



University | School of Physics  
of Glasgow & Astronomy

# The Sun's Atmosphere

Session 2024-24

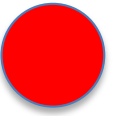
Lecture 9: "Observing and interpreting solar magnetism"

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Monday, 12 February 2024

# The Sun's Atmosphere

| Date       |    | Topic                                | Lecturer         |
|------------|----|--------------------------------------|------------------|
| Mon 08-Jan | 0  | Intro                                | Iain Hannah      |
| Tue 09-Jan | 1  | Overview Structure & Dynamics        | Iain Hannah      |
| Mon 15-Jan | 2  | Particle Acceleration & Transport 1  | Eduard Kontar    |
| Tue 16-Jan | 3  | Particle Acceleration & Transport 2  | Eduard Kontar    |
| Mon 22-Jan | 4  | MHD Basics 1                         | Emma Hunter      |
| Tue 23-Jan | 5  | MHD Basics 2                         | Emma Hunter      |
| Mon 29-Jan | 6  | Radiation Transport 1                | Nicolas Labrosse |
| Tue 30-Jan | 7  | Radiation Transport 2                | Nicolas Labrosse |
| Mon 05-Feb | 8  | Photosphere and Magnetism 1          | Hugh Hudson      |
| Mon 12-Feb | 9  | Photosphere and Magnetism 2          | Hugh Hudson      |
| Mon 19-Feb | 10 | EUV to Infrared Plasma Diagnostics 1 | Sargam Mulay     |
| Mon 26-Feb | 11 | EUV to Infrared Plasma Diagnostics 2 | Sargam Mulay     |
| Tue 27-Feb | 12 | Radio Plasma Diagnostics             | Yingjie Luo      |
| Mon 04-Mar | 13 | Flares & CMEs                        | Iain Hannah      |
| Tue 05-Mar | 14 | X-ray/Gamma-ray Plasma Diagnostics   | Iain Hannah      |
| Mon 11-Mar | 15 | Space Weather                        | Iain Hannah      |
| Tue 06-Feb | T1 | Tutorial 1 - Oral Exam 1 Prep        | Iain Hannah      |
| Tue 20-Feb | T2 | Tutorial 2 - Project Intro           | Iain Hannah      |
| Tue 12-Mar | T3 | Tutorial 3 - Oral Exam 2 Prep        | Iain Hannah      |
| Mon 18-Mar | T4 | Tutorial 4 - Project Prep            | Iain Hannah      |



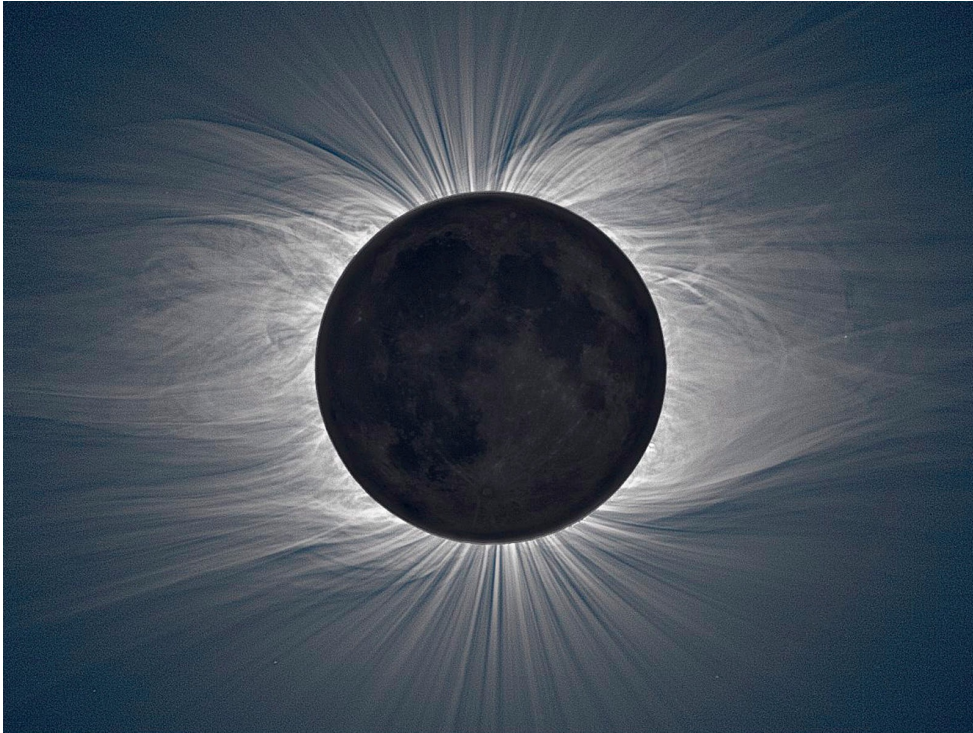
# Useful homework questions based on this material

