

# FAST satellite wave observations in the AKR source region

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**Abstract.** The Fast Auroral SnapshoT (FAST) satellite has made observations in the Auroral Kilometric Radiation (AKR) source region with unprecedented frequency and time resolution. We confirm the AKR source is in a density depleted cavity and present examples in which cold electrons appeared to have been nearly evacuated ( $n_{hot} > n_{cold}$ ). Electron distributions were depleted at low-energies and up-going ion beams were always present. Source region amplitudes were far greater than previously reported, reaching  $2 \times 10^{-4}$  (V/m)<sup>2</sup>/Hz (300 mV/m) in short bursts with bandwidths generally  $< 1$  kHz. Intense emissions were often at the edge of the density cavity. Emissions were near or below the cold plasma electron cyclotron frequency in the source region, and were almost entirely electromagnetic. The  $|E|/|B|$  ratio was constant as a function of frequency and rarely displayed any features that would identify a cold plasma cutoff or resonance.

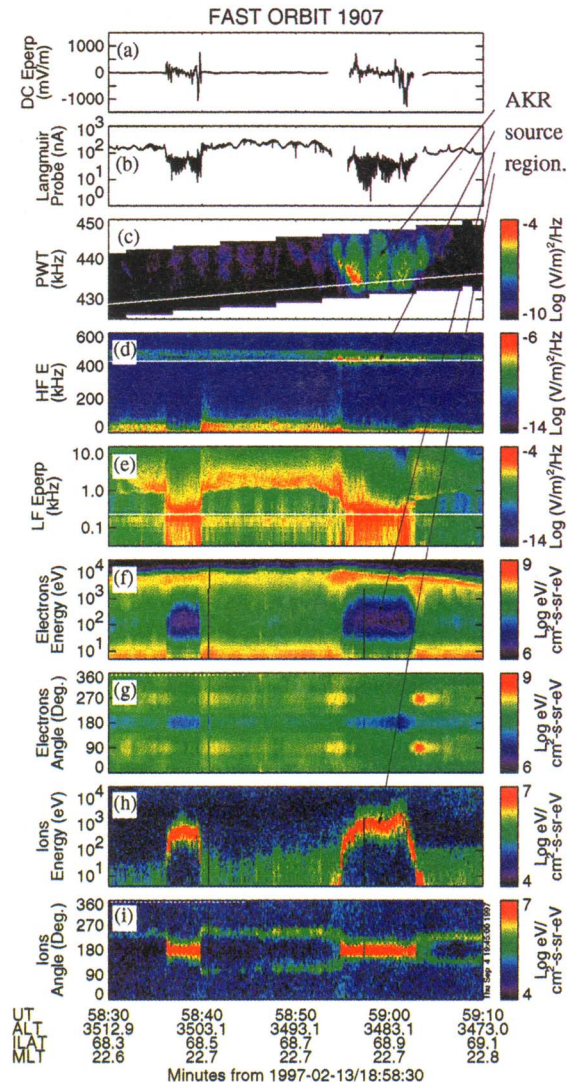
## Introduction

One of the outstanding features of the auroral zone are powerful radio emissions known as Auroral Kilometric Radiation [Gurnett, 1974]. Similar emissions have been observed from all of the magnetized planets [Zarka, 1992]. AKR propagates primarily in the R-X mode and is known to originate on auroral field lines in regions of depleted density where down-going, electrons are accelerated [Benson and Calvert, 1979, and references therein].

AKR is believed to be generated near the electron cyclotron frequency ( $f_{ce}$ ) from the free energy of the non-Maxwellian auroral electron distribution. A loss cone instability amplified by a subtle relativistic effect was one of the first widely accepted growth mechanisms [Wu and Lee, 1979], later expanded into the "cyclotron maser instability". Viking observations suggest that AKR is generated by electron distributions with  $d f/dv_{\perp} > 0$  [Louarn, et al., 1990]. These electron distributions had an absence of low-energy electrons and displayed a trapped electron population. Relativistic modifications were shown to be important in regions where the "hot" plasma density of the precipitating electrons dominates the plasma dispersion [Pritchett, 1984, Pritchett and Strangeway, 1985].

