

The Structure of the Solar Atmosphere

Hugh Hudson

UC Berkeley and University of Glasgow

Scope of the Hudson sessions

- 1) Basic principles, language familiarization
- 2) Flares as seen in the lower solar atmosphere
- 3) CMEs and space weather from a flare perspective
- 4) Practicum: EUV spectroscopy with EVE

Relevant Lectures I'm Sorry I Missed!

Week 0	MONDAY 02	TUESDAY 03	WEDNESDAY 04	THURSDAY 05	FRIDAY 06
9:00 - 10:15		Reception	Giovanni Pinzón Fundamentals in Stellar Structure	Giovanni Pinzón Fundamentals in Stellar Structure	Gustavo Guerrero Solar Dynamo
10:15 - 10:30		Break	Break	Break	Break
10:30 - 11:30		Giovanni Pinzón Fundamentals in Stellar Structure	Leonardo Castañeda Emission Process	Benjamín Calvo-Mozo Introduction to radiative transfer	Leonardo Castañeda Emission Process
11:30 - 11:45		Break	Break	Break	Break
11:45 - 12:45		Leonardo Castañeda Emission Process	Cristina Mandrini Solar Activity and Structure: Magnetic Field	Gustavo Guerrero Solar Dynamo	Benjamín Calvo-Mozo Introduction to radiative transfer
12:45 - 14:00		LUNCH	LUNCH	LUNCH	LUNCH
14:00 - 15:15		Cristina Mandrini Solar Activity and Structure: Generalities	Santiago Vargas Introduction to IDL	Cristina Mandrini Solar Activity and Structure: Magnetic Field	Juan Carlos Martínez-Oliveros Sunpy, the future of solar physics
15:15 - 15:30		Break	Break	Break	Break
15:30 - 16:45		Santiago Vargas Introduction to IDL	Dominik Utz Introduction to SolarSoft	Dominik Utz Introduction to SolarSoft	Dominik Utz Introduction to SolarSoft
16:45 - 17:00		Discussion	Discussion	Discussion	Discussion

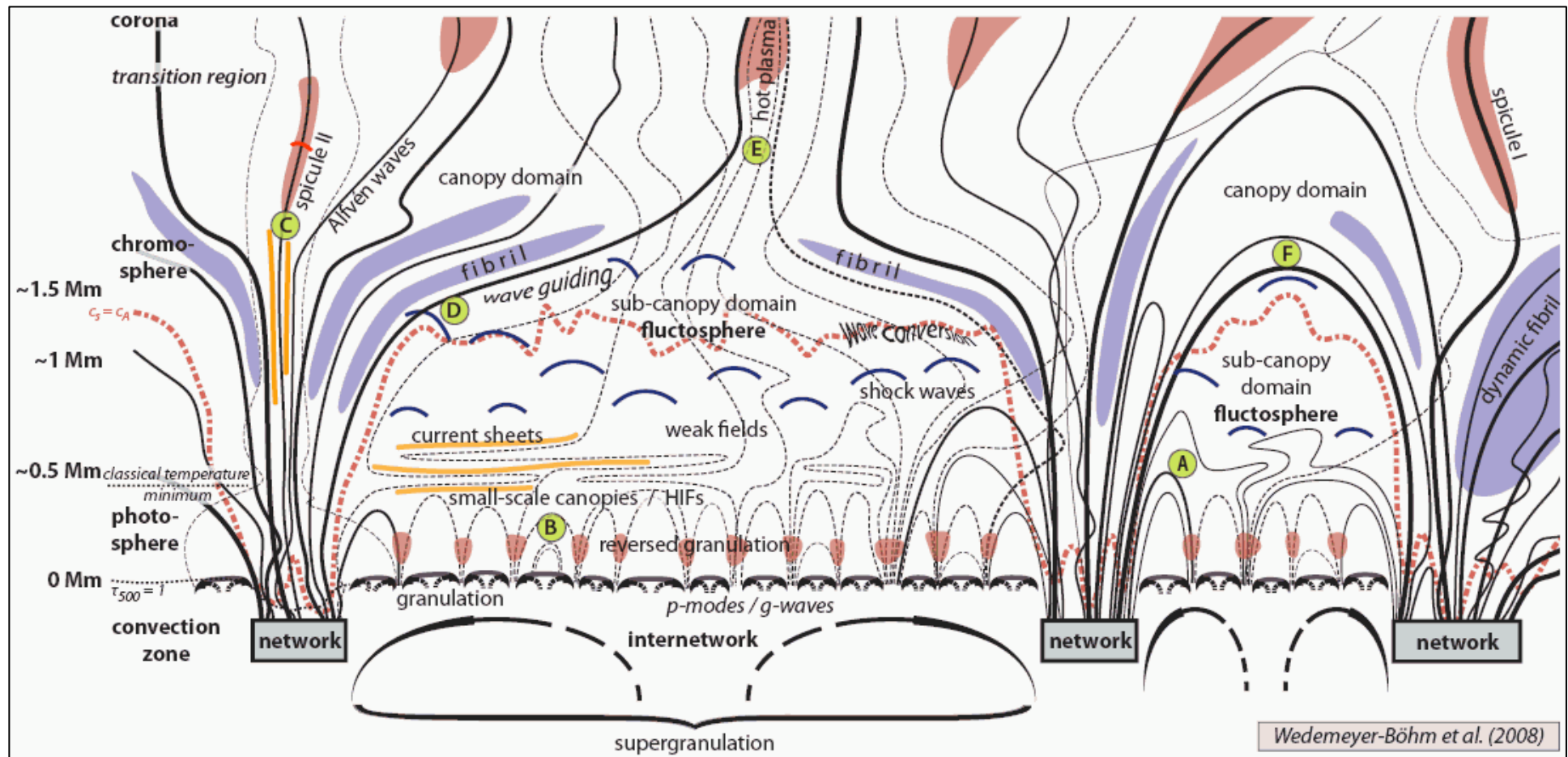
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Outline of presentation

- General structure of solar atmosphere
 - Gravity, radiation, conduction, turbulence
- Model atmospheres
 - Semi-empirical, MHD, exospheric
- Plasma physics
- Simple physics of a flux tube
 - Loop concepts
- Magnetized structures

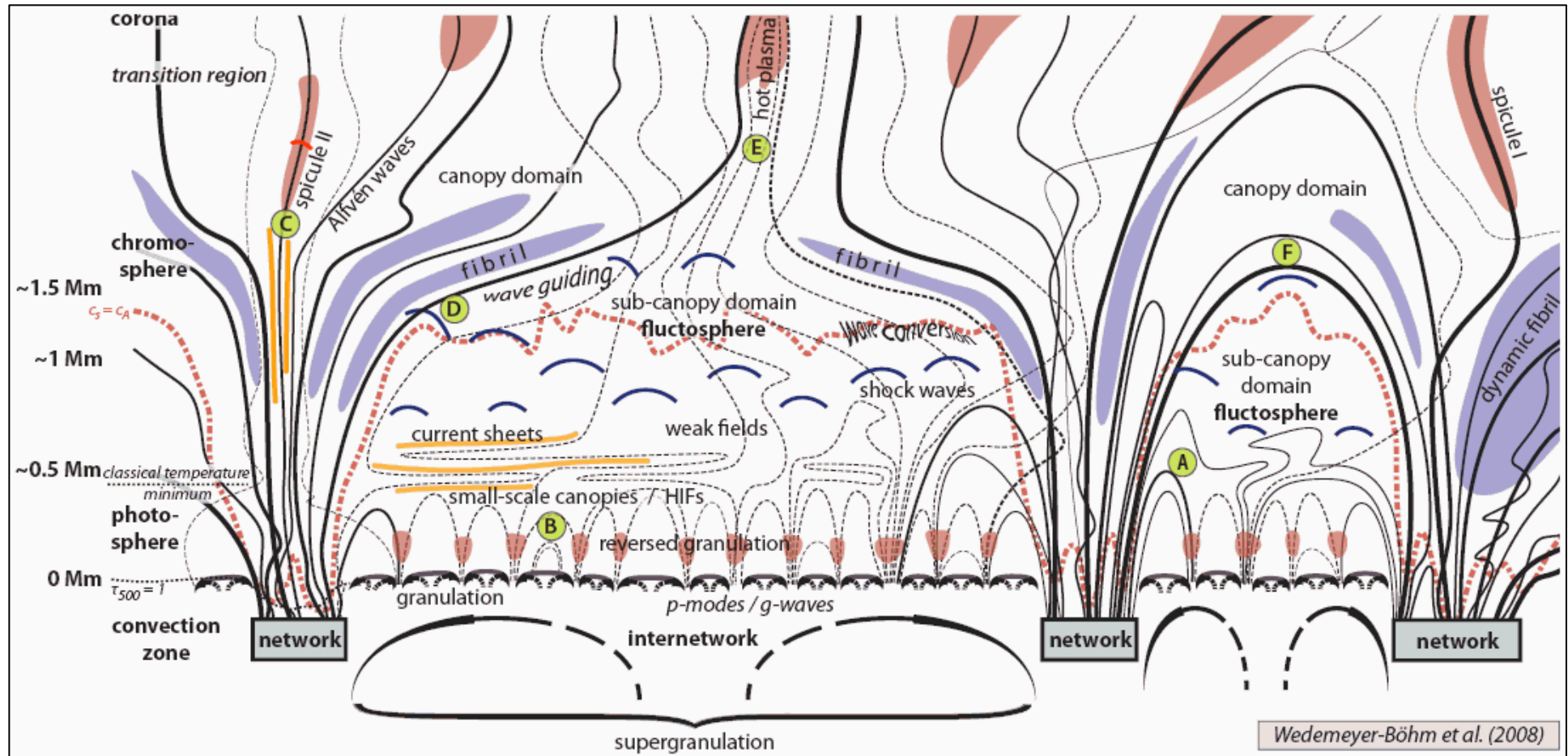
Structure of the atmosphere



Corona
Transition region
Chromosphere
Photosphere

Wedemeyer-Bohm 2004

Structure of the atmosphere



Wedemeyer-Böhm et al. (2008)

Corona

Transition region

Chromosphere

Photosphere

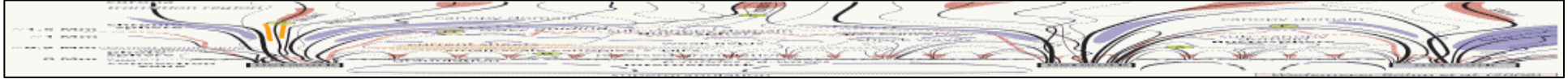


“Interface region”

(=> new IRIS spacecraft)

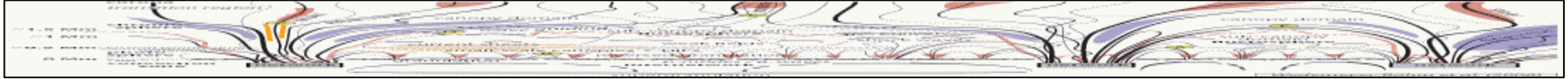
Wedemeyer-Bohm 2004

Structure of the atmosphere



(on a roughly correct aspect ratio)

