

Overview of Transient Luminous Events

(Red Sprites, Elves, Blue Jets, Halos, Pixies, Trolls, etc.)

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University of California Berkeley

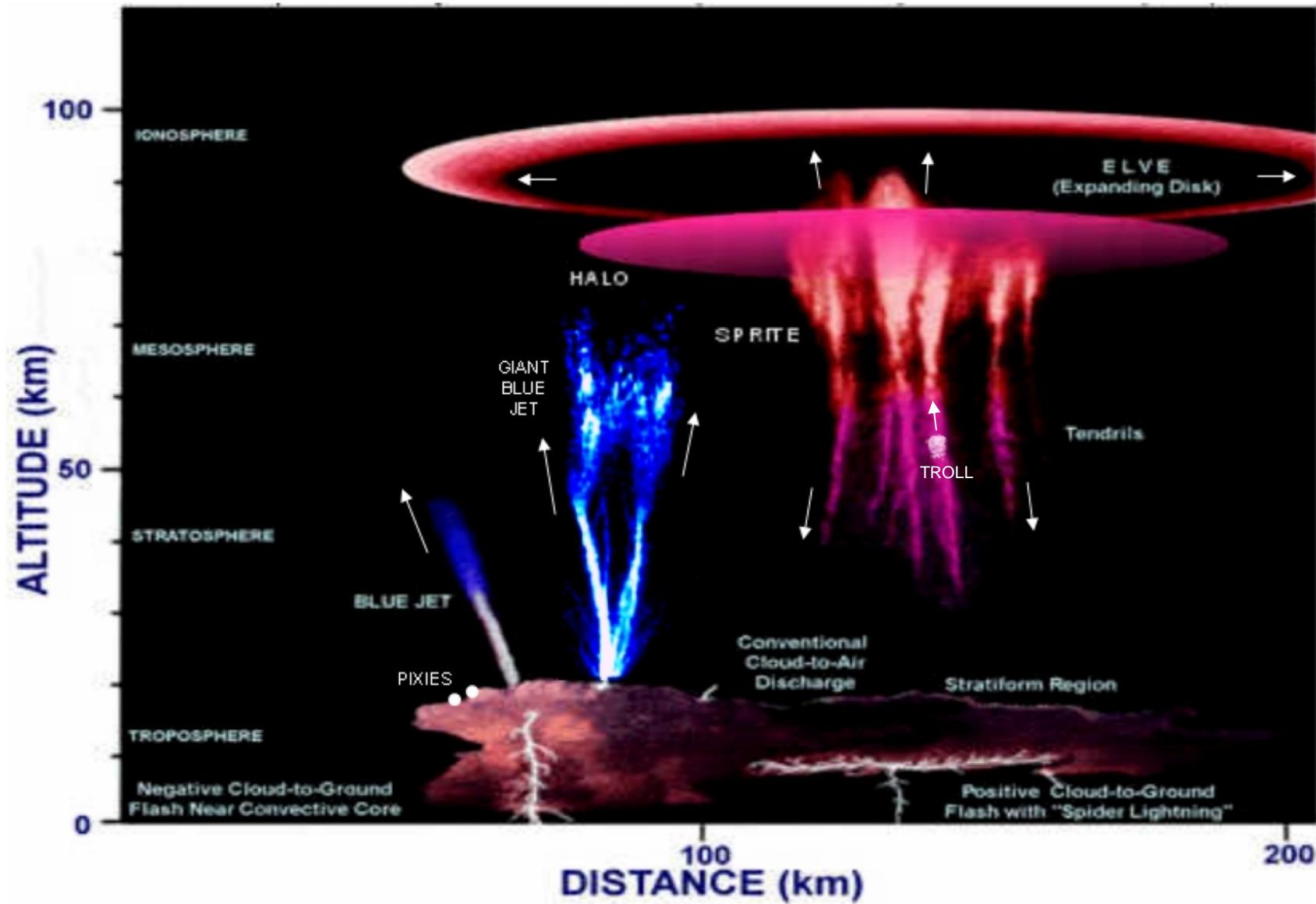
15 February 2005

Outline

1. Observations
2. General Characteristics of the Major Classes of Events
3. Mechanisms

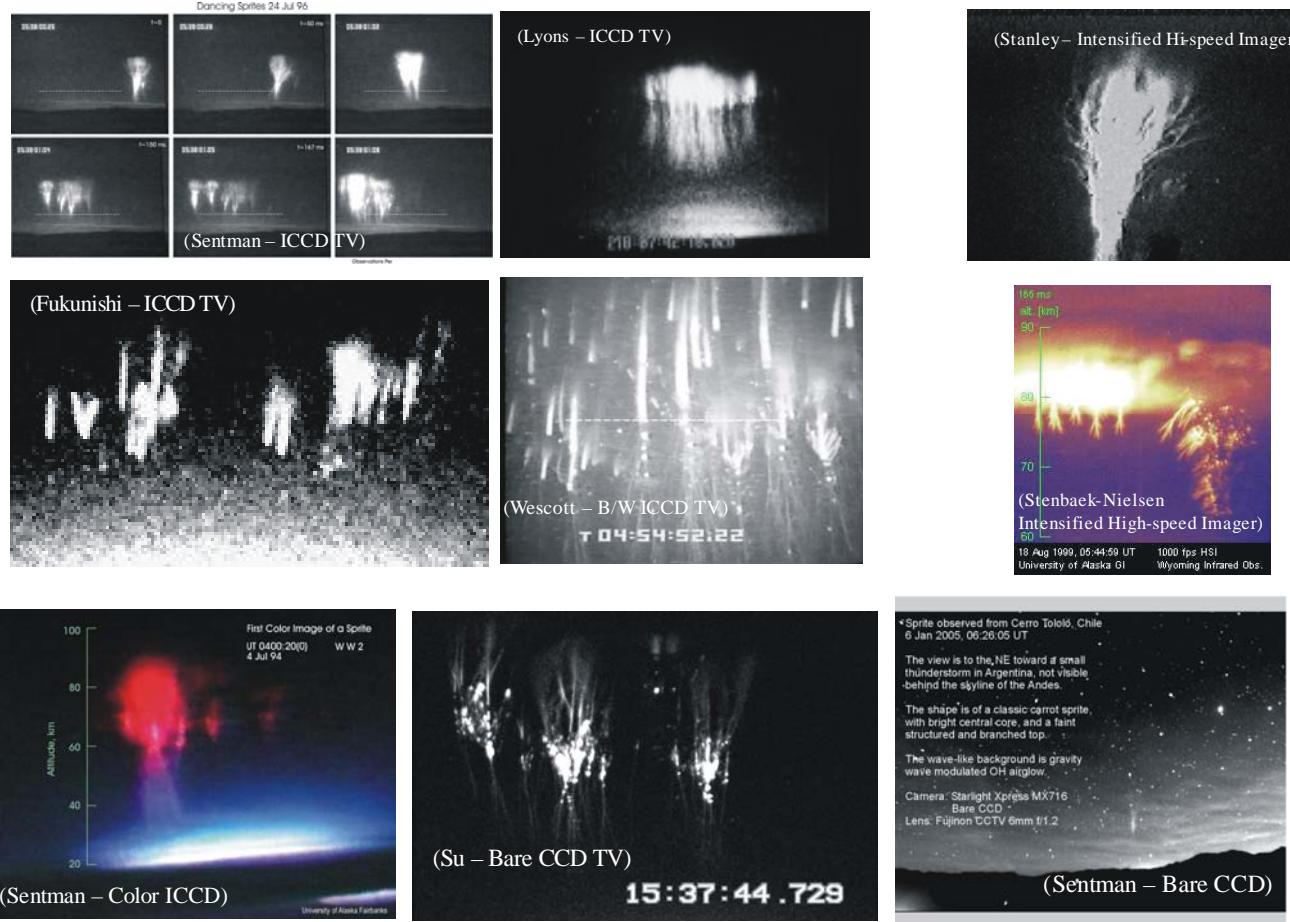
A bibliography of selected papers is included.