- What are the forces acting on plasma in a coronal magnetic loop, rough quantitative estimate?
- For what value of  $B$  can a flux tube escape from the tachocline and emerge in one cycle?
- What is the electrical potential at the photosphere?
- At what stereoscopic angle can Stokes (I,V) do better than Stokes (I,Q,U,V)?

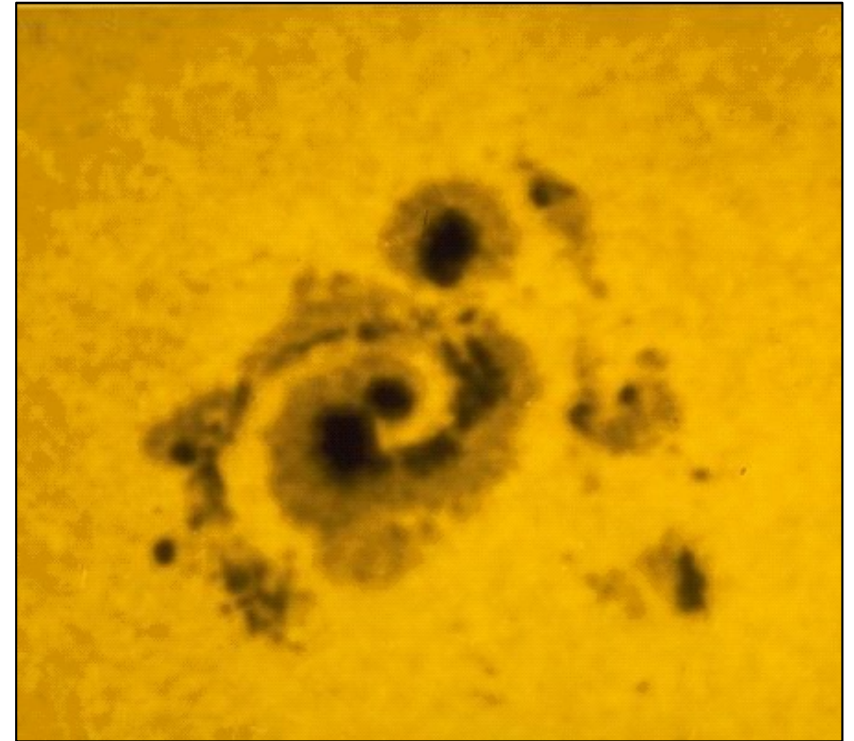
*These are questions to be answered by rough approximation, and the necessary facts should be here on the slides. I will post discussion of these items at*