Gravity is strong



- The effect of gravity can be seen via this truer scale – but still an incorrect one
- Flows and other factors create the complicated structures in Wedemeyer's cartoon
 - Gravity is a powerful agent; $g \sim 274 \text{ m/s}^2$ at the photosphere itself



The “spicule forest” results in vertical scales much larger than the $\sim 2,000 \text{ km}$ of standard models

Radiation is strong

- The photospheric radiation field is intense, and has dynamical consequences because of finite optical depth
- Reasonably competent 1D static models (T, n vs height) can be made with gravity and radiative-transfer theory alone: the “semi-empirical models”
- Such models are forced to fit the spectrum, and thus provide access to basic astrophysical knowledge such as elemental abundances
- The independent parameter in a semi-empirical model is the optical depth τ , normally scaled at the 5000 Å continuum

Thermal conduction is strong

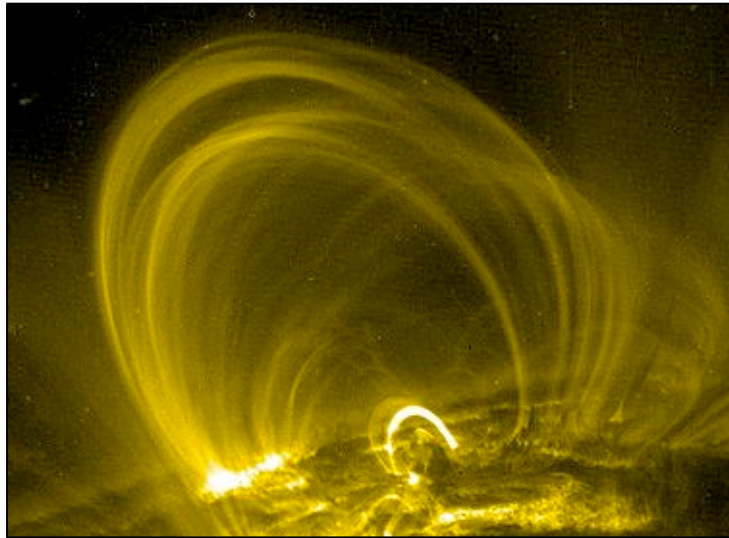
- Classical theory of thermal conductivity in a plasma (“Spitzer conductivity”) gives the energy flux as

$$F_c = -\kappa_0 T^{5/2} dT/ds,$$

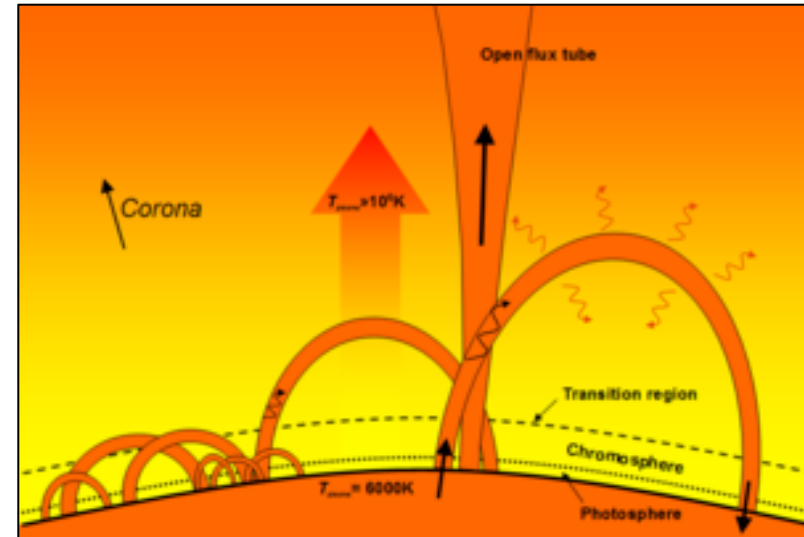
suitable for fluid (MHD) descriptions; more complex transport theory involves **B**

- In the collisionless domain, the Spitzer formula gives way to other concepts because the electron distribution function becomes non-Maxwellian
- In the fluid approximation, this formula gives leads to the “RTV” and other scaling laws

Coronal loops



As observed by TRACE



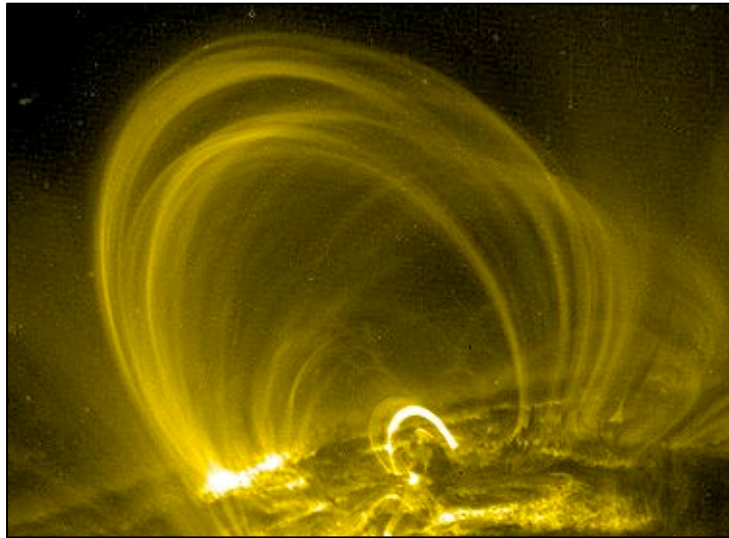
As imagined in a cartoon

- 1) What look like loops are really not stable structures
- 2) Theory tends to ignore the interface region*
- 3) Simple conduction theory allows the derivation of “scaling laws”

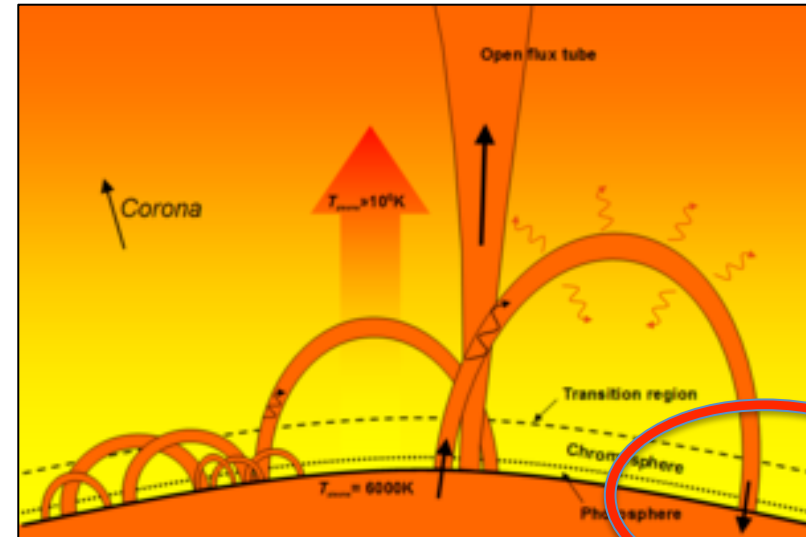
*Recall



Coronal loops



As observed by TRACE



As imagined in a cartoon

Note an
oddy!

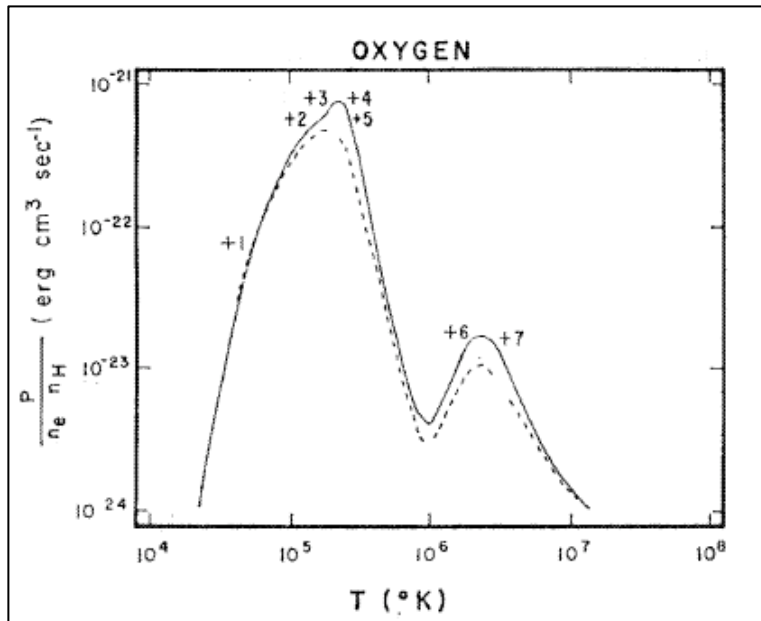
- 1) What look like loops are really not stable structures
- 2) Theory tends to ignore the interface region*
- 3) Simple conduction theory allows the derivation of “scaling laws”

*Recall

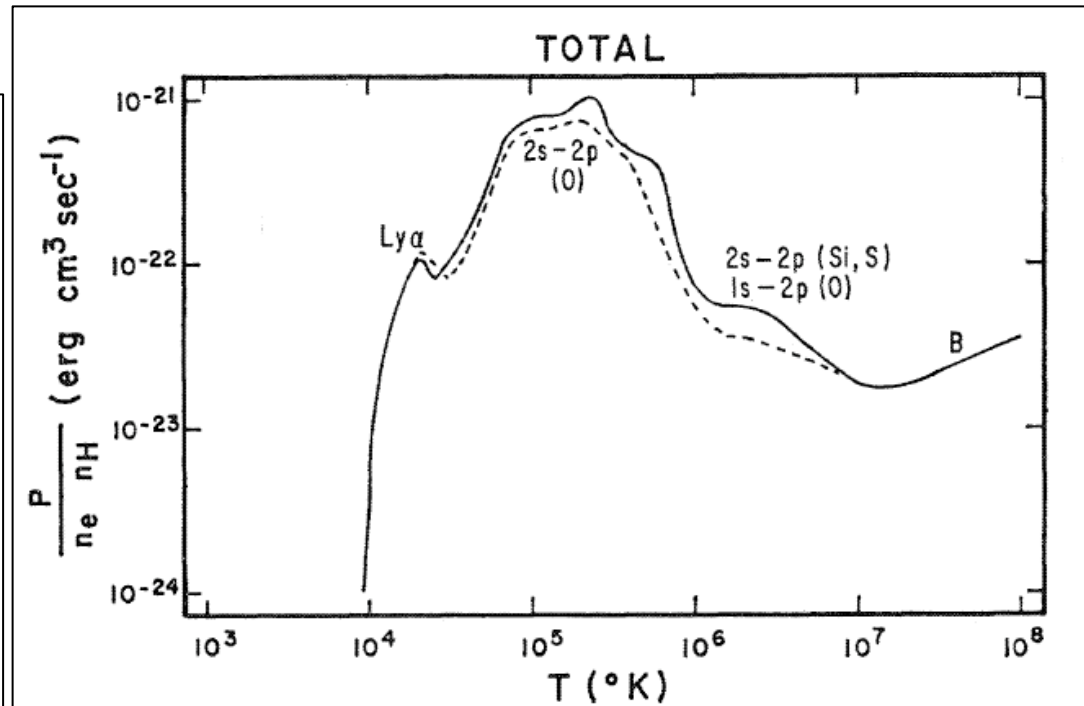


Coronal radiation

(Cox & Tucker, 1969)