Varieties of Transient Luminous Events in the Upper Atmosphere



(Elaboration of figure by Lyons et al. 2000)

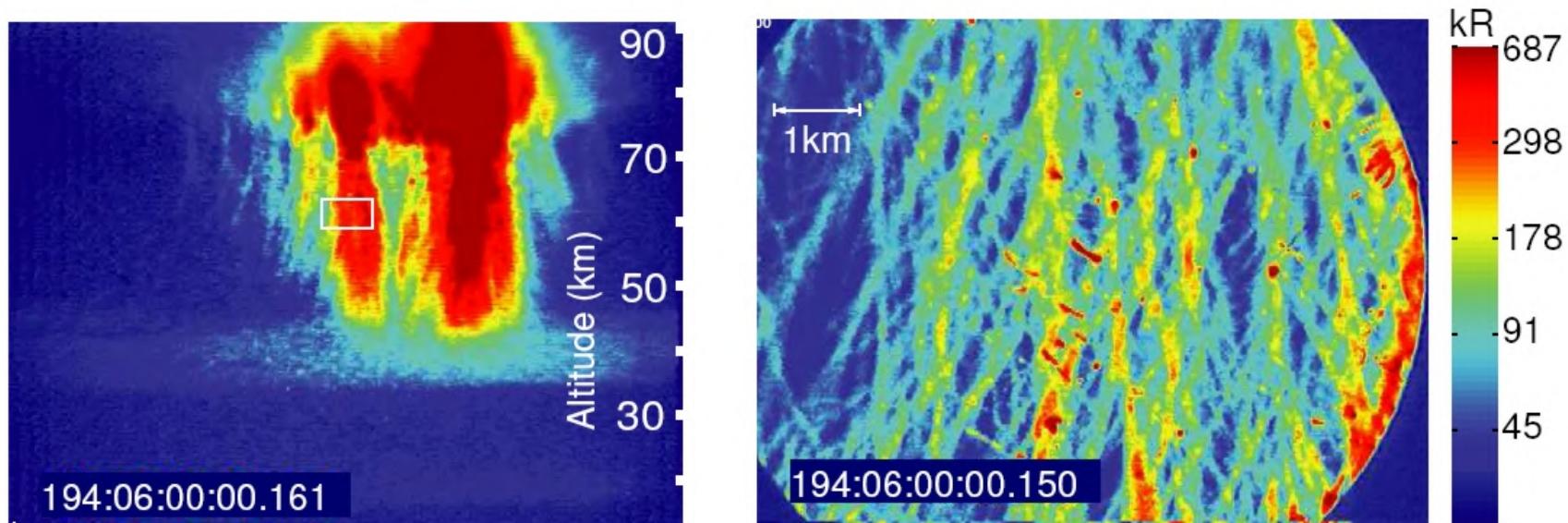
Sprite Gallery – Images From a Variety of Sources Obtained With a Variety of Cameras



Early sprite imagers were intensified CCD TV cameras. Recent research has simultaneously moved in two directions: (1) high speed >1000 fps cameras, and (2) inexpensive bare CCD imagers, both TV and integrating systems. Sprites are bright enough ($\gg 1$ MR) that both types of imagers work well.

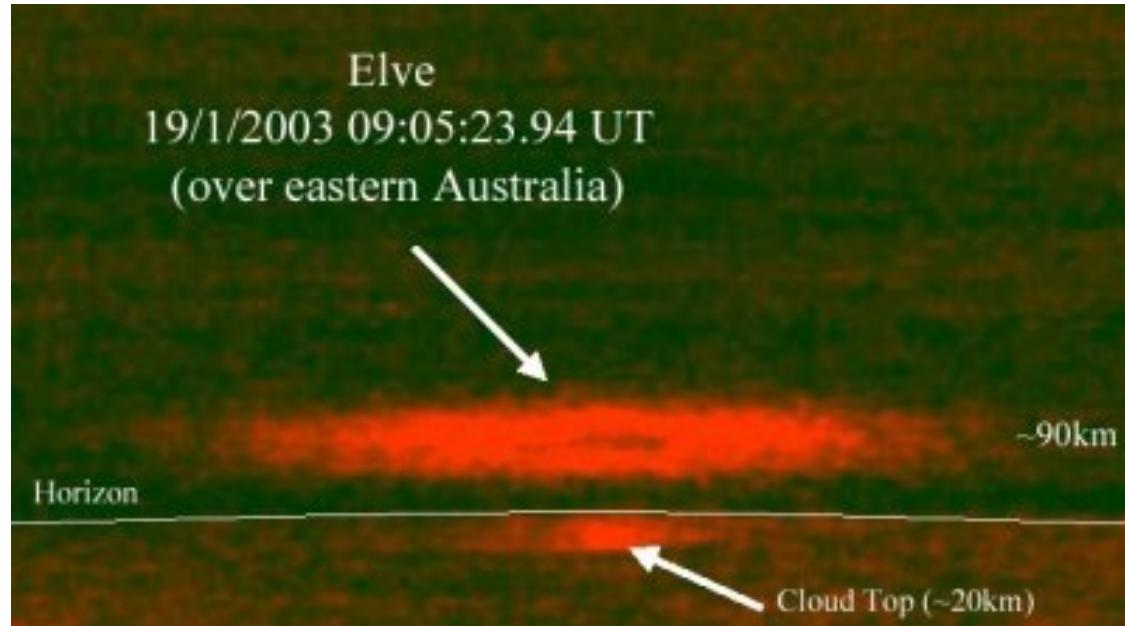
Telescopic Imaging of Sprites

- Wide (left panel) and narrow (right panel) field of view images of a bright sprite event [*Gerken et al.*, GRL, 27, 2637, 2000]:

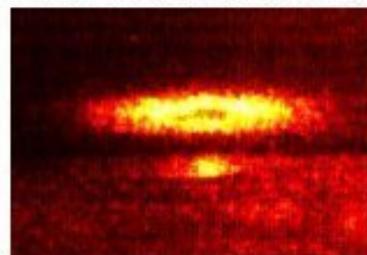


- The measured streamer diameters are 61-145, 150, 196 m, for altitude ranges 60-64, 76-80, 81-85 km, respectively [*Gerken et al.*, 2000].

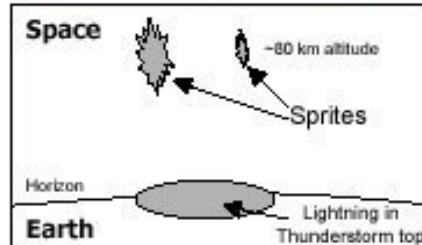
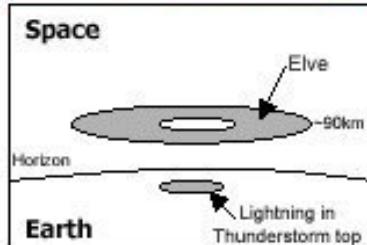
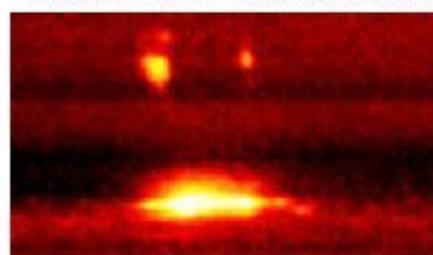
STS-107 Observations of ELVES and Sprites



Elve over South Pacific

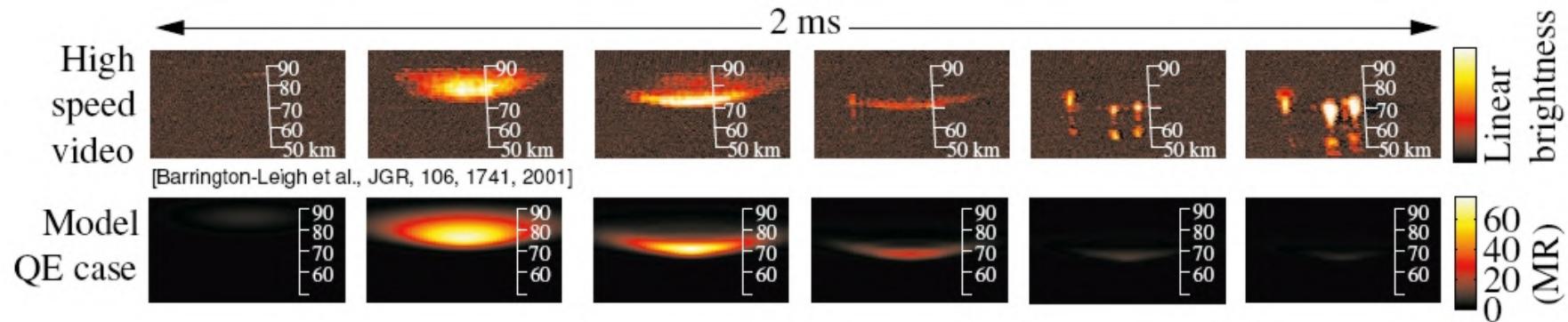


Sprites over SE Australia



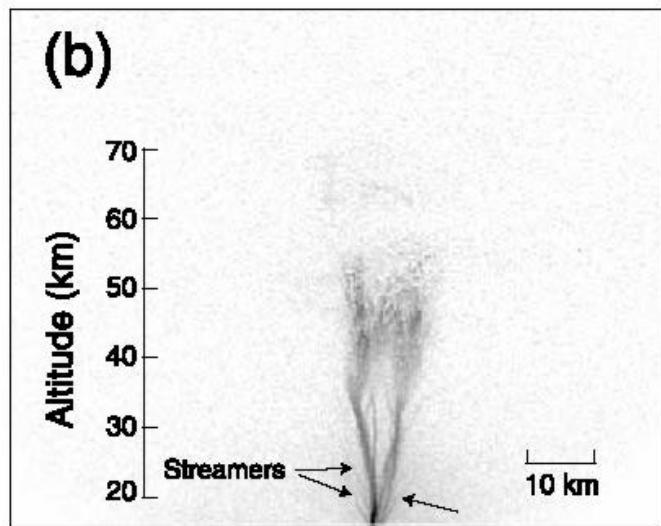
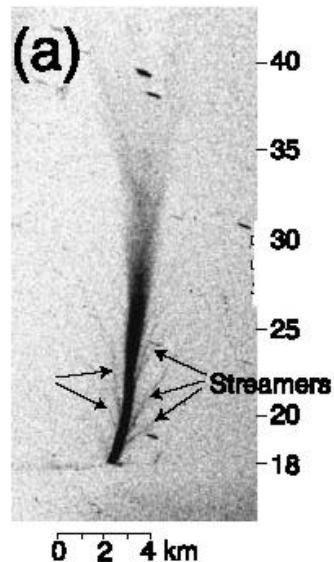
Sprite Halos

- Barrington-Leigh *et al.* [JGR, 106, 1741, 2001] conducted one-to-one comparison between high-speed video observations of sprites and a fully electromagnetic model of sprite driving fields and optical emissions. Sprite halos are brief descending glows with lateral extent 40-70 km, which sometimes observed to accompany or precede more structured sprites. The analysis conducted by Barrington-Leigh *et al.* [2001] for the first time identified sprite halos as being produced entirely by quasi-electrostatic thundercloud fields.