Viking satellite observations of the AKR source region [Bahnen et al., 1987; de Feraudy et al., 1987; Roux et al., 1993] showed intense emissions with a spectral peak within  $\sim 2.5\%$  of the cold plasma electron cyclotron frequency ( $f_{ce}$ ), occasionally extending below  $f_{ce}$ . The largest amplitudes were reported to be  $\sim 10^{-6}$  (V/m)<sup>2</sup>/Hz ( $\sim 50$  mV/m).

In this article, we present electric and magnetic field observations in the AKR source region from the FAST satellite. The electron distributions are described in a companion paper [Delory et al., 1998] and the instruments are described elsewhere [Ergun et al., 1998; Carlson et al., 1997]. FAST observations include the first digital waveform capture in the AKR source region. These data



**Figure 1.** (a) The D.C. electric field. (b) Langmuir probe current. A 2 nA current roughly corresponds to  $\sim 1$  cm<sup>-3</sup> cold density or  $\sim 0.1$  cm<sup>-3</sup> hot density. (c) Wave power as measured by the Plasma Wave Tracker. The white line is  $f_{ce}$ . These data have spin modulations. The PWT covers a  $\sim 16$  kHz band, the center frequency updated every 1 s. (d) High-frequency electric field power. The white line is  $f_{ce}$ . (e) Low-frequency, perpendicular wave power. The white line is  $f_{CH+}$ . (f-g) Electron energy flux versus energy and time and versus pitch angle and time. (h-i) Ion energy flux.

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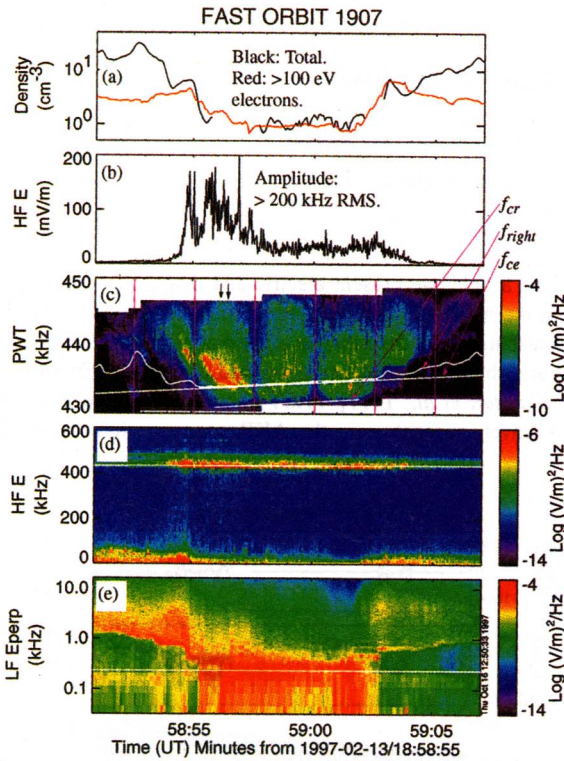
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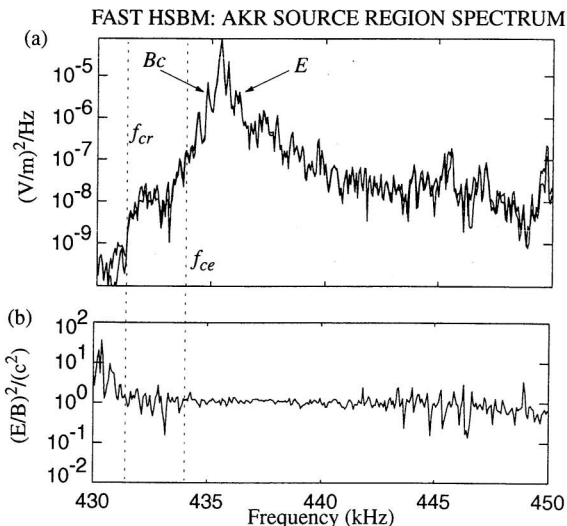
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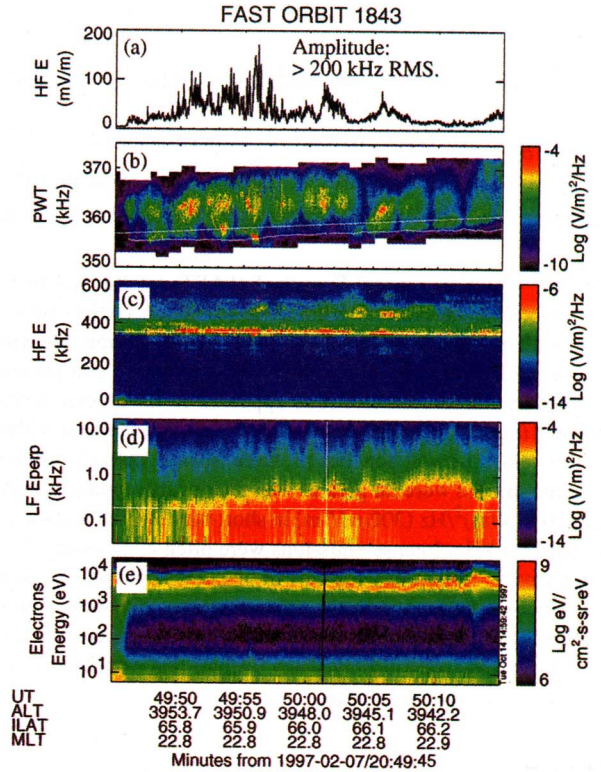