Radiation from O ions only

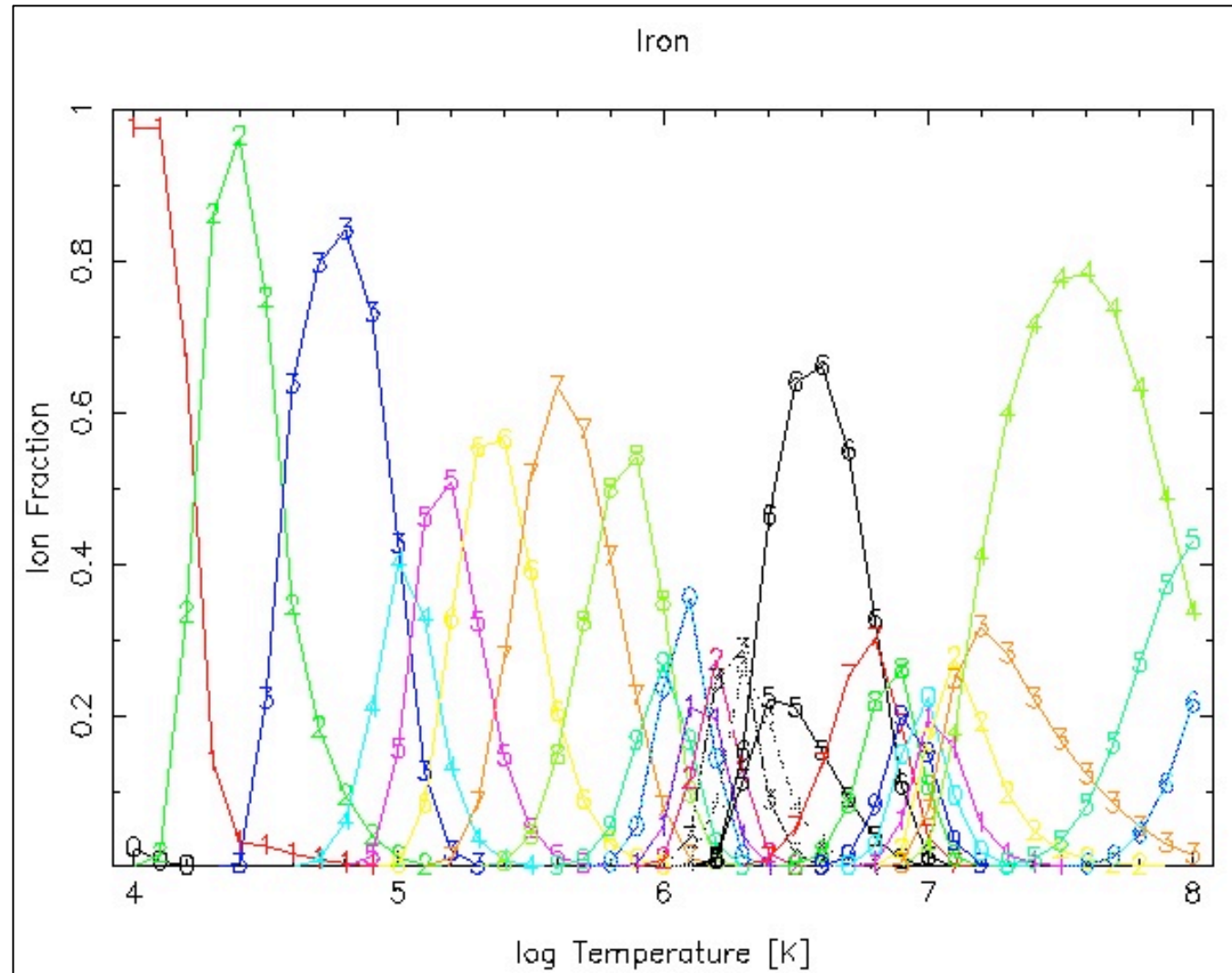


Radiation from a complete plasma:

- Optically thin
- Coronal (not Saha) equilibrium
- Photospheric abundances
- At high temperatures, brems
- Why no Fe?

For fuller and more modern information,
see CHIANTI (via SolarSoft)

Ionization equilibrium



Scaling laws

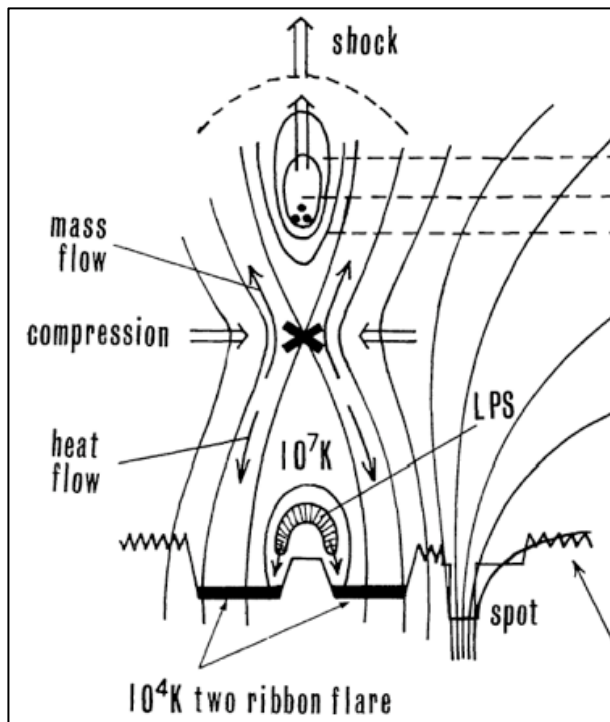
- In conductive equilibrium, “RTV* scaling” gives $T \sim (pL)^{1/3}$, where T and p are temperature and pressure at loop top
- In cooling equilibrium, “Serio scaling” gives $T \sim n^2$ as the loop depressurizes in the aftermath of an injection of energy and mass
- At the base of the solar wind, the pressure depends upon the (unknown) coronal heating function. For models it is simply assumed

* RTV stands for Rosner, Tucker & Vaiana (1978)

The interface region

- This terminology matches the acronym of the IRIS spacecraft, the next solar observatory to be launched
- The interface region separates two highly different physical domains. Generally
 - the **sub-photosphere** is a cold, (single-temperature) fluid, optically thick, collisionally thick, and with high plasma beta
 - the **corona** is hot, optically thin, not very collisional, and with low plasma beta
- We therefore expect that much of the physics that relates the corona to the interior (eg., heating) originates in this region

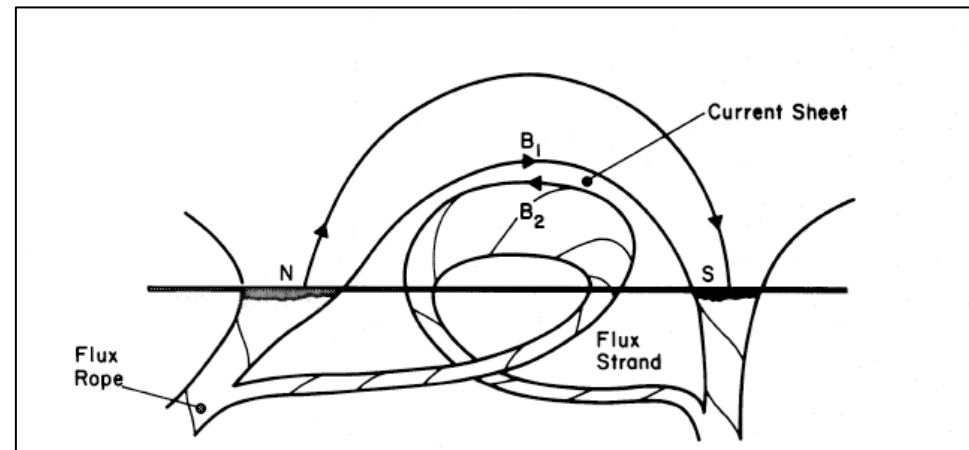
What is the behavior of a coronal loop?



Hydrodynamic: Hirayama, 1974

- Note the depression of the chromosphere, the result of coronal overpressure

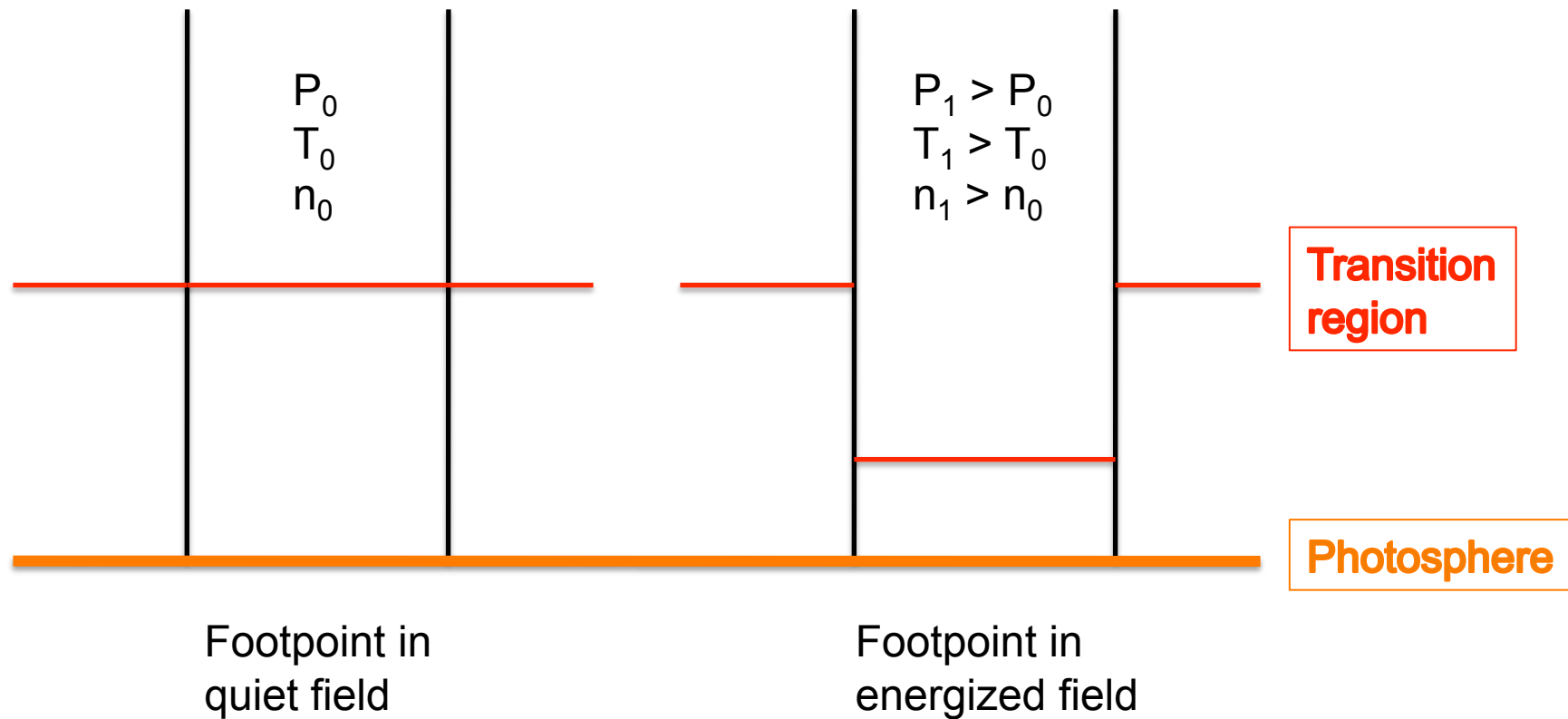
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$



