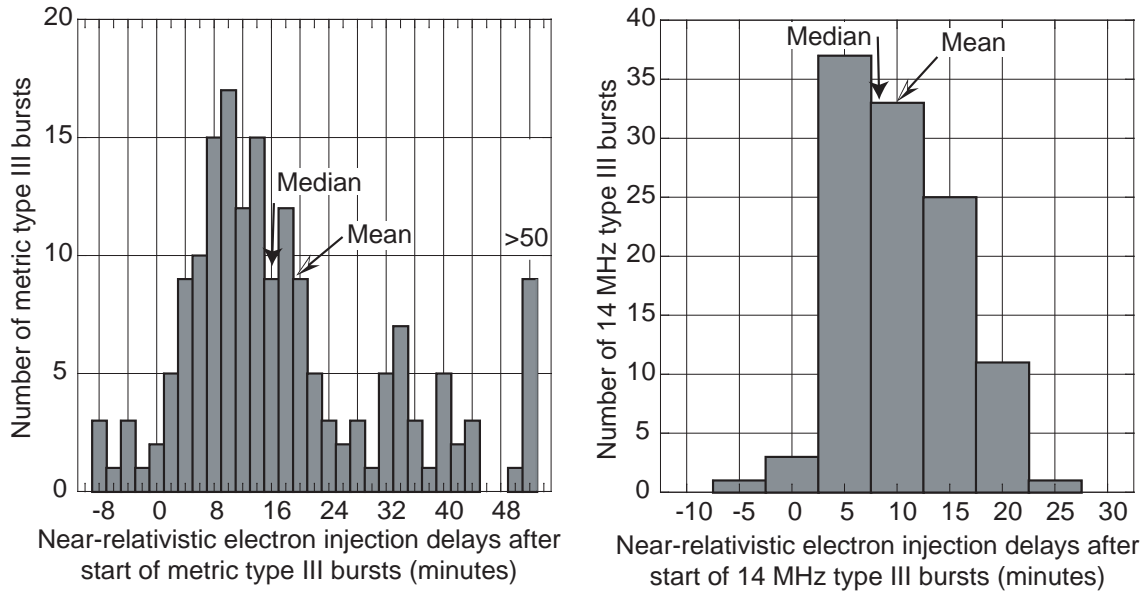


Fig. 3. The delay between the injection of the near-relativistic electron beams and the start time of the metric type III radio bursts (left) and the peak time of the decametric type III radio bursts (right). The time ( $t=0$ ) in each panel indicates the start time/peak time of the metric/decametric type III radio burst.

### RELATIONSHIP BETWEEN ELECTRON INJECTIONS AND CMEs

Simnett et al. (2002) showed that the injection of the near-relativistic electron beams is delayed with respect to the launch of the CME, that the height of the CME at the time of electron injection is a few  $R_o$ , and that the speed of the CME is related to both the intensity and the spectral hardness of the electron event. The event list used by Simnett et al. (2002) contained 79 of the electron events included in this study, and our extended list confirms those results. The left panel of Fig 4 shows the height of the CME at the time of electron injection (the product of the CME velocity times the delay of the electron injection after the CME launch) and the right-most bin represents the total number of CME heights  $> 5 R_o$ . The median height of the CME when the energetic electrons are injected is  $2.4 R_o$ , in agreement with Simnett et al. (2002). The right panel shows that the delay between the injection of the electrons and the launch of the CME is roughly inversely proportional to the velocity of the CME. Some scatter is expected, because a perfect correlation ( $R=1.0$ ) would correspond to a single injection height (which is unlikely).