Blue Jets

Blue Jets – Positive(?) Streamers from Cloud Tops



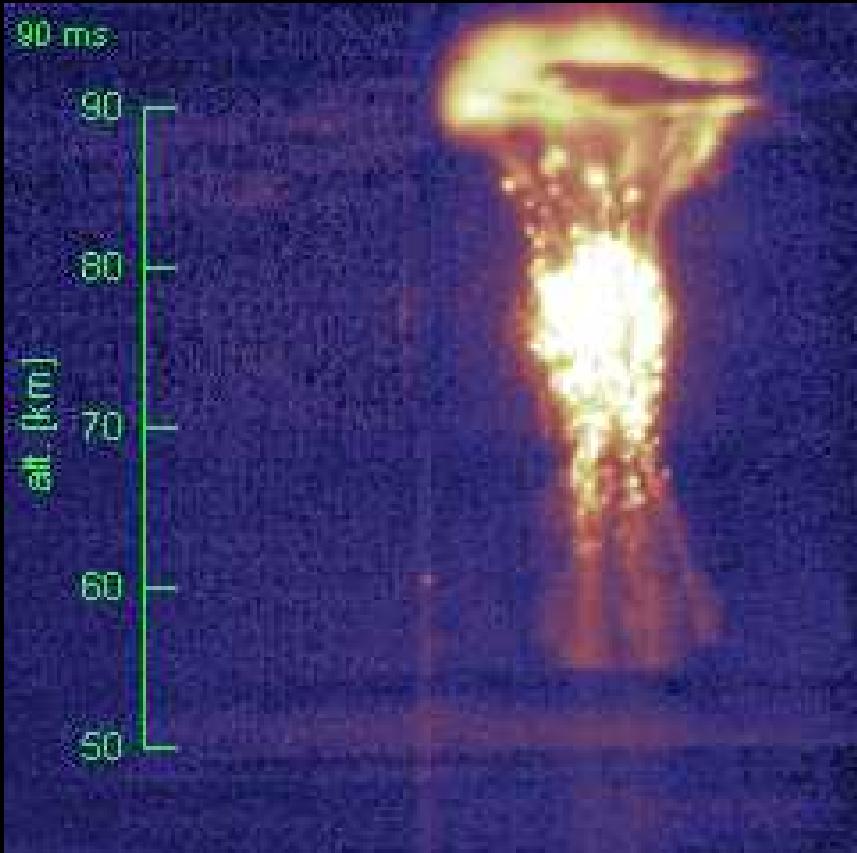
Wescott et al. 1999

Pasko et al. 2000



A Few Case Studies of Sprites

TLEs exhibit a complex variety of forms with considerable variation of structure and dynamics among individual events. Most of the various forms that have been given names (sprites, c-sprites, elves, halos, trolls), as well as numerous other features that have not yet been studied, are present in the following time-lapse sequences of sprites recorded at 1000 fps on 18 August 1999 as part of the NASA Sprites99 Balloon Campaign.



18 Aug 1999, 04:49:20 UT 1000 fps HSI
University of Alaska GI Wyoming Infrared Obs.

044920_B.avi

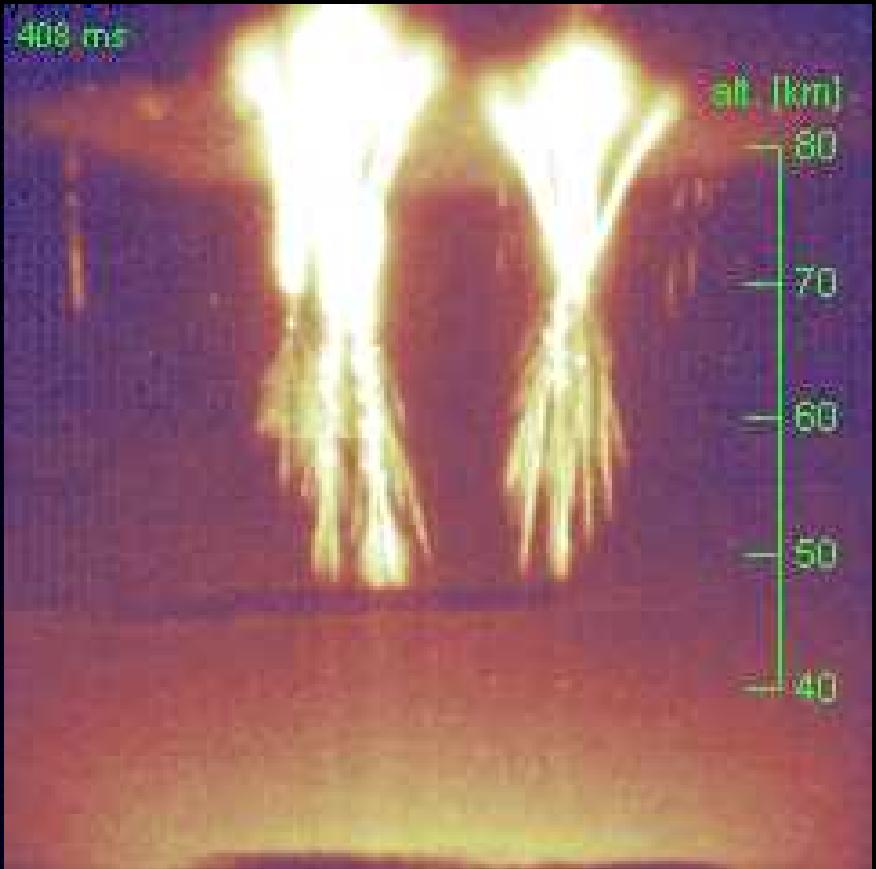
Halo followed by offset double sprites,
weak downward “smoke” billows
(clip)



18 Aug 1999, 05:05:43 UT 1000 fps HSI
University of Alaska GI Wyoming Infrared Obs.

050543_B.avi

Halo, sprites, downward “smoke” billows
(clip)

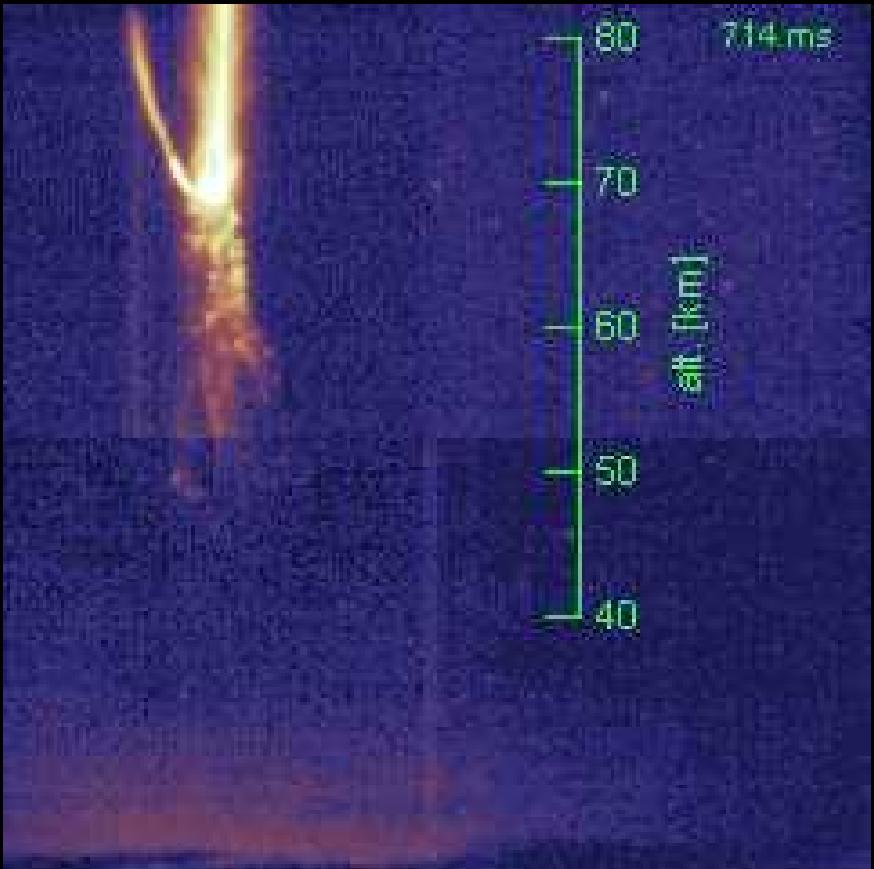


18 Aug 1999, 05:17:11 UT 1000 fps HSI
University of Alaska GI Wyoming Infrared Obs.

051711_B.avi

Halo, Sprites, persistent ascending
“trolls” in tendrils

[\(clip\)](#)