Electrodynamic: Piddington, 1961

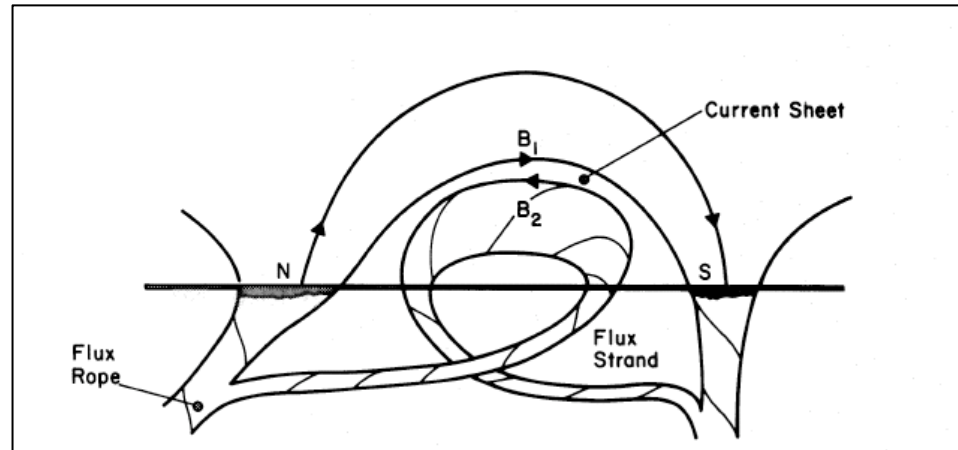
- Note the sub-photospheric twists, which inject currents into the atmosphere according to Ampere's Law

Loop energization



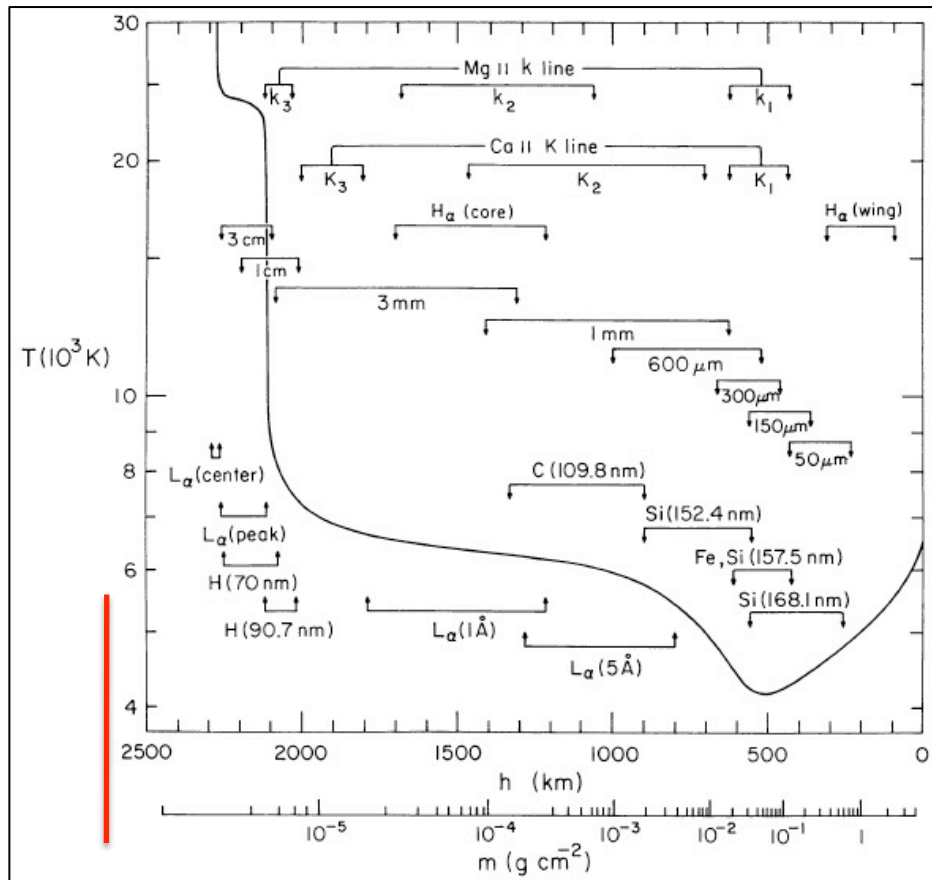
An increase of coronal plasma temperature forces new material up from the chromosphere, increasing both pressure and density

Coronal currents and dynamic “loops”



- *Parallel currents in the corona (ie, $\mathbf{J} \parallel \mathbf{B}$) result from twists present in sub-photospheric fibril fields that emerge through the surface.*
- *At the opposite ends of a loop, a conflicting condition on the injected current must generally prevail.*
- *No static MHD solution of this problem seems possible*

A semi-empirical model



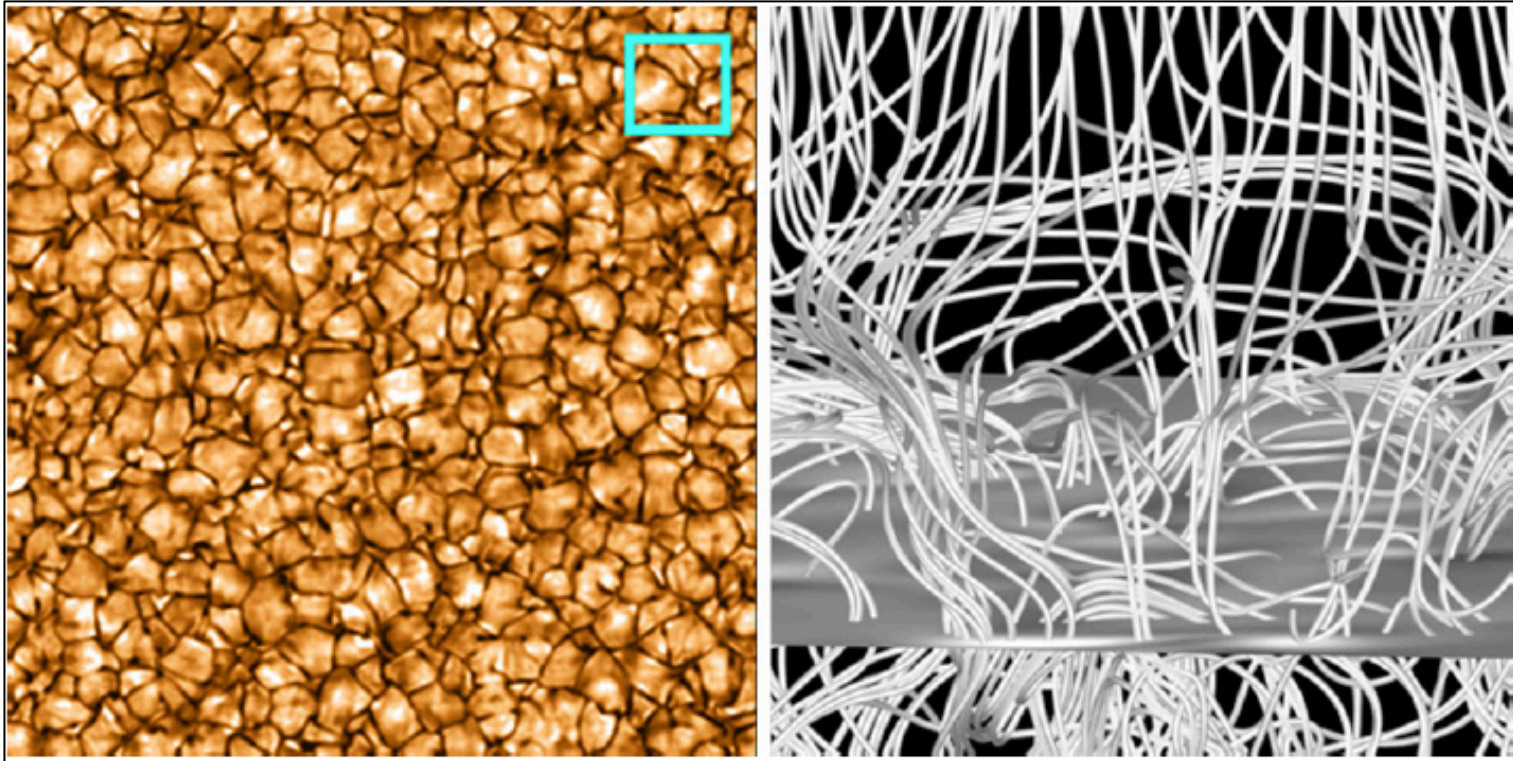
Vernazza et al., 1981

- This is a standard example. Note the height and column-density scales: really the opacity τ is the independent variable here
- Note also the array of spectral features, including the IR and mm-wave continuum (cf. ALMA)
- The photosphere ($t_{5000\text{\AA}} = 1$) has a column density of $2\text{-}3 \text{ g/cm}^2$ (about one Thomson length) and a density of $\sim 10^{13} \text{ kg/m}^3$
- The corona seems too low! This is because of dynamical effects

What is really happening in the interface region?