**Figure 2.** The AKR source region. (a) The total electron density (black) derived from fits to the whistler/Langmuir dispersion. The hot (>100 eV) electron density is in red. (b) The >200 kHz, RMS amplitude of  $E$ . (c) Wave power as measured by the Plasma Wave Tracker. The white lines are  $f_{\text{right}}$ ,  $f_{ce}$ , and  $f_{cr}$  ( $eB/m\gamma$ ), with an uncertainty of  $\pm 400$  Hz. The vertical lines mark when the antenna was nearly parallel to  $B_0$ . (d) The high-frequency wave power. The white line is  $f_{ce}$ . (e) The perpendicular, low-frequency wave power. The white line is  $f_{CH+}$ .

have three axes of electric field ( $E$ ) and one axis of magnetic field ( $B$ ) at  $0.5 \mu\text{s}$  resolution. We also present data from a relatively new instrument called a Plasma Wave Tracker, which measures the power within a  $\sim 16$  kHz frequency band which includes  $f_{ce}$ . It has  $\sim 100$  Hz frequency resolution and  $\sim 16$  ms time resolution.



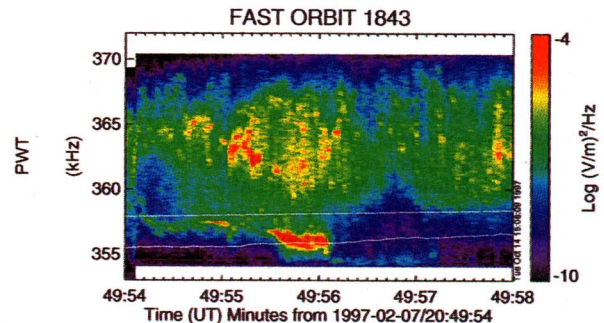
**Figure 3.** The AKR source region spectrum of  $E$  and  $B_c$  which were nearly identical and difficult to distinguish. The inter-calibration of  $E$  and  $B$  was made from free space emissions. The FWHM of the peak was less than 100 Hz.



**Figure 4.** The AKR source region. (a) The >200 kHz, RMS amplitude of the  $E$ . (b) Wave power as measured by the Plasma Wave Tracker. The white lines are  $f_{ce}$  (above) and  $f_{cr}$  (below) with an uncertainty of  $\pm 400$  Hz. (c) The high-frequency electric field wave power. The white line is  $f_{ce}$ . (d) The perpendicular, low-frequency wave power. The white line is  $f_{H+}$ . (e) The electron energy flux versus energy. The fluxes at <60 eV were photoemissions.