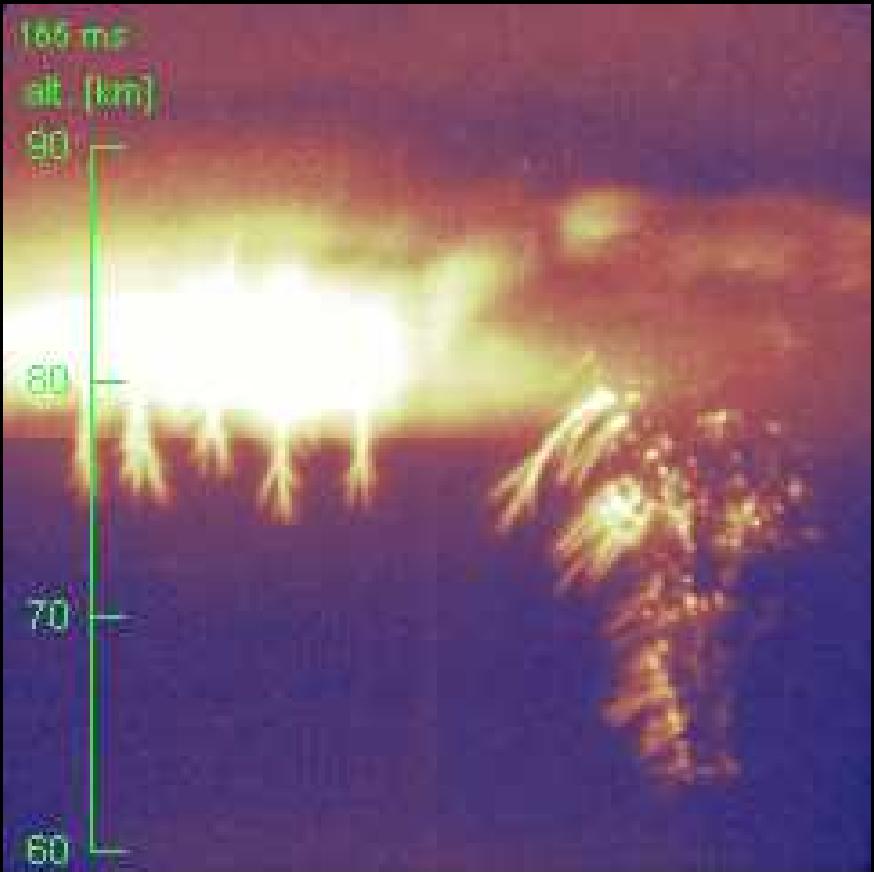


18 Aug 1999, 05:10:34 UT 1000 fps HSI
University of Alaska GI Wyoming Infrared Obs.

051034_B.avi

Sprites, ascending “trolls,” second sprite
and rebrightening, weak “smoke” billows

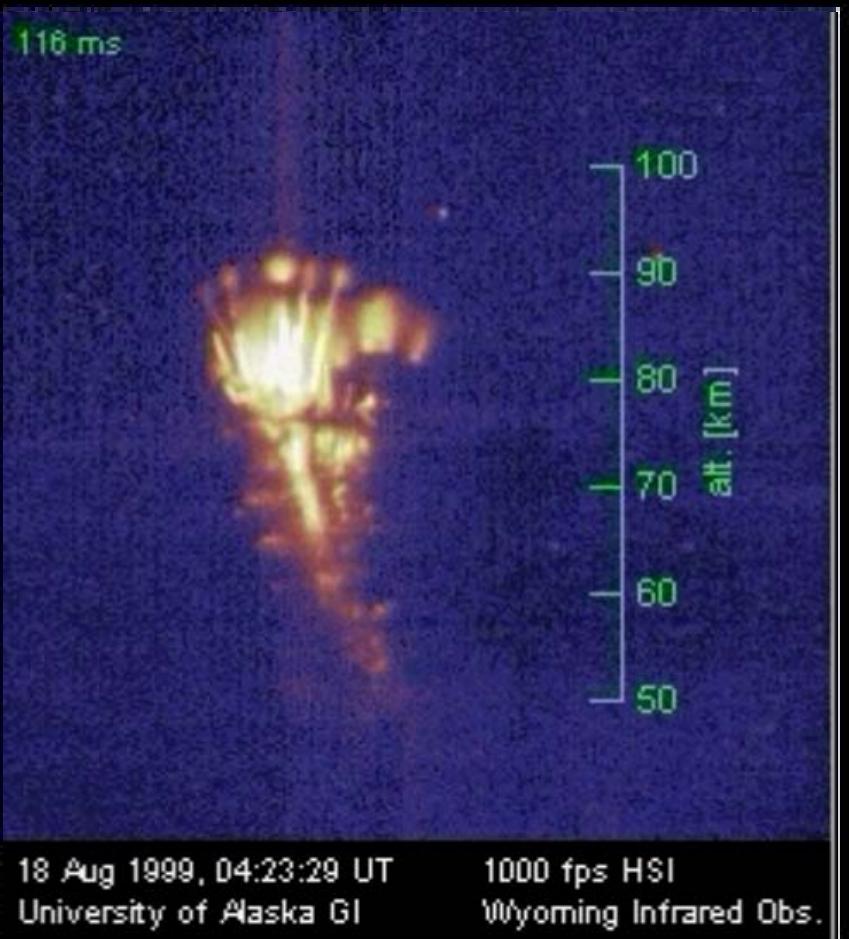
[\(clip\)](#)



052417_B.avi

Complex sprite followed by “palm tree”
ascending from below, sideways
discharges, “smoke” billows
[\(clip\)](#)

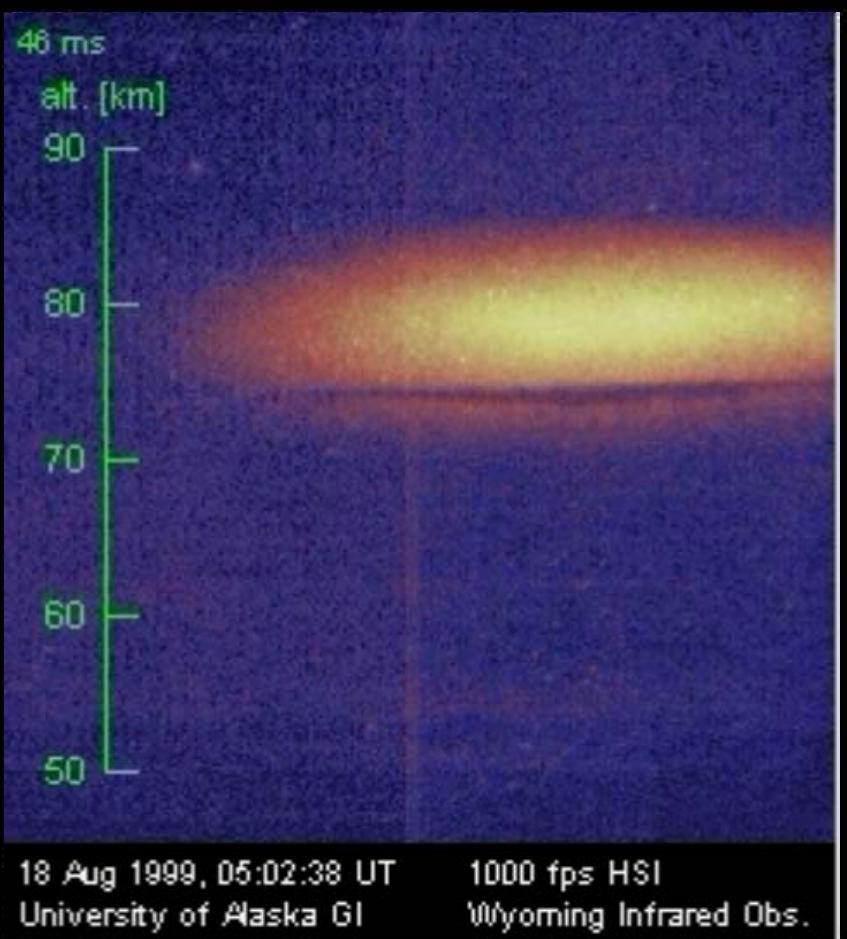
Fast downward precursor streamer
preceding sprite, second upward
precursor, retracing in second sprite
features in first sprite.
[\(clip\)](#)



042329_B.avi

Sprite, downward “smoke,” late developing upward “stalks.”

[\(clip\)](#)



050238_B.avi

Simple isolated sprite halo.