- Surface convection launches wave energy into it
 - Shock waves will form because of dispersion
 - Differential magnetization exists
 - Non-convective electric fields will occur locally
 - Ionization will depart from equilibrium locally
- In short, kinetic physics will ultimately be needed. But, in the meanwhile, we have single-fluid MHD approaches

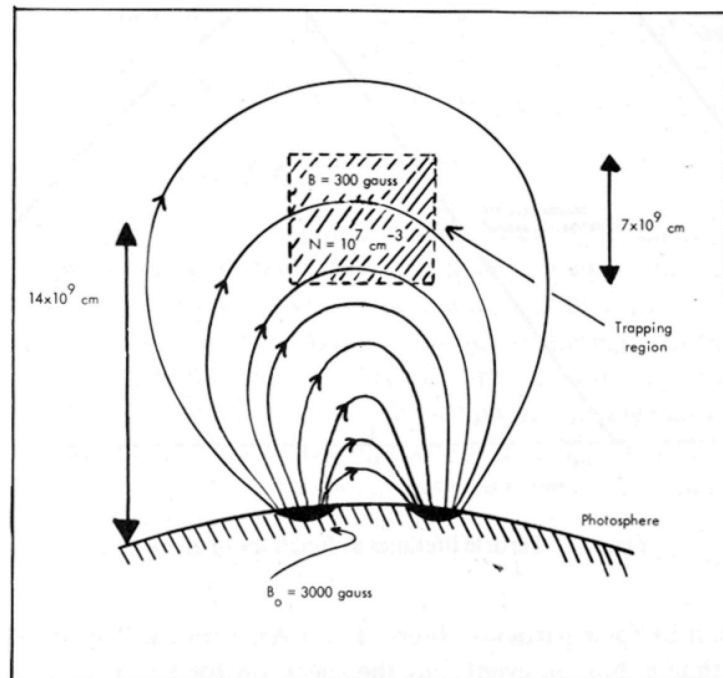
Interior-to-corona modeling of the interface region



Perspective view (right) of the magnetic field in a small domain of an interior-to-corona MHD simulation (Abbett & Fisher, 2012). Note the magnetic complexity in the interface region

Exospheric models?

- Basic plasma physics starts with particles, treated at the microscopic level kinetically
- The distribution function of these particles, in general, is not Maxwellian (cf. Scudder theory of “velocity filtration”)
- If the plasma has low collisionality, these complications must be dealt with (e.g., Størmer theory and adiabatic-motion theory)



- Størmer* theory in the coronal context (Elliot, 1973): coronal non-thermal energy storage without twist
- Do adiabatic drift motions make this idea impractical?

* *Next slide*

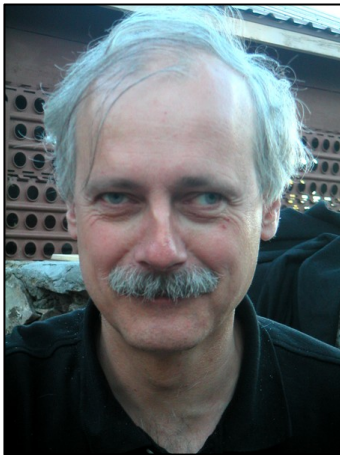
Carl Størmer



With Birkeland, observing the aurora (?) in 1910

Better models, still MHD

- Three-dimensionality
- Time dependence (hydrodynamics)
- Magnetic field (MHD)
- Efficient radiative transfer algorithms



Carlsson



Nordlund

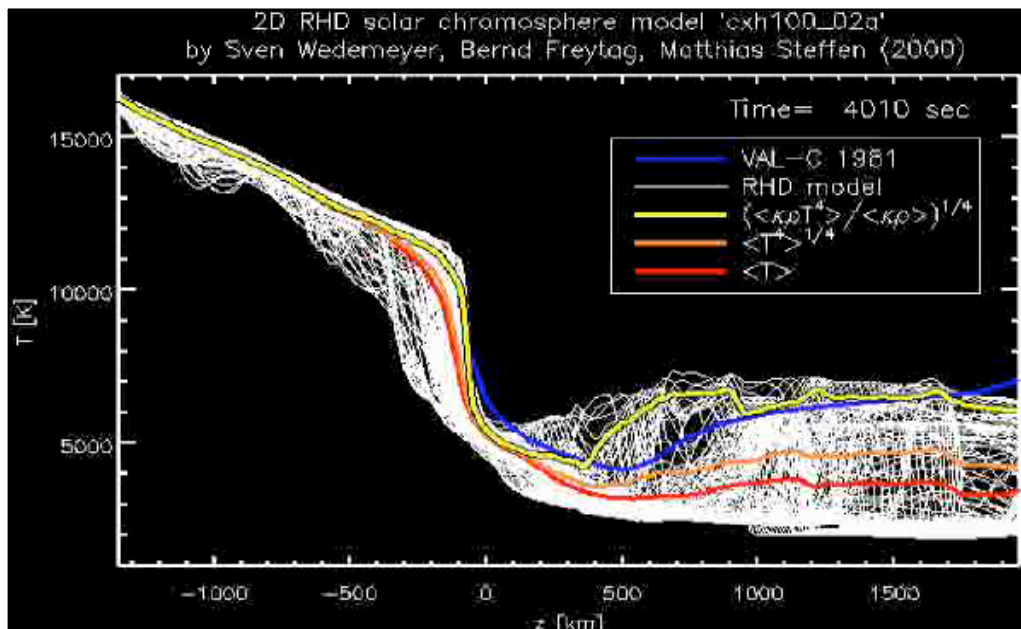
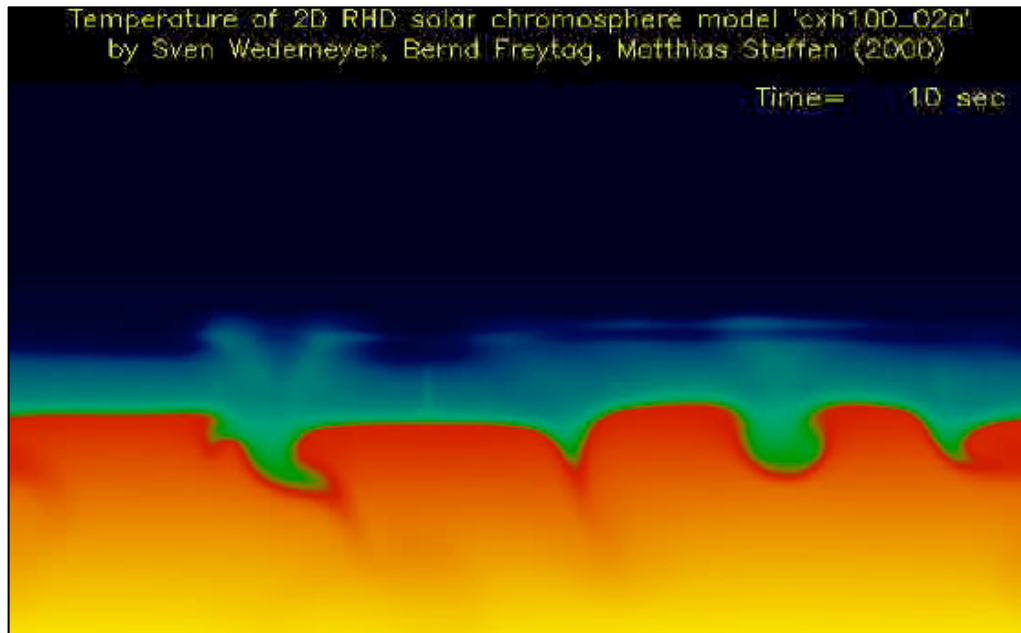


Hansteen



Rutten

Better models



Representative simulations

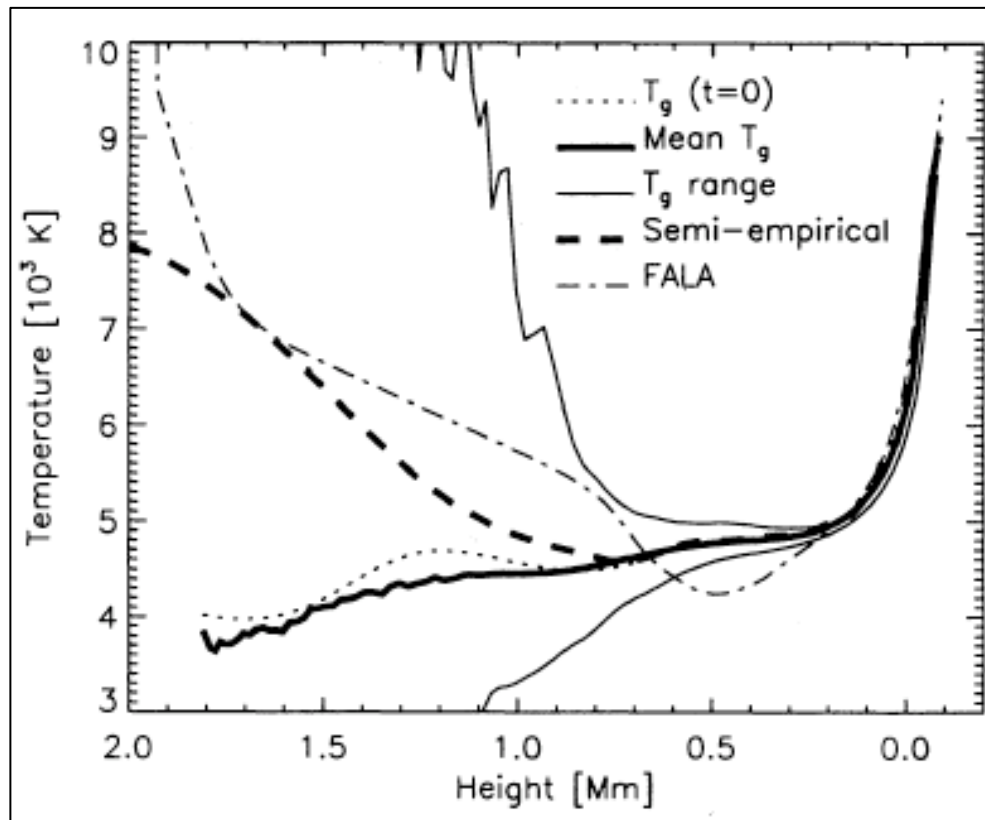
- Shock waves form
- Large temperature excursions result
- The VAL-C model (blue line) misses in important details both in space and time

- Even these “better” models omit 99% of the plasma physics, since they rely upon one-fluid MHD theory