*<http://www.ssl.berkeley.edu/~hudson/presentations/supa.220207/>*

# Solar magnetism



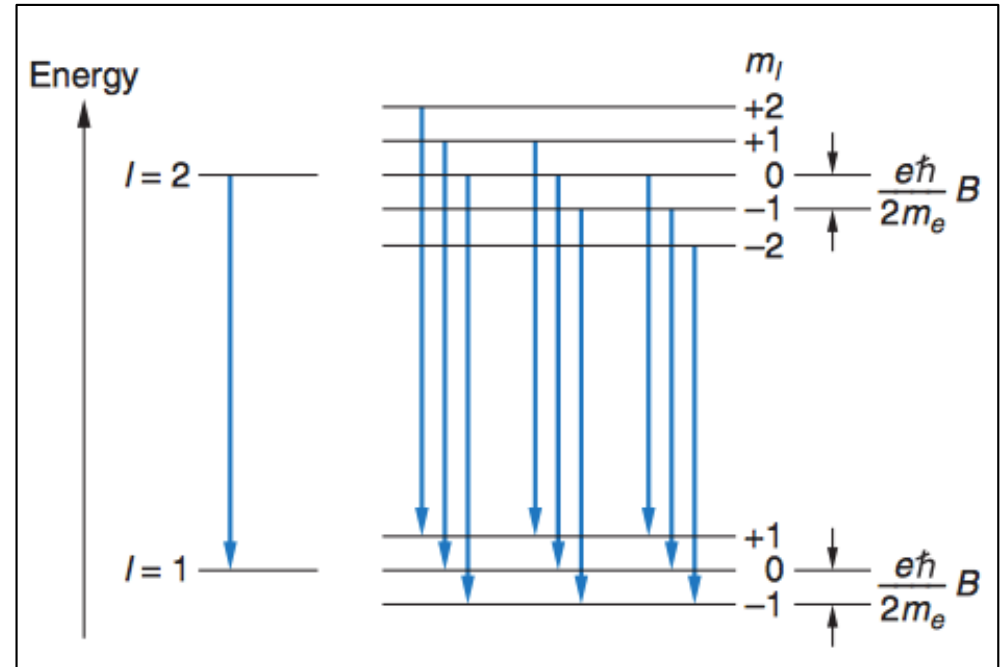
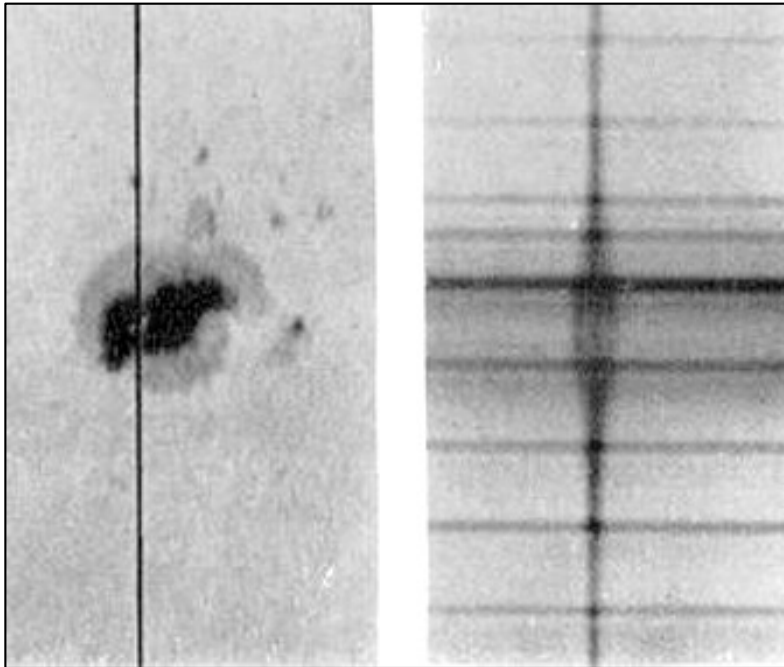
- The fine striations at the poles during an eclipse sure look like a dipole magnetic field.



- Sometimes sunspots have spiral patterns.

• Clues such as these led Hale to confirm the existence of solar magnetism via the Zeeman effect.