## Observations

Figure 1 shows high-resolution data covering  $\sim 40$  s of a near-midnight auroral crossing at  $\sim 3500$  km altitude. The display covers only a small portion of the auroral zone. Magnetometer data (not shown) indicate that the satellite was in the upward current region. We start by examining the D. C. electric field (panel a) and Langmuir probe current (panel b). There were two conspicuous periods,  $\sim 18:58:36$  UT to  $\sim 18:58:40$  UT and  $\sim 18:58:55$  UT to  $\sim 18:59:02$  UT, which had strong electric field turbulence and a noticeable decrease in the Langmuir probe current. The decrease in the Langmuir probe current signifies the auroral density cavity. The D.C. electric fields had electrostatic shocks which indicate that the satellite crossed a potential structure.



**Figure 5.** Plasma Wave Tracker observations show emissions near the relativistic electron cyclotron frequency ( $f_{cr}$ ).

Panels f and g display electron energy flux versus energy and angle. There were  $\sim 1$  keV to  $\sim 10$  keV precipitating electrons throughout (panel f). In the two density cavities there were no electron fluxes below  $\sim 1$  keV; fluxes below  $\sim 60$  eV were spacecraft photoelectrons. The distribution was mostly isotropic (panel g) with reduced fluxes near  $180^\circ$  pitch angle (up-going) due to the ionospheric loss cone. There were, on several occasions, enhanced electron fluxes at  $90^\circ$  and  $270^\circ$  pitch angles. The most noticeable events were at  $\sim 18:58:54$  UT and  $\sim 18:59:03$  UT. The enhanced perpendicular electron fluxes had lower energies than the precipitating electrons and therefore were trapped [Delory *et al.*, 1998].

Up-going ion beams were seen in both density cavities. The ion energy in the first cavity was below 1 keV. In the second cavity it reached several kilovolts. Outside of the cavities, there were weak fluxes of transversely heated ions. The decreased density, electrostatic shocks, depleted low-energy electron fluxes, and ion beams were consistent with the spacecraft crossing into two regions that had parallel potential drops both above and below. Elsewhere, the spacecraft was below the auroral acceleration region.

Panels d and e show the electric field spectral power density as a function of frequency and time. Panel e displays the frequency span from 32 Hz to  $\sim 16$  kHz at 32 Hz bandwidth and panel d displays D.C. to 600 kHz at 15 kHz bandwidth. The white lines are  $f_{ce}$  and the H+ cyclotron frequency ( $f_{CH+}$ ). The data in panels d and e were derived from two antennae and are independent of the satellite's spin phase.

Panel c displays Plasma Wave Tracker (PWT) data. The frequency axis is from 425 kHz to 450 kHz with 64 Hz resolution. The PWT receives data from only one antenna so it is modulated by the satellite's spin. It has  $\sim 100$  Hz jitter. The data show no enhancement of wave power in the first density cavity but do show strong AKR in the second density cavity that penetrated down to and below  $f_{ce}$ . The second density cavity,  $\sim 18:58:55$  UT to  $\sim 18:59:02$  UT, was an AKR source region. It was characterized by increased wave power near the  $f_{ce}$ , decreased plasma density, an absence of low-energy electron fluxes, and an ion beam.

There were intense VLF emissions covering almost the entire 40 s period in Figure 1e. The emissions concentrated at or above the lower hybrid frequency ( $f_{lh} \sim 1$  kHz) outside of the density cavities.  $f_{lh}$  abruptly decreased to near the  $f_{CH+}$  in the cavities.

Figure 2 shows an expanded view of the AKR source region. The black trace in panel a is the total electron density ( $n_{tot}$ ) calculated from fits of  $E_{\parallel}^2/E_{\perp}^2$  to the resonance cone angle of the whistler/Langmuir dispersion over the 1 kHz - 16 kHz frequency range [Strangeway *et al.*, 1998]. The fits used cold fluid dispersion with a thermal correction. We estimate an uncertainty of  $\sim 25\%$  for densities greater than  $4 \text{ cm}^{-3}$ , and  $\sim 50\%$  for densities of  $\sim 1 \text{ cm}^{-3}$ . Fits with an RMS deviation greater than 50% were omitted. Strong thermal effects, electron acoustic emissions, or beam mode emissions could cause deviations in the fits.