Courtesy Sven Wedemeyer-Bohm

A basic problem

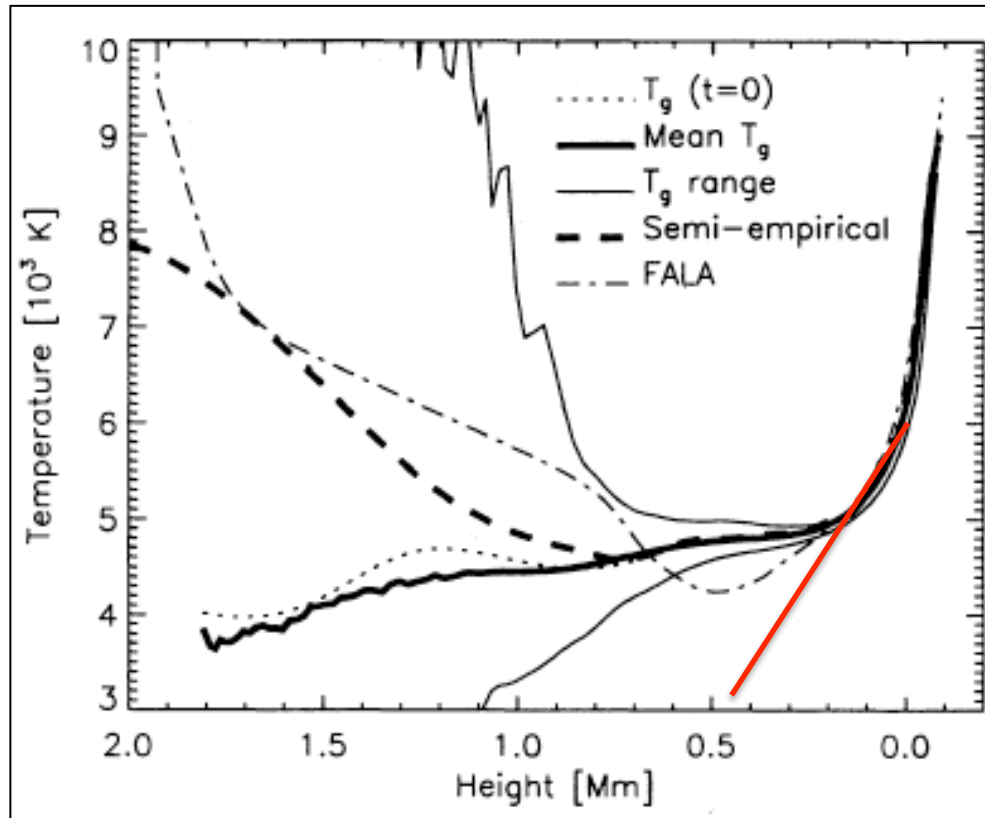
Why is the chromosphere so hot?



Carlsson & Stein, 1995

A basic problem

Why is the interface region so hot?



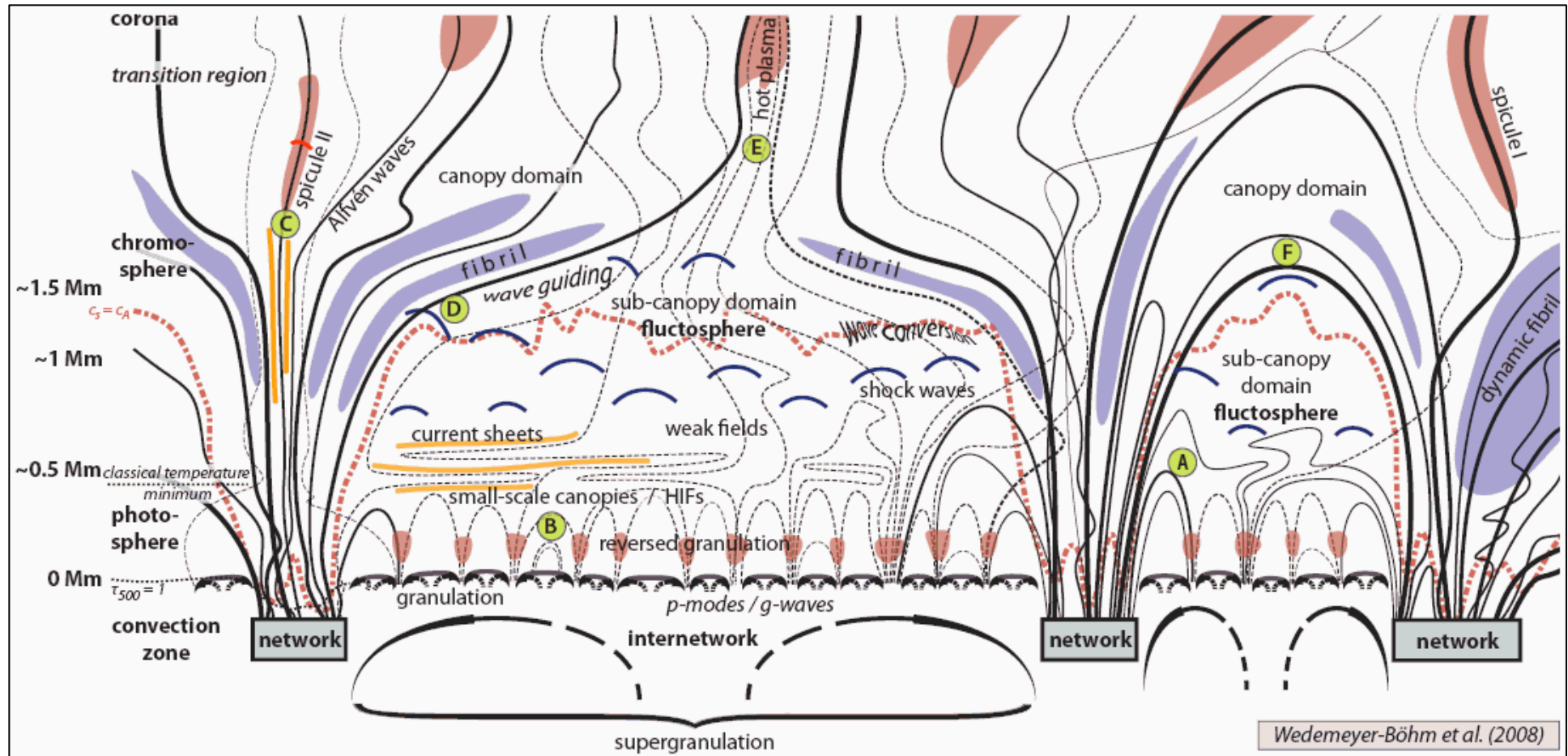
- These simple models immediately miss the temperature minimum!
- In fact, the temperature minimum should asymptotically approach the microwave background, 3 K:

$$T_{eq}(r) = \frac{(1 - \Omega_{\odot}) \times 3 + \Omega_{\odot} \times T_{\odot}}{4\pi} \text{ K}$$

$$r = R_{\odot} \rightarrow \Omega_{\odot} = 2\pi \rightarrow T_{eq} \approx T_{\odot}/2$$

Carlsson & Stein, 1995

Structure of the atmosphere



Photosphere

Chromosphere

Transition region

Corona

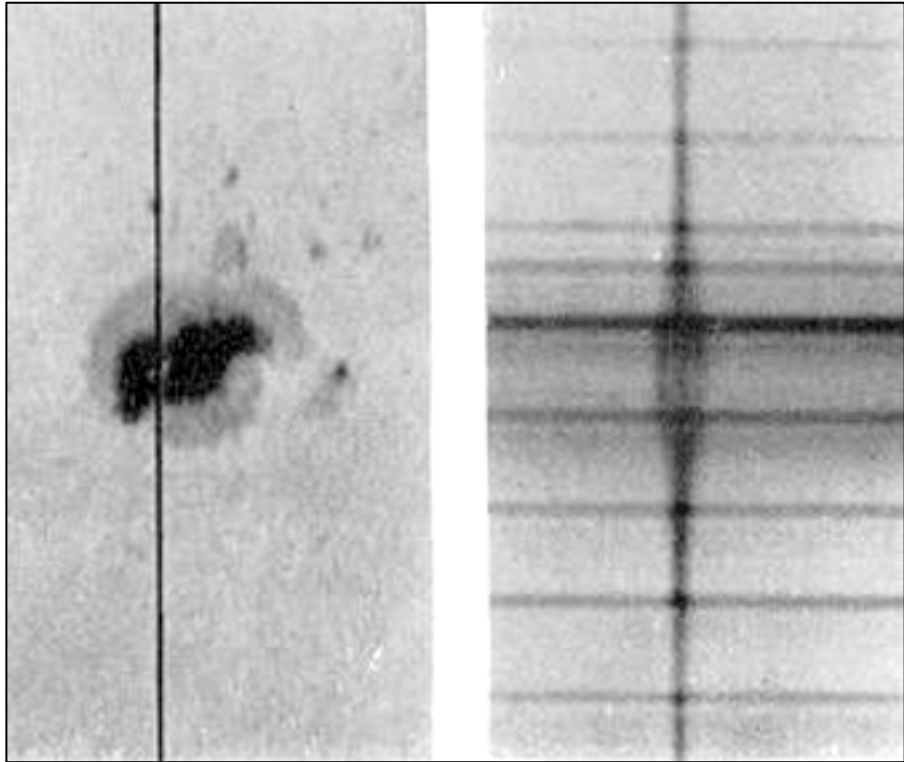


“Interface Region”

(=> new IRIS spacecraft)