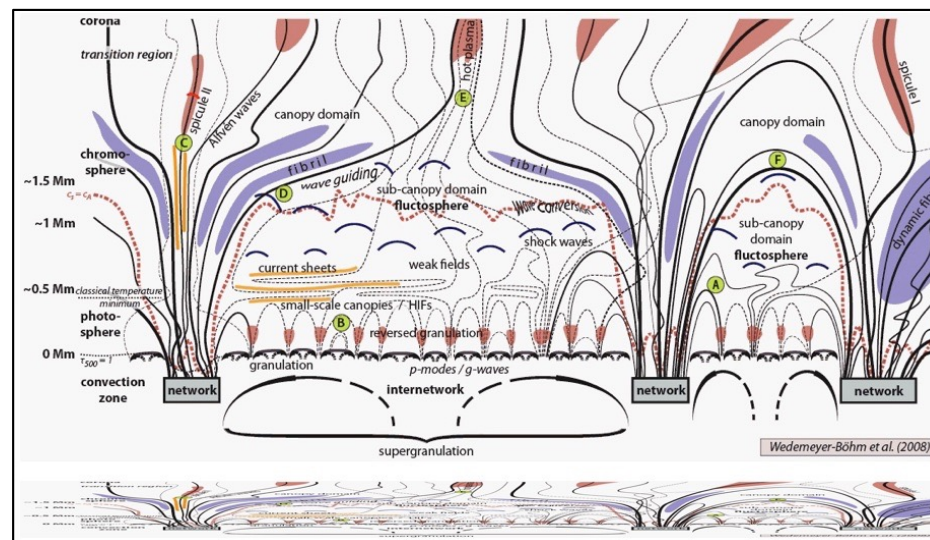
# The Zeeman effect



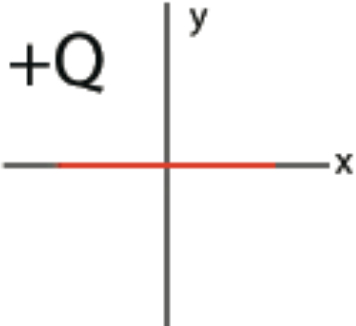
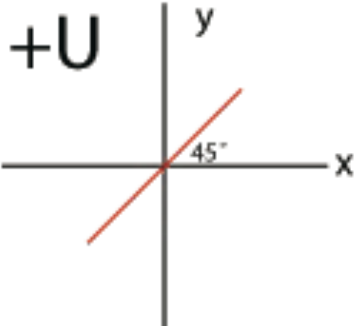
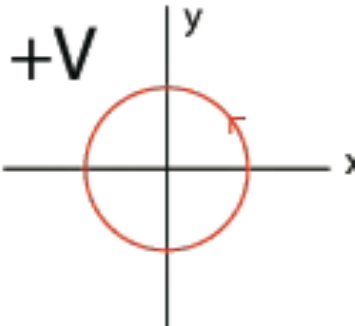
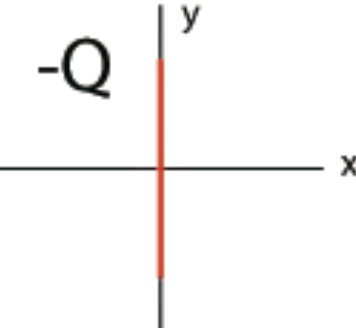
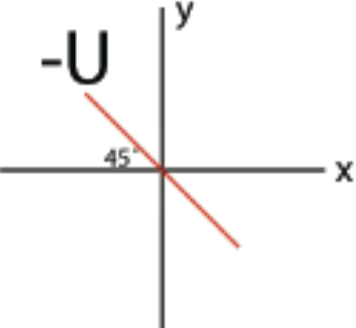
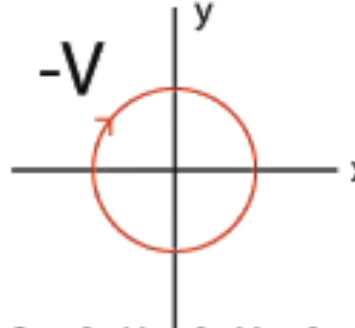
- Magnetic fields distort atomic energy levels in a polarized manner.
- For sunspot 0.1 T fields, the splitting is visible to the naked eye in a spectroscope. The “Bohr Magneton” is  $5.788 \times 10^{-5}$  eV/T.
- The splitting depends upon the angle  $\langle B, \text{eye} \rangle$ : the  $\sigma$  components (linear polarization) reflect the perpendicular field, whereas the  $\pi$  components (circular) reflect the line-of-sight field.