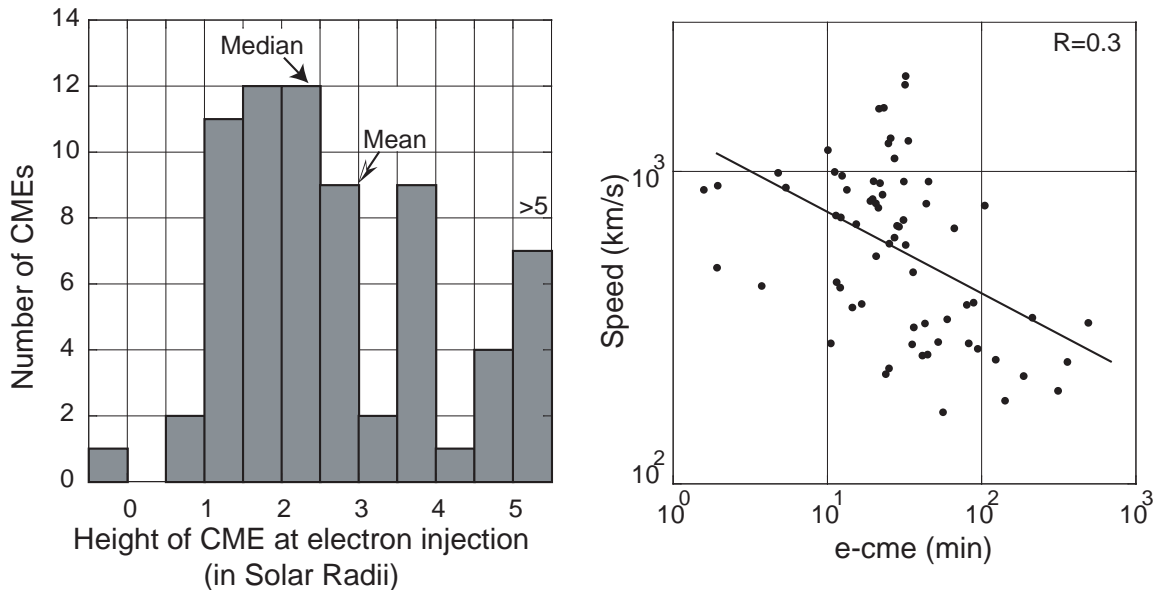


Fig. 4. Left panel. Height of the CME at the time of near-relativistic electron injection. Right panel. Delay between the injection of the energetic electrons and the launch of the CME on the abscissa with the speed of the CME on the ordinate.

## DISCUSSION AND CONCLUSIONS

It has been known for years that the high energy ions (Kahler, 1994) and relativistic electrons (Simnett, 1974; Cliver et al., 1982) are accelerated and released into the interplanetary medium, delayed with respect to related EM emissions in large solar flares. It was suggested that the ions, and possibly electrons, are accelerated by a CME driven shock and that the delay is due to the shock propagating to open field lines. Only recently has it been clearly shown that the near-relativistic electrons show a similar delay (Krucker et al., 1999; Haggerty and Roelof, 2002) and that the delay is inversely correlated with the speed of the CME, that the height of the CME at the time of the injection of the electrons is a few  $R_{\odot}$ , and that the intensity and spectral hardness of the near-relativistic electron events is directly related to the speed of the CME (Simnett et al., 2002). In this updated survey of five years of scatter-free beam-like near-relativistic electron events we have confirmed here that there is no evidence near 1 AU for escaping near-relativistic electrons as a constituent of the electron population that produces the type III radio bursts. It has been suggested that the propagation of a disturbance around the sun from the flaring region to the well connected region might be responsible for the delays observed in the near-relativistic electron injection (Krucker et al., 1999). However, we have shown that the angle between the flaring region based on the observed  $H\alpha$  longitude and the foot point of an Archimedean spiral connecting to 1 AU (based on the local solar wind speed at 1 AU) do not order the delays observed in the injection of the near-relativistic electron events. We have shown that there is a good relationship between beam-like near-relativistic electron events and SOHO/LASCO observed coronal transients (mainly traditional CMEs), that the height of the CME at the time of the electron injection is a few  $R_{\odot}$ , and that the delay in the injection of the electrons after the CME launch is inversely proportional to the speed of the CME. We therefore conclude that the delays are not due to a disturbance propagating along the disk of the Sun from the flaring region to a well connected Archimedean foot point, but are most likely due to the upward propagation of the coronal transient and subsequent shock acceleration and injection onto open field lines.

## ACKNOWLEDGMENTS

We thank the EPAM, LASCO, and WIND/WAVES consortia. The work performed at the Johns Hopkins University Applied Physics Laboratory has been supported by NASA under task 009 of contract NAS5-97271 and LWS grant NAG512479. Computing facilities at the University of Birmingham were supported by Starlink.

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Manuscript received 19 October 2002; revised 21 January 2003; accepted 22 January 2003