Wedemeyer-Bohm 2004

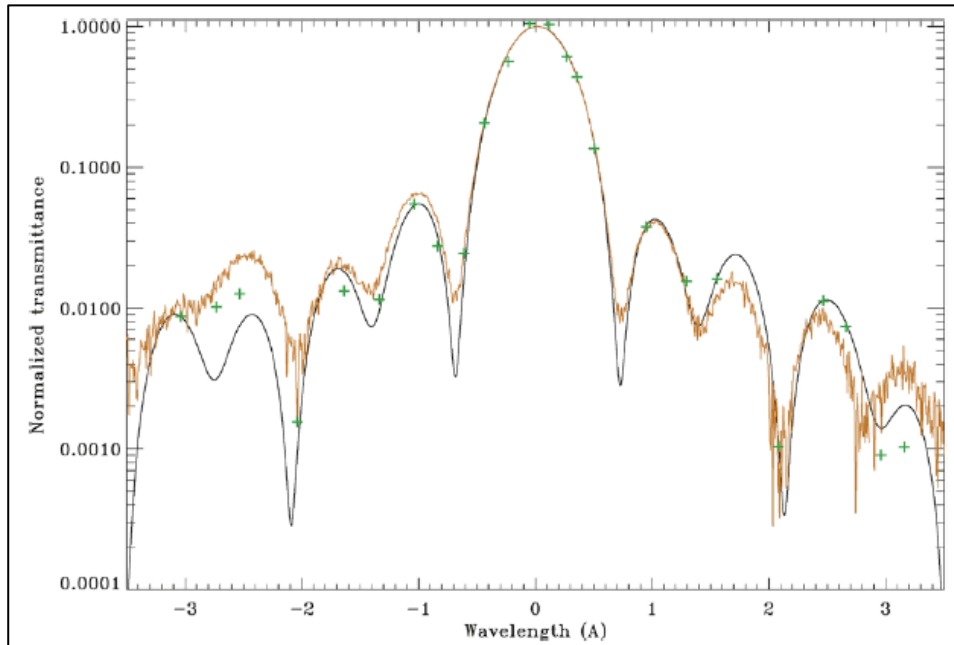
Coronal magnetism 1: Basic physics



Hale et al. 1919

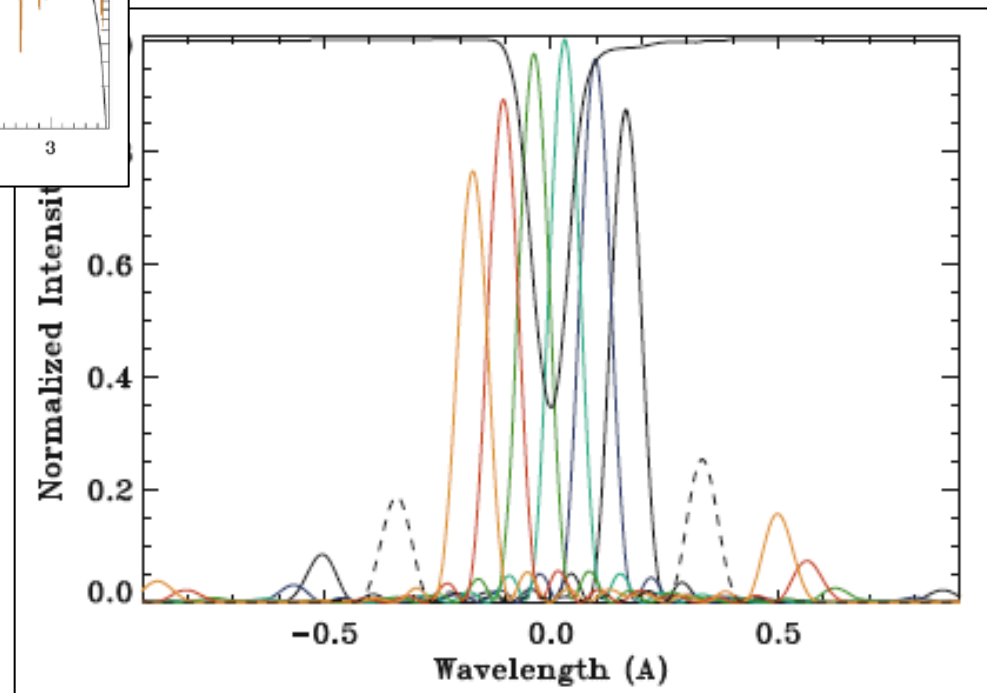
- Zeeman splitting allows for the remote sensing of magnetic field
- The full Stokes parameters allow one to determine the line-of-sight and transverse components
- This is possible in the photosphere, and maybe in the chromosphere and corona, but with caveats

Coronal magnetism 2: HMI spectral response

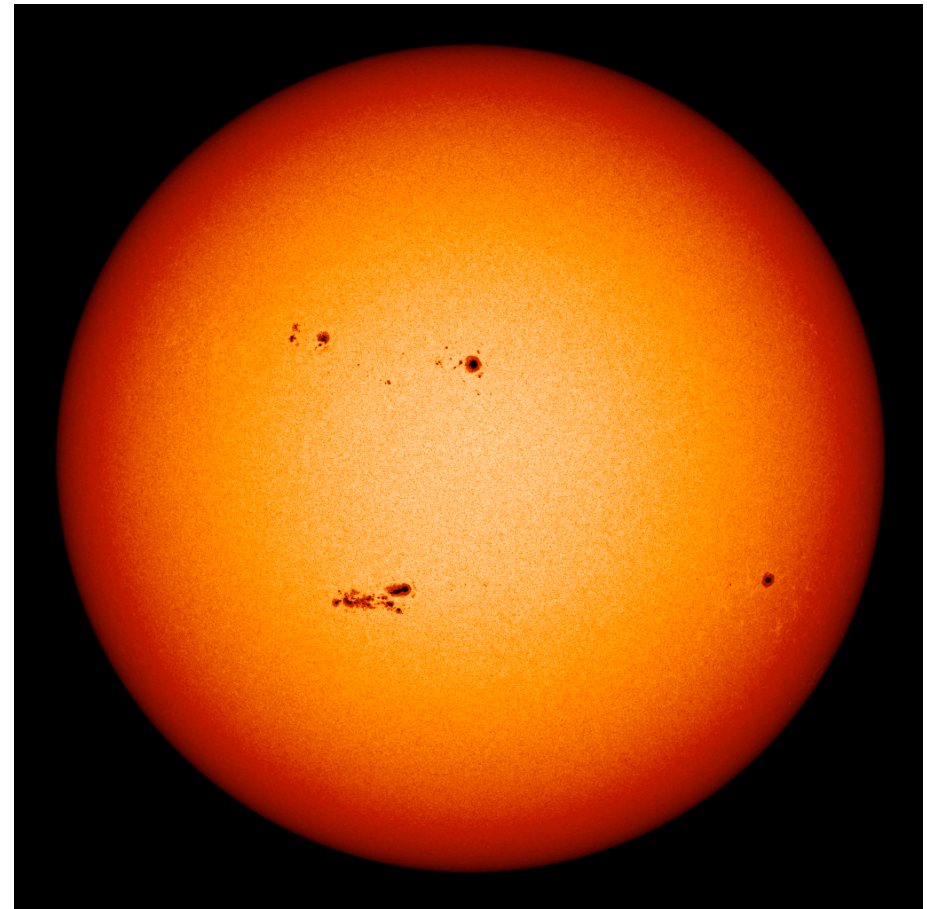
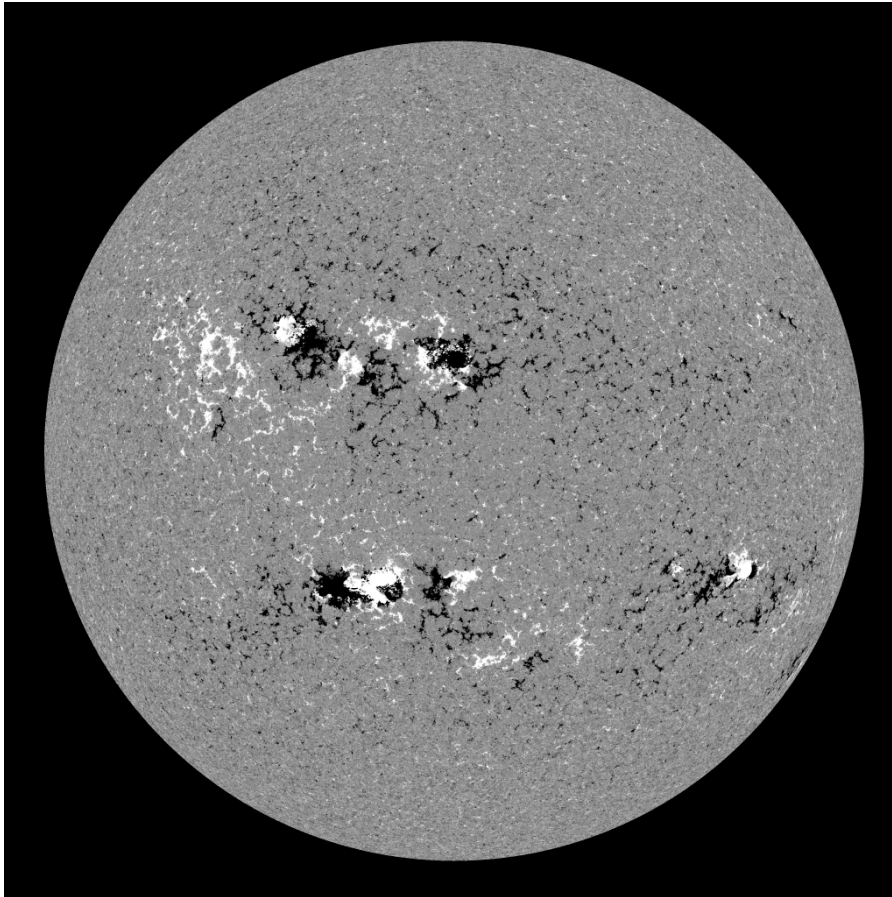