# How do we absorb this discovery?

- Hale showed that we can observe photospheric magnetism via Zeeman splitting
- Eclipse observations had already revealed “bar magnet” dipole light patterns
- How to link these phenomena?
  - Mathematical extrapolation (the dipole etc.)
  - New physical observations of the corona



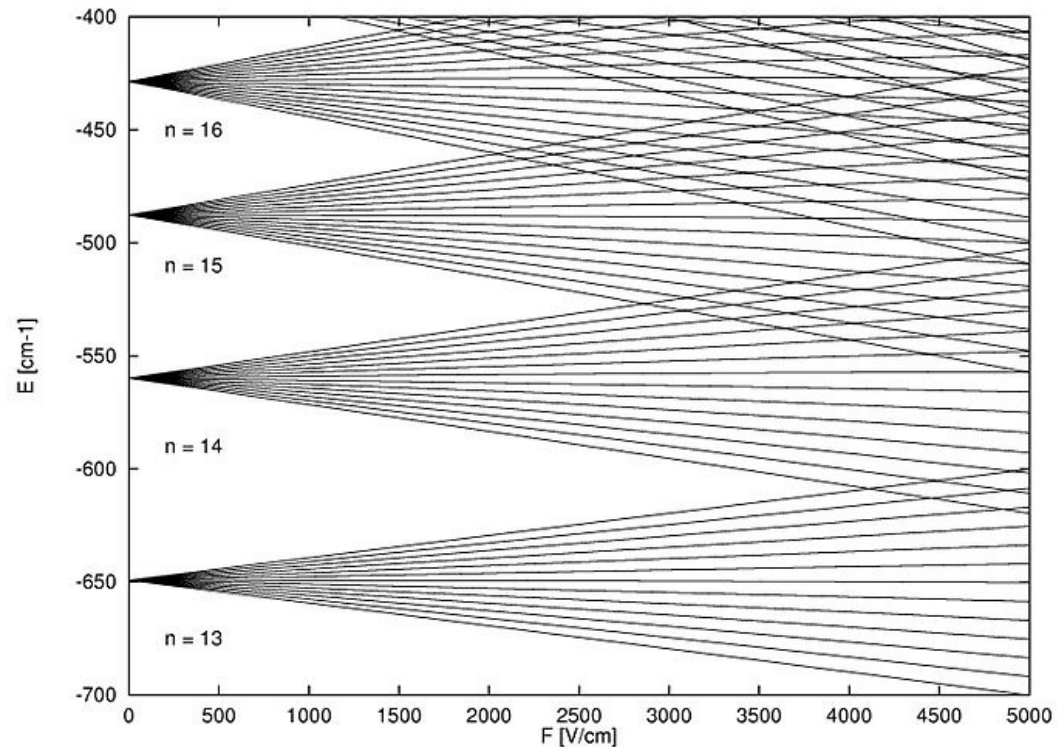
# Three of the Stokes parameters

| 100% Q  | 100% U   | 100% V  |
|---|--|---|
| <p><b>+Q</b></p>  <p><math>Q &gt; 0; U = 0; V = 0</math></p> <p>(a)</p>  | <p><b>+U</b></p>  <p><math>Q = 0; U &gt; 0; V = 0</math></p> <p>(c)</p>  | <p><b>+V</b></p>  <p><math>Q = 0; U = 0; V &gt; 0</math></p> <p>(e)</p>  |
| <p><b>-Q</b></p>  <p><math>Q &lt; 0; U = 0; V = 0</math></p> <p>(b)</p> | <p><b>-U</b></p>  <p><math>Q = 0; U &lt; 0; V = 0</math></p> <p>(d)</p> | <p><b>-V</b></p>  <p><math>Q = 0; U = 0; V &lt; 0</math></p> <p>(f)</p> |

(the fourth Stokes parameter is just the intensity I)