[\(clip\)](#)

Observations from Chile of Sprites over Argentina Storms

Observations of sprites from Cerro Tololo in the early morning hours of 15 Jan 2005 UT looking eastward over the Andes. The source of the sprites is a small convective system in Argentina at a distance of ~500 km. A crescent moon illuminates the landscape early in the clip. OH nightglow and gravity waves are prominent. About two dozen sprites are visible just above the mountain ridge on the right half of the images. (Shown here: first frame of video [clip](#))



15-Jan-2005 01:48:57 UT
Cerro Tololo, Chile

MX716 9mm f/1.4
University of Alaska

Camera 1:

Starlight Xpress MX716
CCD Sony ICX249AL
752x580x8.6 μ m
16 bit samples, 2x2 binning
10s integration, 4s readout
unfiltered

Lens:

Fujinon CCTV 9mm f/1.4
FOV: 40Hx30V deg

Control:

IBM Thinkpad
IEEE 802.11b wireless
MaxIm DL

Camera 2:

KT&C PC164C TV camera
640x480x6.8 μ m
30 fps NTSC

Lens:

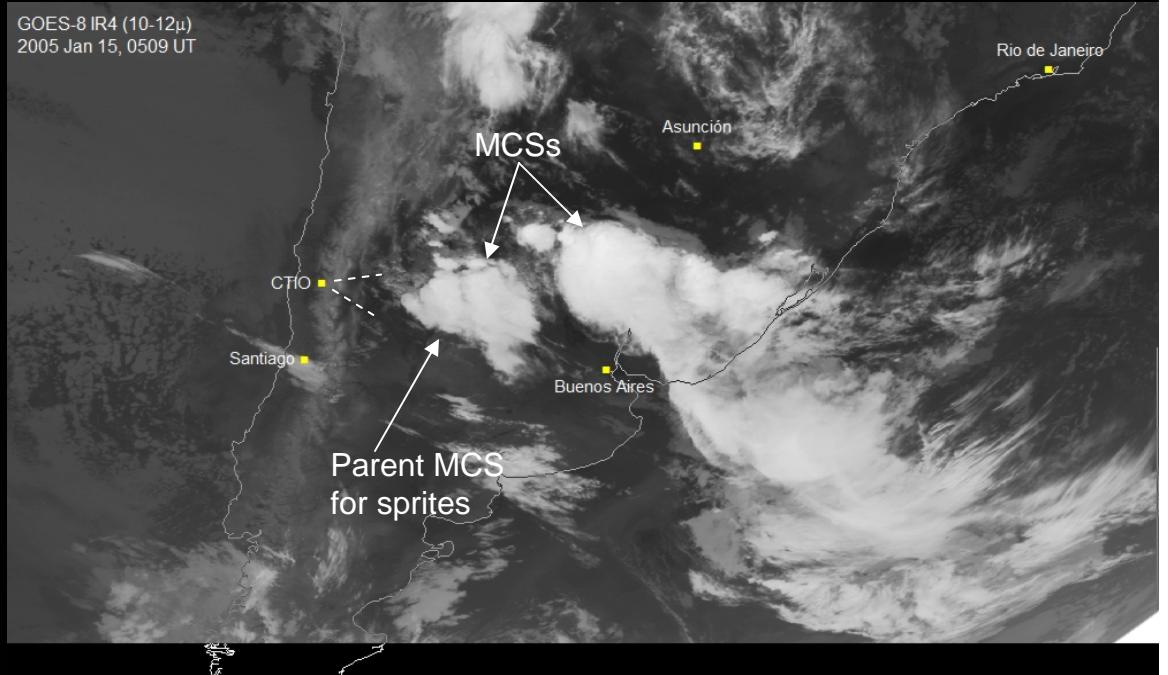
Fujinon CCTV 6mm f/1.2

Recorder:

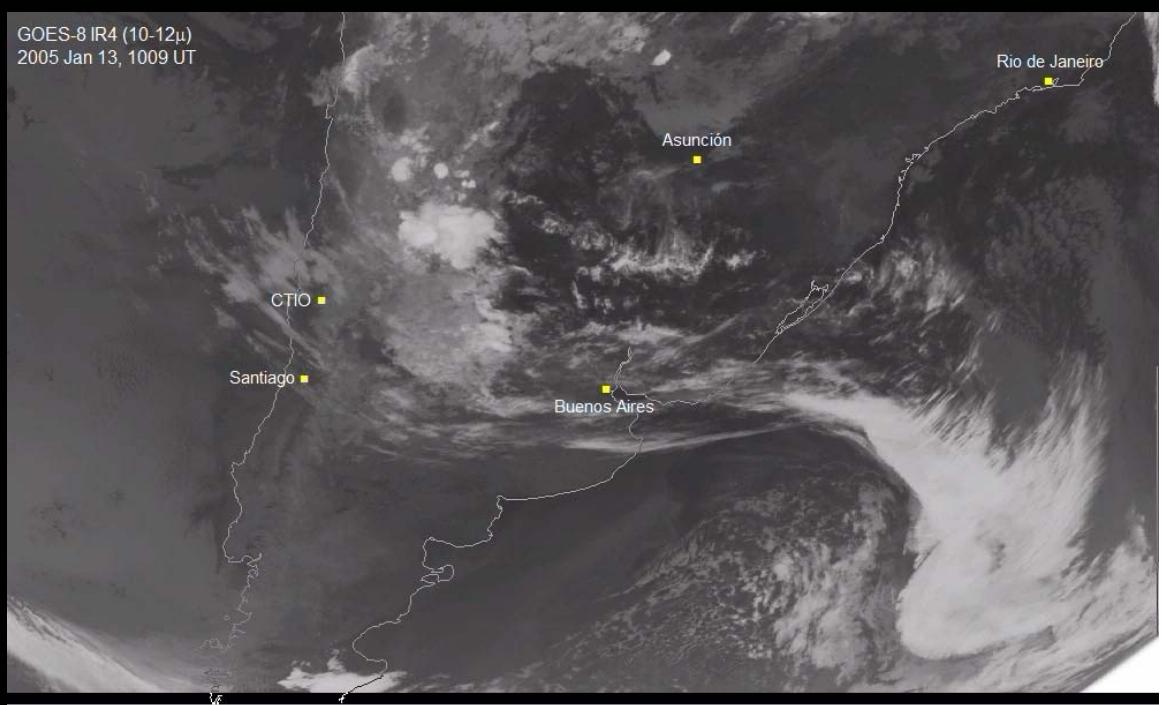
Sony Digital Video Walkman
miniDV recorder

GOES-8 IR4 (10-12 μ m) weather images of southern South America at the time of the observations.

Above: 15 Jan 0509 UT, at the time of maximum sprite activity. There are two moderately sized Mesoscale Convective Systems (MCSs) in Argentina; one ~500 km to the east of CTIO and one to the north of Buenos Aires. CTIO, several large cities, and the parent storm of the sprites are marked. The FOV of the primary camera is indicated.



Below: Video clip of weather patterns beginning two days before the observations.



Spectroscopic Identification of Sprite Optical Emissions

Sprite and its emission spectra measured using a TV (30 fps) imager and slit spectrograph showing characteristic N₂ red 1PG and blue 2 PG emission line features.

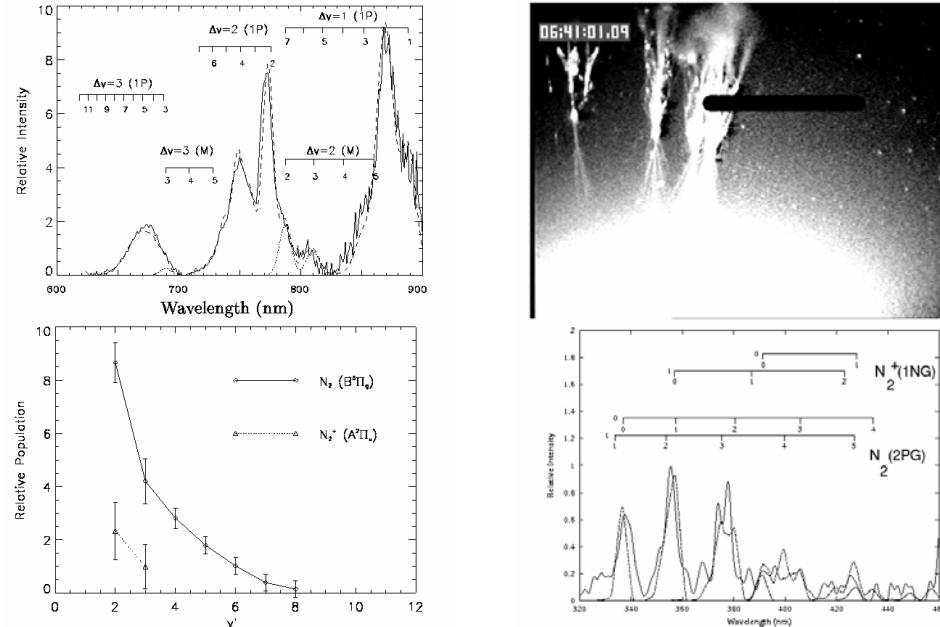
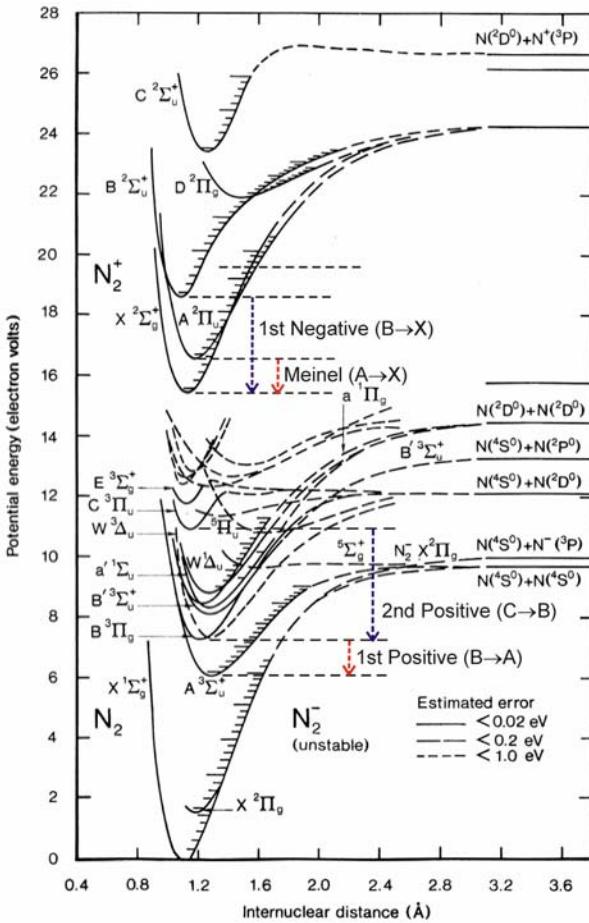