Fixed filter

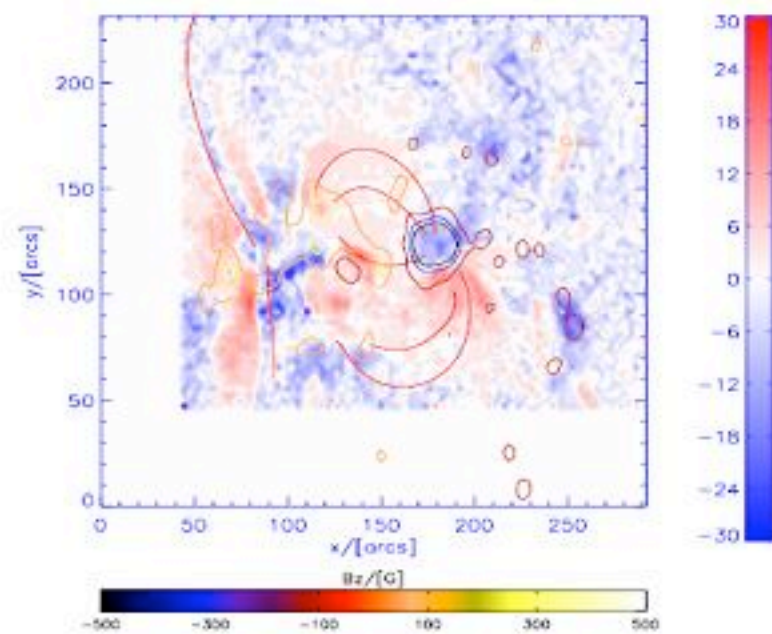
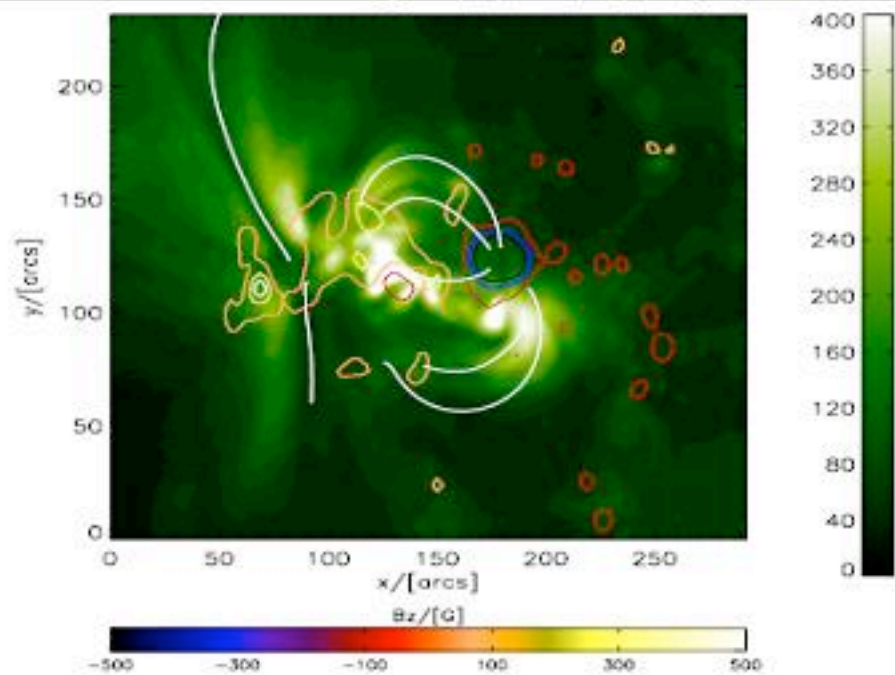
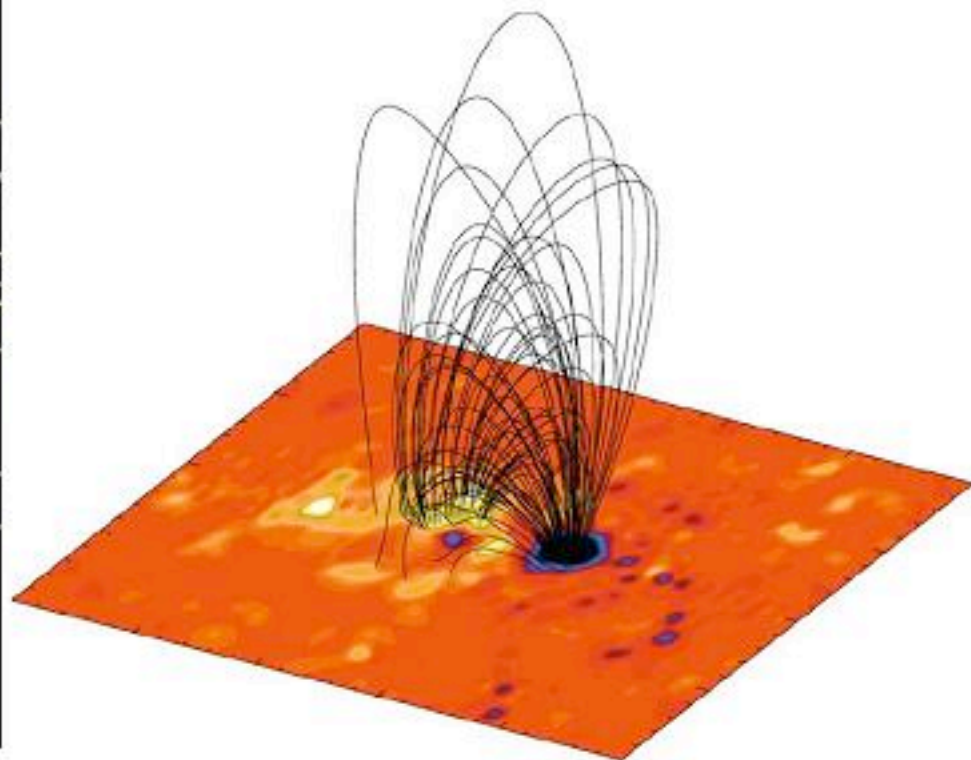
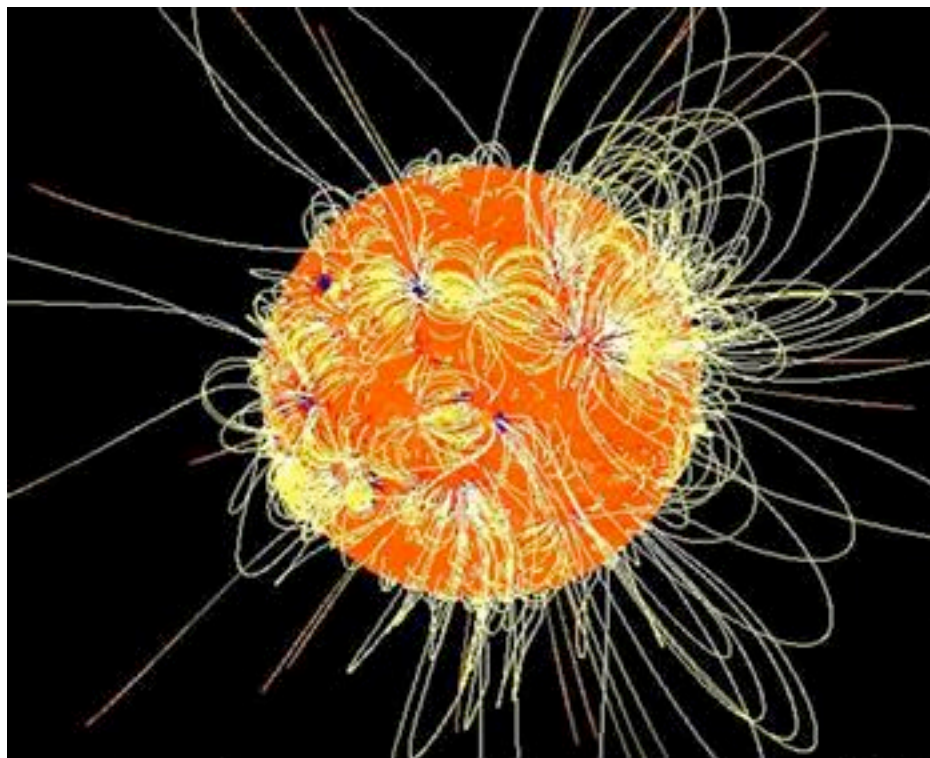
Tunable filter and Fe I
6173.3433 Å line



Coronal magnetism 3: HMI magnetograms



Line-of-sight field and intensity, 2 July 2012



Recall energy storage as currents

MHD version of Ampère's Law

$$\mathbf{j} = \nabla \times \mathbf{B} / \mu$$

Twisting the field produces “free energy” in the form of its inductive field.

Assume ~ steady state, with negligible gravitational forces and pressure gradients (low-beta corona). Then

$$\mathbf{J} \times \mathbf{B} = 0 \quad \text{Force-free condition, i.e., the field and current are aligned}$$

This means that $\mathbf{curl} \mathbf{B} = \alpha \times \mathbf{B}$ with the “force-free parameter” α being constant in the field direction

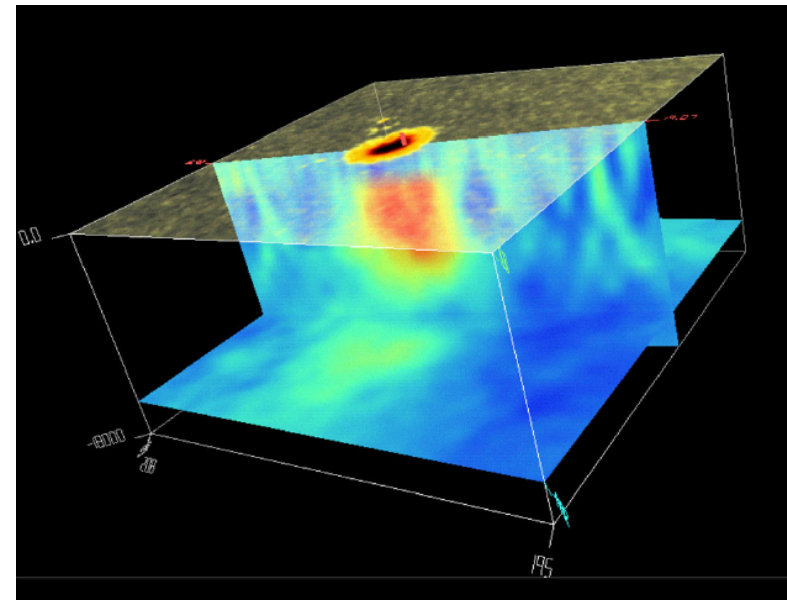
Coronal magnetism 5: Extrapolation techniques

$$\nabla \times \mathbf{B} = \alpha \mathbf{B}$$

- *Potential Field Source Surface (PFSS)* assumes $\alpha = 0$ (no electric currents)
- *Linear Force Free* assumes $\alpha = \text{constant}$
- *Nonlinear Force Free* assumes $\alpha = \alpha(x, y)$
- *Geometrical* relies upon striations viewed in the image as a quantitative guide (Malanushenko et al. 2009)

Solar active regions

- Sunspots, faculae, plage
- *Three-dimensionality* (Wilson depression, hot-wall effect, “loops”)
- p-mode absorption



Scherrer et al, 2012

Conclusions

- The nature of the solar atmosphere, and its physics, remain ill-understood
- The “interface region” is particularly important, even though often ignored in coronal theories
- We have many new tools now that will let us explore the behavior of the interface region more fully – ie, its future is bright
- The MHD approximation must be viewed with great caution

Stephanie's plasma pages

<http://sprg.ssl.berkeley.edu/~hudson/plasma/webpage/plasma.html>

What do semi-empirical models, such as those of Fontenla et al (2009) suggest for plasma conditions? Note that these are (non-magnetized) *fluid* models, and so there is limited scope.

Model number	Feature modeled	P @ $2 \times 10^5 \text{K}$	B
		dyne cm^{-2}	Gauss
Model 1001:	Quiet Sun inter-network	0.235	100
Model 1005:	Facula (i.e. very bright plage)	1.62	300
Model 1006:	Sunspot umbra	3.86	1000

Stephanie's plasma pages

<http://sprg.ssl.berkeley.edu/~hudson/plasma/webpage/plasma.html>

[Column density](#)

[Thermal velocities](#)

[Alfvén speed](#)

[Plasma beta](#)

[Debye parameter](#)

[Gyrofrequencies](#)

[Plasma Frequencies](#)

[Electron collision frequency](#)

[Proton collision frequency](#)

[Debye length](#)

[Electron inertial length](#)

[Ion inertial length](#)

[Electron skin depth](#)

[Ion skin depth](#)

[Electron mean free path](#)

[Proton mean free path](#)

[Hall conductivities](#)

[Pedersen conductivities](#)

[Cowling conductivity](#)

[Tensor electrical conductivities](#)

[Magnetization](#)

Resources

- M. Stix, 1989 “The Sun, an Introduction” (basic material on quiet Sun)
- D. Billings, 1966 “A Guide to the Solar Corona” (background on solar corona)
- A. Hundhausen, 1972, “Coronal Expansion and the Solar Wind” (basic theory of solar wind)
- A.G. Emslie et al. 2012, “High-Energy Aspects of Solar Flares,” (overview of flares): SSR vol. 159
- F. Chen, 1984, “Introduction to Plasma Physics and Controlled Fusion” (plasma physics text)
- H. Hudson, 2011, “Global Properties of Solar Flares,” SSR 158, 5
- Web resources
 - Living Reviews <http://solarphysics.livingreviews.org/>
 - Nugget collections
http://sprg.ssl.berkeley.edu/~tohban/wiki/index.php/RHESSI_Science_Nuggets et al.
 - Stephanie’s plasma pages
<http://sprg.ssl.berkeley.edu/~hhudson/plasma/webpage/plasma.html>