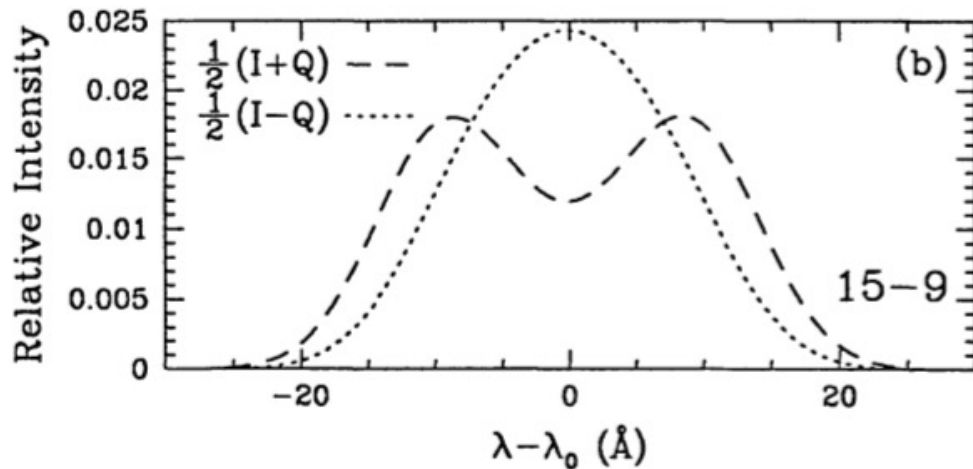
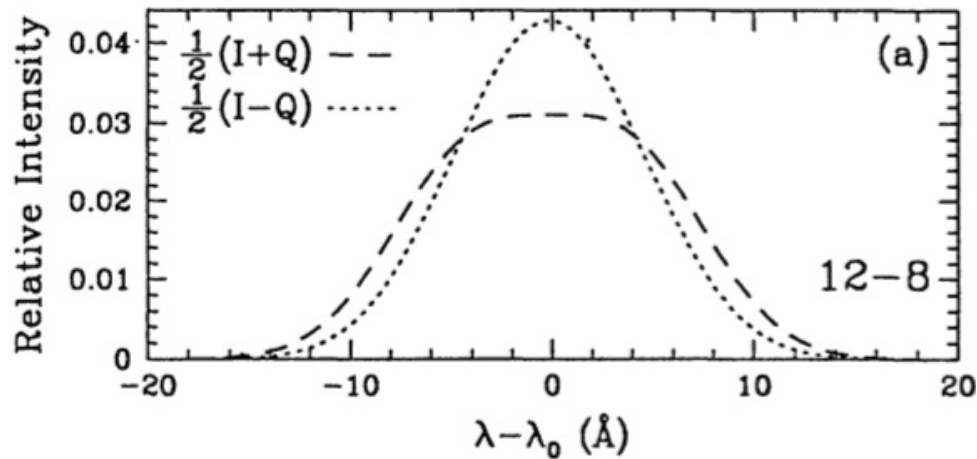
# E observations

- **The Stark effect** results from level shifts due to the presence of an electric field (the generic term).
- **Pressure broadening** is the Stark effect at the atomic level in a dense plasma.