Table 1. Several neutral and ionized N₂ emissions observed in sprites, blue jets, and elves, as well as the aurora. The first observations of sprites showed only N₂(1PG) emission, which requires the lowest threshold energy of the N₂ states that emit optically. The threshold energies reported are based on filling the lowest vibration energy level of the electronic state. This table is based on Table 4.7 of Vallance Jones [1974].

| Name | Upper State | Lower State | Lifetime | Quench Alt. | Energy |
|-----------------------------------|---|---|----------|-------------|----------|
| N ₂ (1PG) | N ₂ (B ³ Pi _g) | N ₂ (A ³ Sigma _u ⁺) | 6 μs | 53 km | 7.50 eV |
| N ₂ (2PG) | N ₂ (C ³ Pi _u) | N ₂ (B ³ Pi _g) | 50 ns | 30 km | 11.18 eV |
| N ₂ ⁺ (1NG) | N ₂ ⁺ (B ² Sigma _u ⁺) | N ₂ ⁺ (X ² Sigma _g ⁺) | 70 ns | 48 km | 18.56 eV |
| N ₂ ⁺ (M) | N ₂ ⁺ (A ² Pi _u) | N ₂ ⁺ (X ² Sigma _g ⁺) | 14 μs | 85-90 km | 16.54 eV |
| N ₂ (VK) | N ₂ (A ³ Sigma _u ⁺) | N ₂ (X ¹ Sigma _g ⁺) | 2 s | 145 km | 6.31 eV |

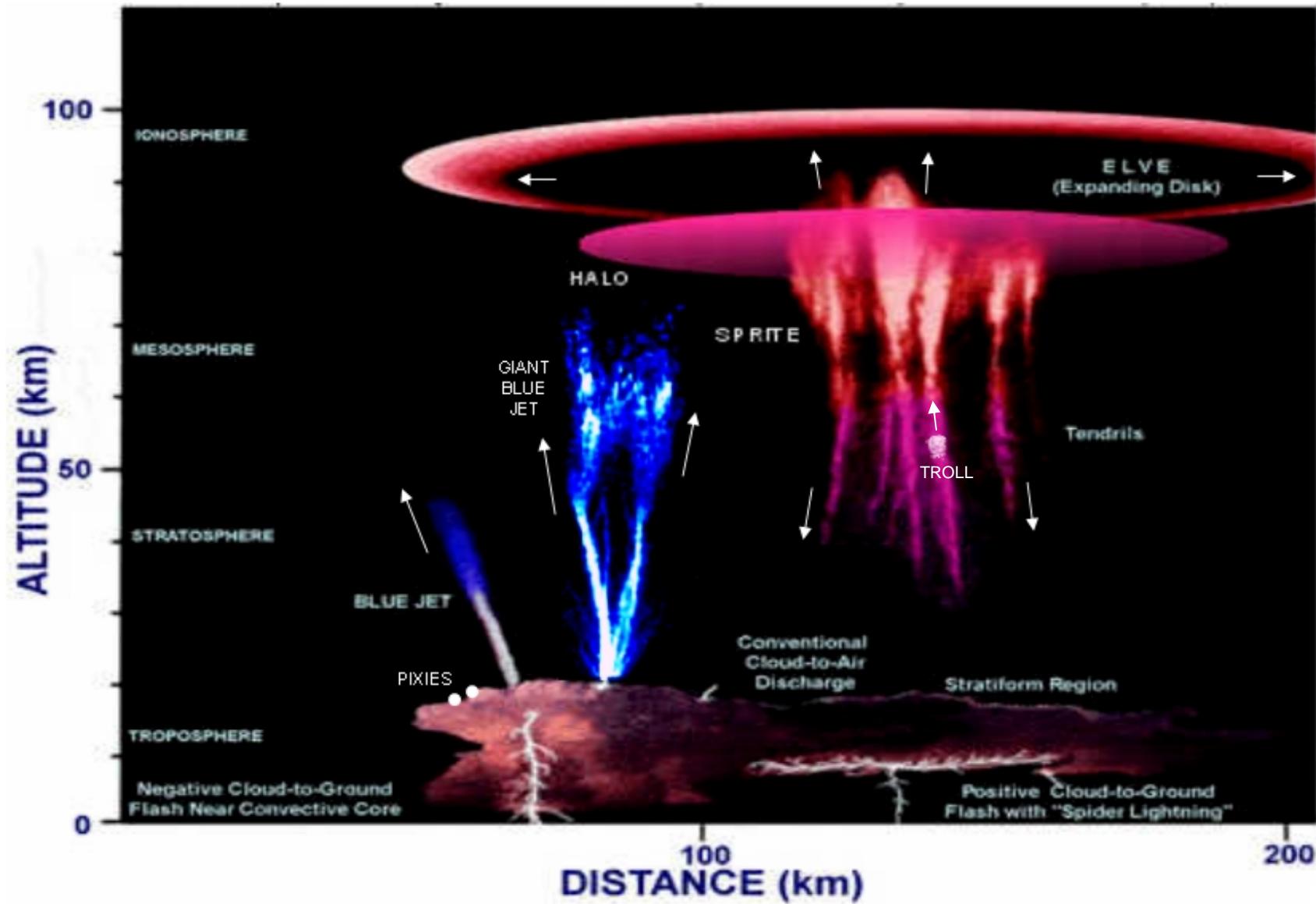
From Heavner et al. (2000)

N₂ energy diagram.



A noteworthy and unexpected feature of the measured spectrum is the absence of strong ion lines. This indicates that most of the optical brightness originates in neutral excited electronic states, although high speed photometric observations have detected a brief (< 1ms) burst of N₂⁺ 1NG ion emissions at sprite onset .

Varieties of Transient Luminous Events in the Upper Atmosphere



(Elaboration of figure by Lyons et al. 2000)

| Principal Types of Transient Luminous Events in the Upper Atmosphere Associated with Thunderstorms/Lightning | | | | | | |
|---|-----------------|-----------------------|--|---|----------|----------------------------------|
| Type of TLE | Altitude Regime | Transverse Dimensions | Spatial Characteristics | Apparent Motion | Duration | Inventory of Observations (est.) |
| Sprites | ~ 50-90 km | ~1-20 km | Top (>80 km) diffuse Bottom (<70 km) structured | Top-upward Bottom-downward | few ms | >10,000 |
| Elves | ~ 100 km | > 100 km | Diffuse | Lateral Expansion | few ms | 100s |
| Blue Jets | ~ 18-45 km | few km | Structured | Upward | 100s ms | <100 |
| Giant Blue Jets | ~ 18-75 km | few km | Structured | Upward | 100s ms | <10 |
| Halos | ~ 75 km | ~ 50 km | Diffuse | Downward | ~ ms | 1000s |
| Trolls | ~ 60-70 km | ~ kms | Structured | Upward (Within decaying sprite tendrils) | 100s ms | 100s |
| Pixies | ~ 15-18 km | ~ 100s m | Compact | Stationary (Stormcloud tops) | 100s ms | 10s(?) |

Low Energy (few eV) Electron Fluid Model for Gas Discharges

Diffusion-drift equations – Describe a wide variety of electrical discharges in neutral gases.

$$\begin{aligned}\frac{\partial n_e}{\partial t} + \nabla \cdot n_e \vec{v}_d - D_e \nabla^2 n_e &= (\nu_i - \nu_a)n_e - \beta_{ep} n_e n_p + S_{ph} \\ \frac{\partial n_p}{\partial t} &= \nu_i n_e - \beta_{ep} n_e n_p - \beta_{np} n_n n_p + S_{ph} \\ \frac{\partial n_n}{\partial t} &= \nu_a n_e - \beta_{np} n_n n_p \\ \nabla \cdot \vec{E} &= \frac{e}{\epsilon_0} (n_p - n_e - n_n)\end{aligned}$$

The first three equations are the continuity equations for electrons, positive and negative ions, respectively, and the last equation is Poisson's equation self-consistently linking the spatial structure of the electric field with the total charge distribution, including especially the charges created and lost during the discharge process. In the above, \vec{E} is the electric field, n_e , n_p , and n_n are electron, positive ion, and negative ion number densities, \vec{v}_d is the electron drift velocity, ν_z is the electron attachment frequency, β_{ep} and β_{np} are the electron-positive ion and negative-positive ion recombination coefficients, D_e is the electron spatial diffusion coefficient, S_{ph} is the electron-ion pair production function from photoionization, and e is the absolute value of the electron charge. The electron drift velocity is defined in terms of the mobility by $\vec{v}_d = -\mu_e \vec{E}$.