The red trace in Figure 2a is the density of  $> 100$  eV electrons measured by the electron electrostatic analyzer. This trace represents the "hot" electron density ( $n_{hot}$ ) with an accuracy of  $\sim 20\%$ . In this example, the electron density appears to have been primarily supported by energetic electrons,  $n_{tot} \approx n_{hot}$ . The data are consistent with no thermal core below  $\sim 100$  eV.

Figure 2b displays the RMS amplitude of the  $> 200$  kHz, omnidirectional electric field signal averaged over  $\sim 50$  ms. The signal is dominated by AKR and has no satellite spin dependence and therefore represents the AKR amplitude "envelope". The AKR in this example reached amplitudes as high as 200 mV/m. The most intense emissions came in bursts of less than 1 s, during which the satellite traveled less than  $\sim 6$  km.

Panels c, d, and e in Figure 2 are expanded views (15 s) of those in Figure 1.  $f_{ce}$  was calculated from the measured magnetic field with an uncertainty of  $\pm 0.1\%$  ( $\pm 400$  Hz). Panel c, the PWT data, has several added features. The magenta lines indicate when the electric field antenna were nearly parallel to the magnetic field. The measured power was much lower at these times indicating the electric field signal was primarily perpendicular. The right cutoff,  $f_{right} = (f_{ce} + (f_{ce}^2 + 4f_{pe}^2)^{1/2})/2$ , where  $f_{pe}$  is the electron plasma frequency, was calculated from the total density in panel a and the measured magnetic field. The relativistic electron cyclotron frequency,  $f_{cr} = eB/m\gamma$ , corrects for the average energy of the  $> 100$  eV population in measured electron distribution. Since the bulk of the electron distribution was nearly mono-energetic,  $f_{cr}$  represents the electron cyclotron frequency of the hot electron population.

There are several interesting features in the AKR source region spectra. The emissions, at times, noticeably extended below  $f_{ce}$  (for example,  $18:59:01.2$  UT), but not below  $f_{cr}$ . Outside of the density cavity, the emissions seemed to obey the right cutoff. The AKR at  $\sim 18:58:54.75$  UT and  $\sim 18:58:55.75$  UT was at the boundary of the cavity. This particular event is strikingly similar to that reported on the Viking satellite [Pottelette *et al.*, 1992]. The density gradient had discrete steps and the most intense AKR was before the density reached its lowest value.

Figure 3a shows electric and magnetic field spectra superimposed. The spectra were calculated from part of a  $\sim 0.25$  s digital waveform capture starting at  $\sim 18:58:56.18$  UT (marked on Figure 2c). The  $E$  and  $B$  spectra were remarkably identical and difficult to distinguish. The FWHM bandwidth of the peak was quite narrow ( $\sim 100$  Hz). The  $|E|/|B|$  ratio (Figure 3b) showed no feature that would identify  $f_{ce}$ , the upper hybrid resonance, or  $f_{right}$ . There was, however, a sharp reduction in wave power below  $f_{cr}$ .

During the digital capture ( $18:58:56.317$  UT), one set of antennae was nearly parallel ( $1.3^\circ$ ) to the ambient magnetic field ( $B_0$ ). At that time, the AKR was peaking at  $> 300$  mV/m. The ratio  $E_{\parallel}/E_{\perp}$  was 0.040 indicating  $E$  was within  $2.3^\circ$  of perpendicular to  $B_0$ . The average  $E_{\parallel}/E_{\perp}$  from point by point de-spinning of the 3-axes digital data over the  $\sim 0.25$  s capture was 0.063 ( $3.6^\circ$ ). Within error, the wave polarization is consistent with  $k_{\parallel} = 0$ .