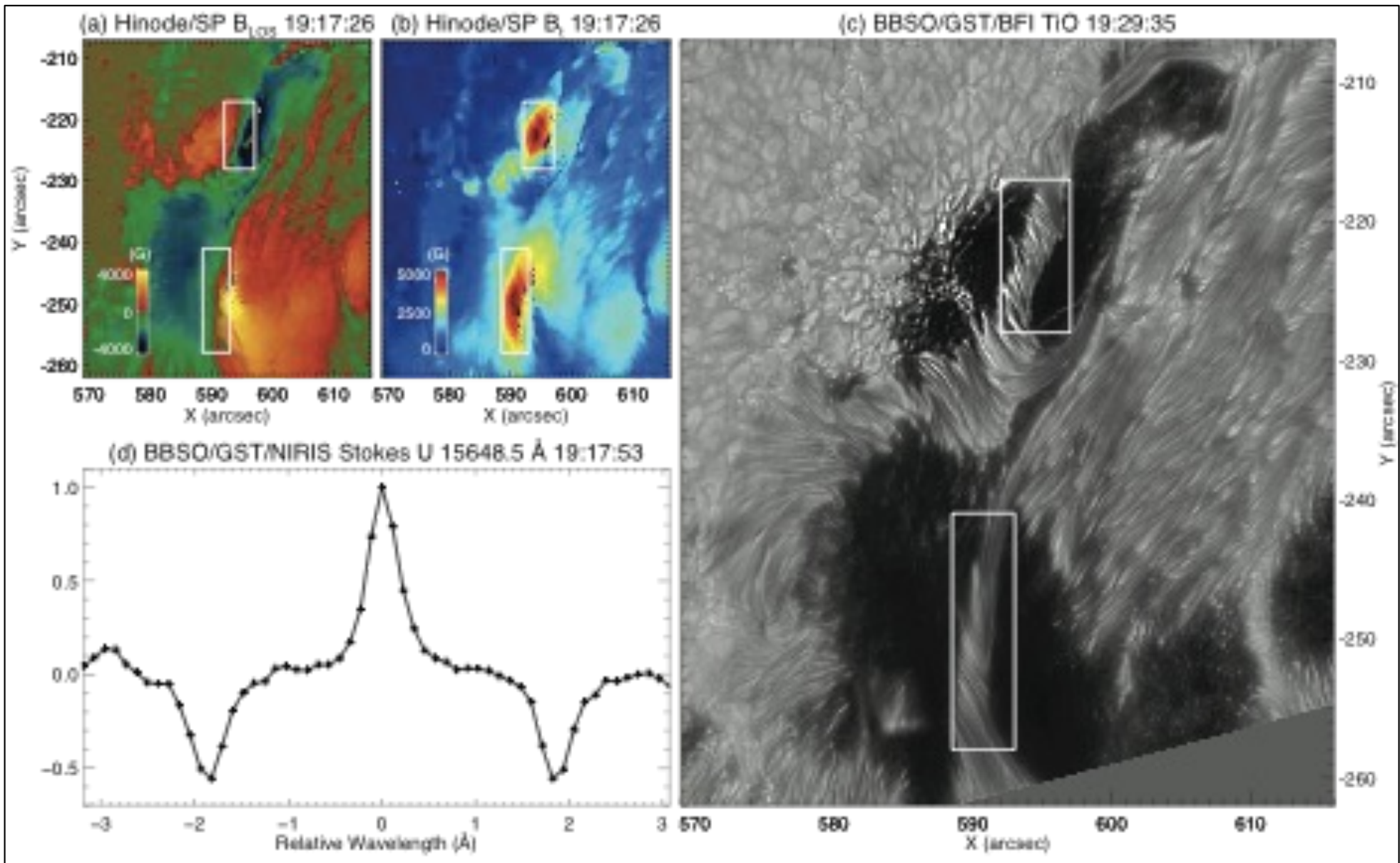
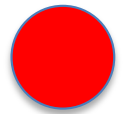
# Stark Effect