The Diffusion-drift model describes the initial stages of electrical breakdown, the formation and propagation of streamers, and the creation and loss of charge through ionization and recombination. The term involving the divergence of the electron current ($\nabla \cdot n_e \vec{v}_d$), which was neglected in many early sprite models, plays an especially critical role in the structural dynamics of streamers. The Diffusion-drift equations are highly non-linear and require numerical techniques to evaluate. Numerical methods that have been applied to its solution include the Corrected Flux Transport and the modified Sharfetter-Gummel algorithms, and more recently finite element (mesh) methods. Most work has been done in 1- or 2-dimensions, and it is only within the past few years that fully 3D studies have appeared.

Key References: *Raizer* (1997), *Liu and Pasko* (2004), papers by *Kulikovsky*.

Evaluation of Electron Distribution - Nonthermal

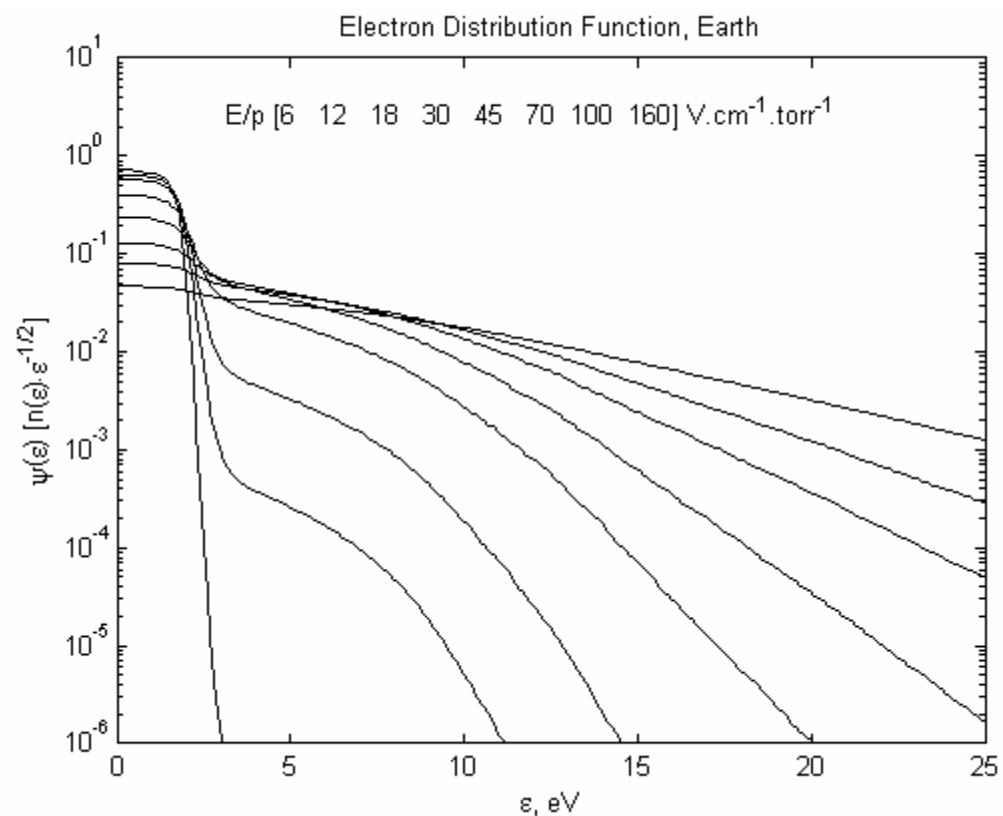
Solution of Time-Stationary Kinetic (Boltzmann) Equation (2-Term SH approximation)

$$\frac{\partial n}{\partial t} = \frac{\partial}{\partial \varepsilon} \left\{ A \varepsilon^{3/2} \frac{\partial}{\partial \varepsilon} \frac{n}{\varepsilon^{1/2}} + \underbrace{\frac{2m}{M} \varepsilon \nu_m n}_{\text{Momentum Transfer}} \right\} + \underbrace{Q(n)}_{\text{Inelastic}} = 0$$

where ε is energy, $n = n(\varepsilon)$, $A = \frac{2e^2 E^2}{3m\nu_m}$, E is electric field, and $\nu_m = \nu_m(\varepsilon)$

At low electric fields electron energy distribution $n(\varepsilon)$ in air is determined primarily by rotational interactions with N_2 and O_2 , and because of the dense number of rotational states it is approximately $n/\varepsilon^{1/2} \sim \exp(-\varepsilon/\varepsilon_0)$. At thermal energies greater than ~ 0.2 eV interactions with N_2 vibrational states become important, and the electron distribution becomes Druyvesteyn-like $n/\varepsilon^{1/2} \sim \exp(-\varepsilon^2/\varepsilon_0^2)$. At high reduced fields $\varepsilon/p > 10$ V/cm/torr that are a significant fraction of breakdown fields ($\varepsilon/p \sim 40$) a high energy tail begins to form above the $\varepsilon \sim 2-4$ eV barrier in the N_2 vibrational cross sections. Electronic excitation of atomic and molecular species, ionization, molecular dissociation, and electron attachment-detachment are driven by the nonthermal tail of the electron distribution.

The form of this distribution function shown on the right is characteristic of nitrogen. Other gases possess different equilibrium distributions.



Energy dependent electron collisions with neutral species play a fundamental role in determining the quasi-equilibrium EEDF.

Elastic collisions isotropize EEDF about the electric field direction. Inelastic collisions involving vibrational and electronic states of neutral species occur at energies greater than a few eV.

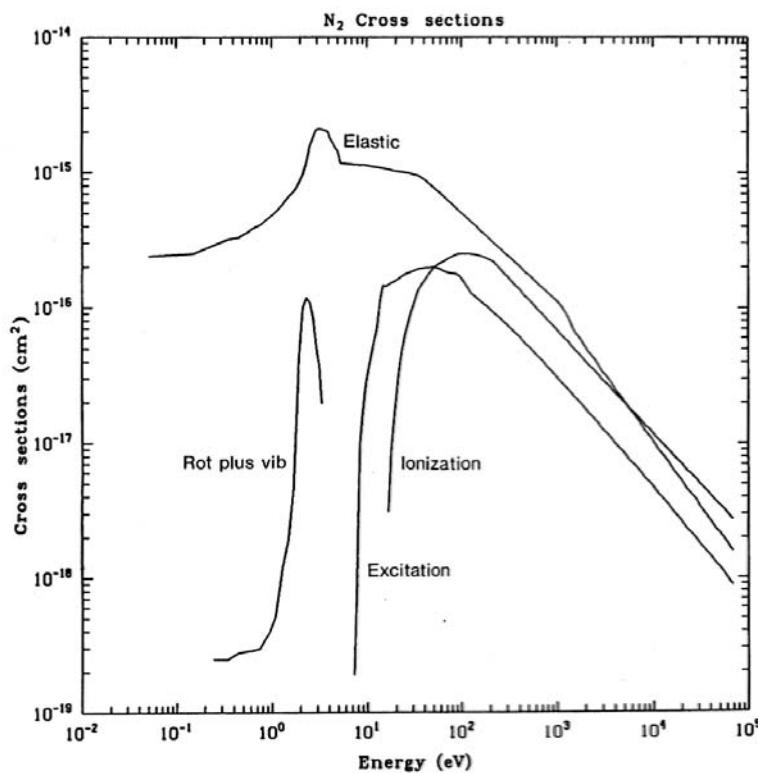


Fig. A4.1 N₂ cross sections.

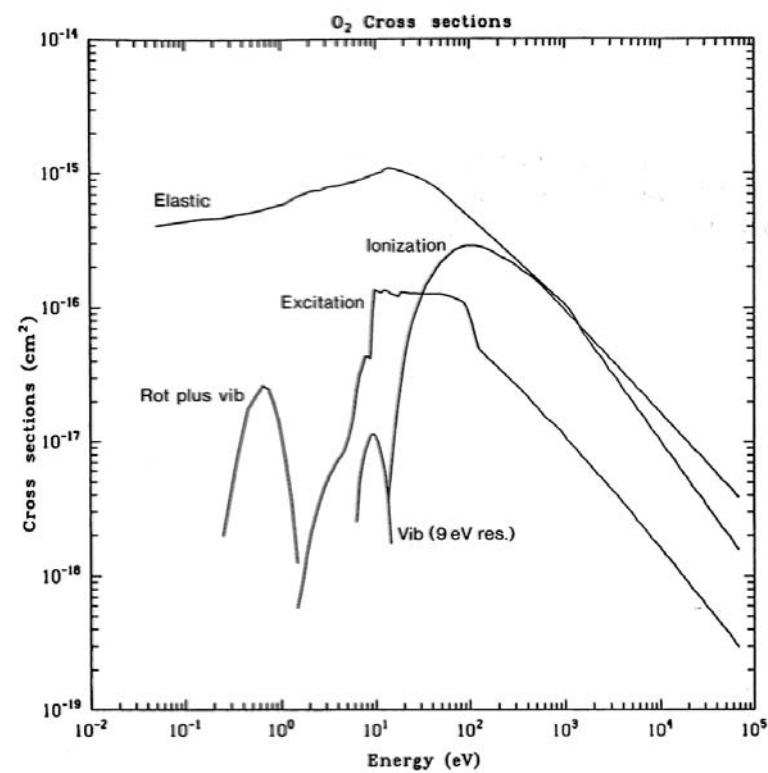


Fig. A4.2 O₂ cross sections.