An AKR source region crossing during a second orbit is displayed in Figure 4 which plots 40 s of the near-midnight auroral zone. After  $20:49:46$  UT, the electric field emissions near  $f_{ce}$  intensify (panels a, b, and c) and there were no measurable electron fluxes below  $\sim 1$  keV (panel e). The total plasma density, calculated from the low-frequency wave emissions, again was consistent with the hot ( $> 100$  eV) plasma density measured by the electrostatic analyzers. The amplitude envelope (panel a) again showed intense bursts reaching almost 200 mV/m that often lasted less than one second. The majority of the AKR emissions (panel b) were 5 kHz to 10 kHz above  $f_{ce}$ , but the most intense burst ( $\sim 20:49:56$  UT) had a spectral peak measurably below  $f_{ce}$  and within error, at  $f_{cr}$ . Figure 5 shows an expanded view of panel b during the intense burst.

## Discussion and Conclusions

Over one hundred orbits were inspected for AKR and more than ten AKR source regions have been examined. We confirm that the most intense AKR occurs in density depleted cavities extending 30 km to 300 km in latitude. All of the source regions were near magnetic midnight and close to the apogee of FAST (4175 km).

Observations from the FAST satellite have detailed the AKR source region with orders of magnitude higher time and frequency resolution than previous missions. As a result, the source region spectra are fully resolved and several new features of the AKR source region have emerged. Electron distributions had precipitat-

ing, energetic electrons with a loss cone and up-going ion beams were always present. The most outstanding features were a void of low-energy ( $< 1$  keV) electron fluxes and an enhanced trapped population. High-resolution data show that the trapped electrons were near the edge of the density cavity rather than inside, clarifying previous reports [Roux *et al.*, 1993].

Detailed examinations revealed four examples of source region densities ( $\sim 1 \text{ cm}^{-3}$ ) that were dominated by  $> 100$  eV electrons, consistent with  $n_{\text{hot}} > n_{\text{cold}}$  (the total density was not determined in the other orbits). How often this occurs is under investigation [Strangeway *et al.*, 1998]. FAST observations thus far indicate  $f_{pe}/f_{ce} < 0.05$ , which is less than previously reported [Hilgers, 1992].

Strong temporal and/or spatial variations were in every source crossing that was examined. The strongest emissions were often confined to the edge or to a small region in the density cavity and typically appeared in bursts  $< 1$  s. The source region did not extend over the entire density cavity. The AKR source had an average spectral power density roughly one order of magnitude higher, and peak power typically two to three orders of magnitude higher than reported by Viking [de Feraudy *et al.*, 1987]. Peak electric field amplitudes were often greater than 100 mV/m and reached 300 mV/m. Source region amplitudes were such that the ratio of electric field energies to thermal energies ( $W = \epsilon_0 E^2 / 2 n_e T_e$ ) were as high as  $2.5 \times 10^{-3}$  which may enable nonlinear wave-wave interactions. The source emissions had bandwidths often  $< 1$  kHz and occasionally  $\sim 100$  Hz. Outside of the source region, the FAST observations have comparable amplitudes to those reported earlier and remote, longer time-scale observations show similar rising and falling tones.

The most intense source region emissions were within  $\sim 2\%$  of  $f_{ce}$  and occasionally below  $f_{ce}$ , but almost always above  $f_{cr}$ . Outside of the density cavity, a clear  $R-X$  cutoff was seen. The  $|E|/|B|$  ratio showed that the source region emissions were electromagnetic, remarkably featureless, and revealed no indication of  $f_{ce}$ ,  $f_{\text{right}}$ , nor the upper hybrid resonance. Only rarely was a possible electrostatic component observed. In one example, the wave polarization was shown to be  $< 3.6^\circ$  from perpendicular to  $B_0$  and, within error, consistent with  $k_{\parallel} = 0$ . This examples extends an earlier analysis which found the source polarization within  $10^\circ$  of  $B_0$  [Hilgers *et al.*, 1992]. These findings, the strongly depleted ( $n_{\text{hot}} > n_{\text{cold}}$ ) density cavity, the near-perpendicular polarization, the emissions extending below  $f_{ce}$  to  $f_{cr}$  and the constant ratio  $|E|/|B|$  are consistent with strong modification of the plasma dispersion from the cold wave approximation [Pritchett and Strangeway, 1985], and wave growth below  $f_{ce}$  with  $k = 0$  [Delory *et al.*, 1988].

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