From Rees, 1989

Ionization, Dissociative Attachment, and Vibrational Excitation Frequencies

Ionization coefficient



Modeled by:

$$\frac{\alpha}{p} = A_i \exp\left(-\frac{B_i}{E/p}\right)$$

Attachment coefficient



Modeled by

$$\frac{\eta}{p} = A_a \left(\frac{E}{p}\right)^{-2} \exp\left(-\frac{B_a}{E/p}\right)$$

Electron mobility

Defined through drift speed

$$v_d = \mu_e E$$

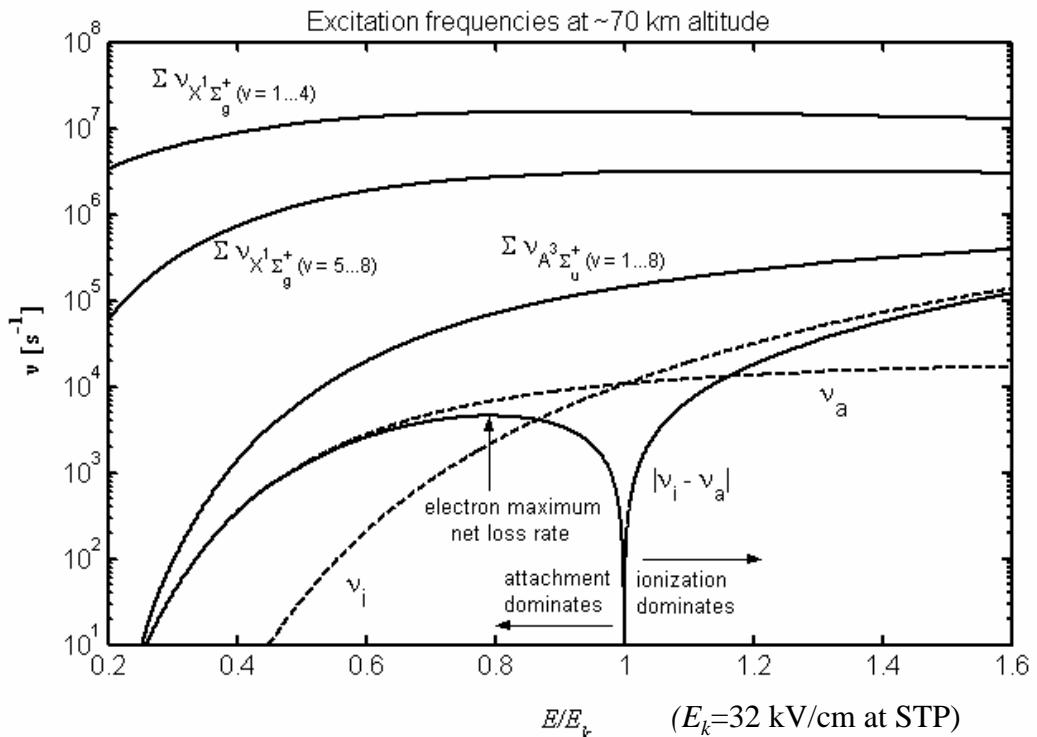
Modeled by

$$\mu_e p = A_m \left(\frac{E}{p}\right)^{-\beta}$$

Ionization and attachment frequencies

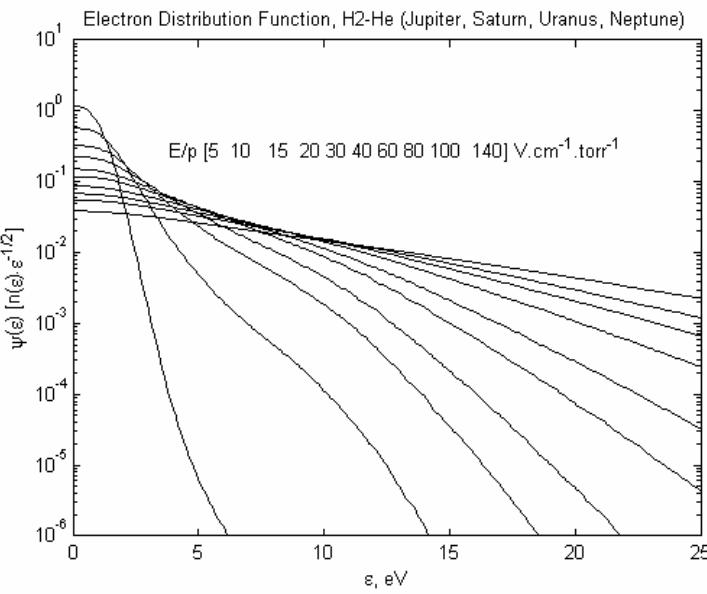
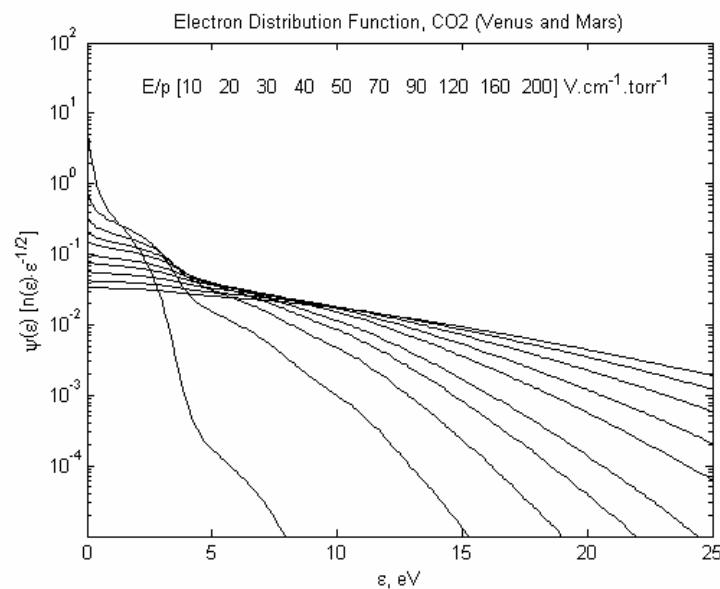
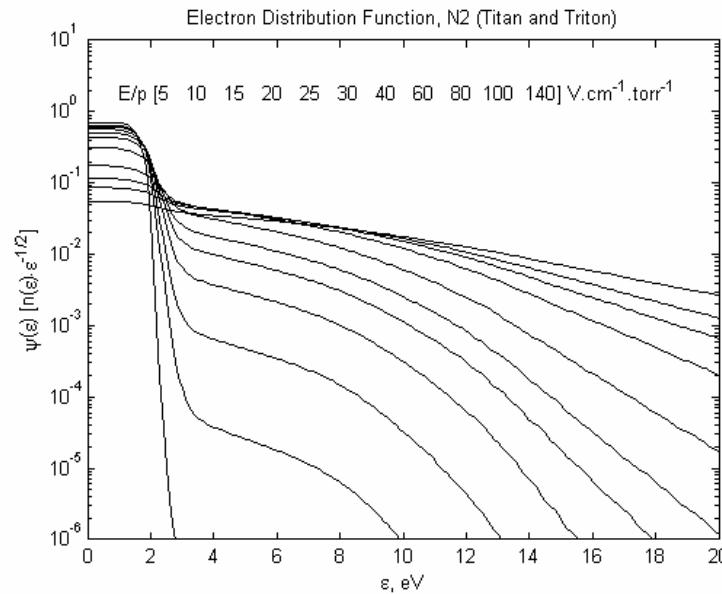
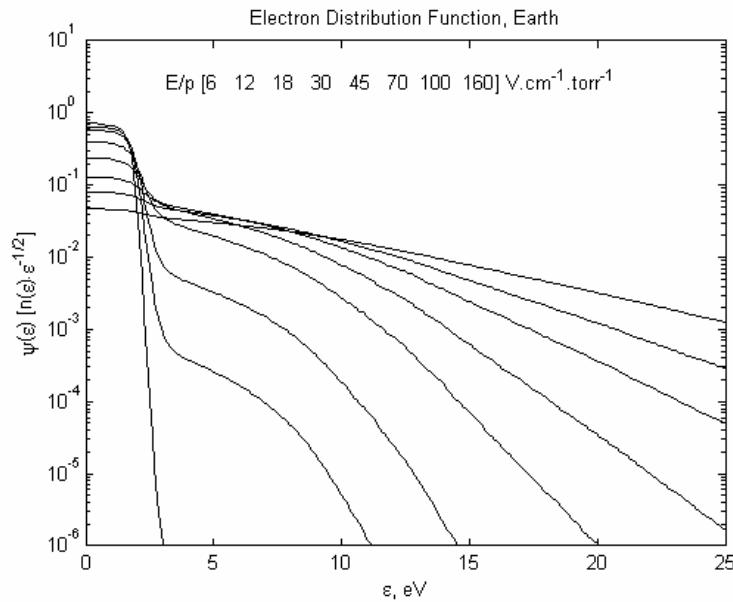
$$\nu_i = \alpha v_d = \alpha (\mu_e p) (E/p)$$

$$\nu_a = \eta v_d = \eta (\mu_e p) (E/p)$$



Vibrational excitations play a significant role in determining the form of the EEDF. In general the excitation frequencies of the vibrational modes of both ground and excited states are much larger than the ionization/attachment frequencies at all undervoltage ($E < E_k$) and modest ($E > E_k$) overvoltage fields.

Electron Distribution Function vs Reduced Electric Field for the Four Major Atmospheric Composition Classes of Planetary Atmospheres



Selected Background References

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