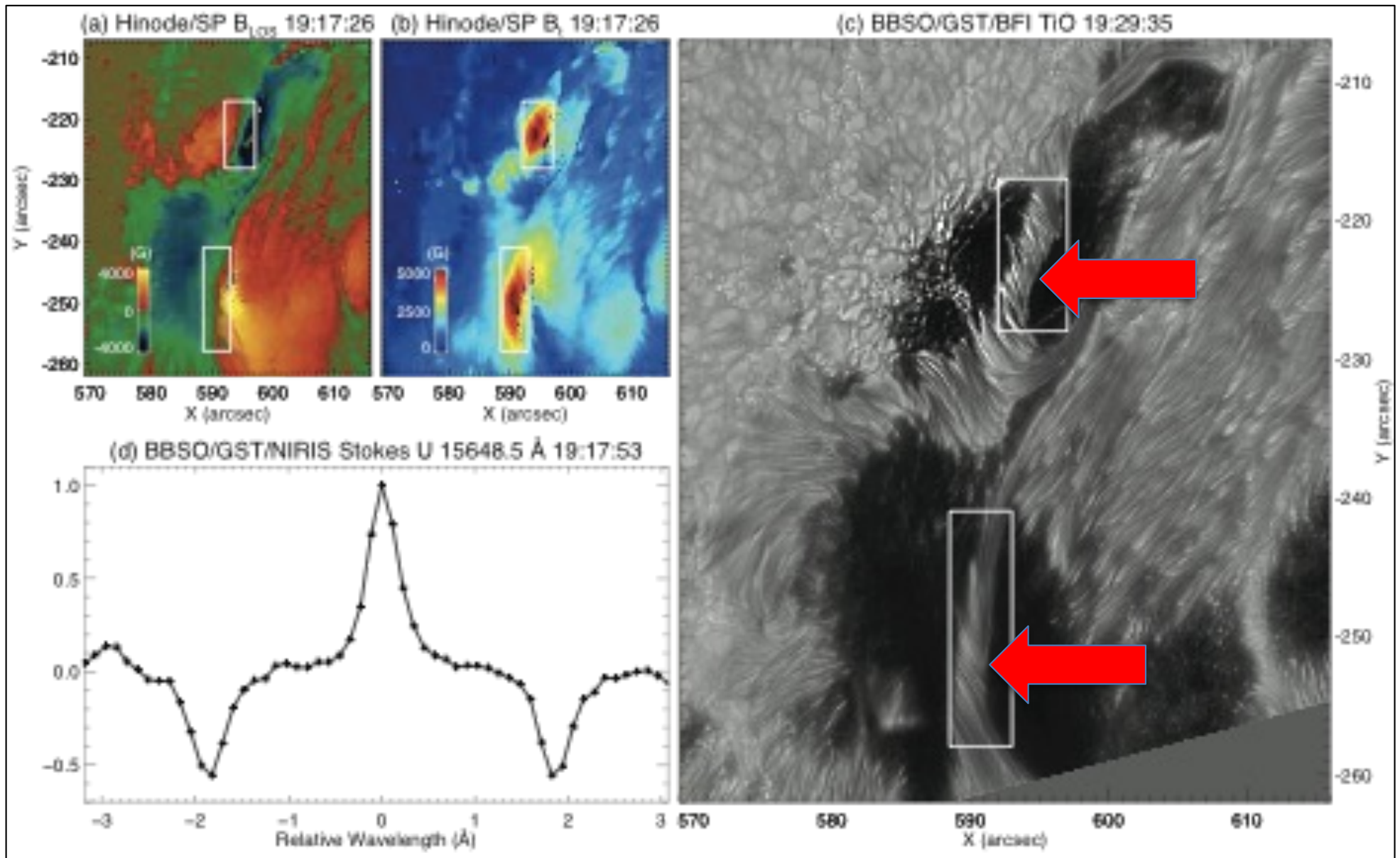
- The effect is strongest for “Rydberg” states (e.g.,  $n = 20$ ), hence ALMA?
- Minimum detectable macroscopic fields may be of order 100 V/m.\*

*\*This is not as crazy as it seems, since flare motions may cause huge electric fields*

# An extreme case (H. Wang)

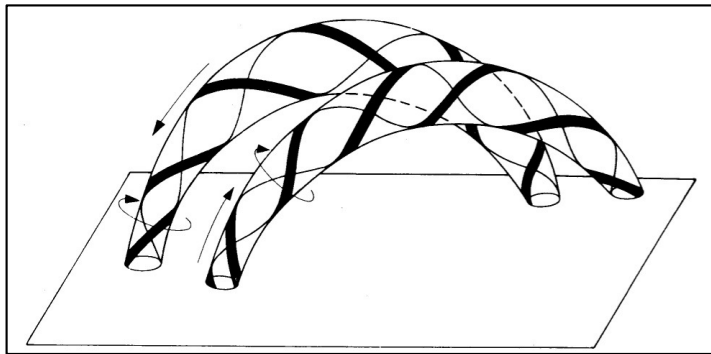


# An extreme case



# Flux ropes

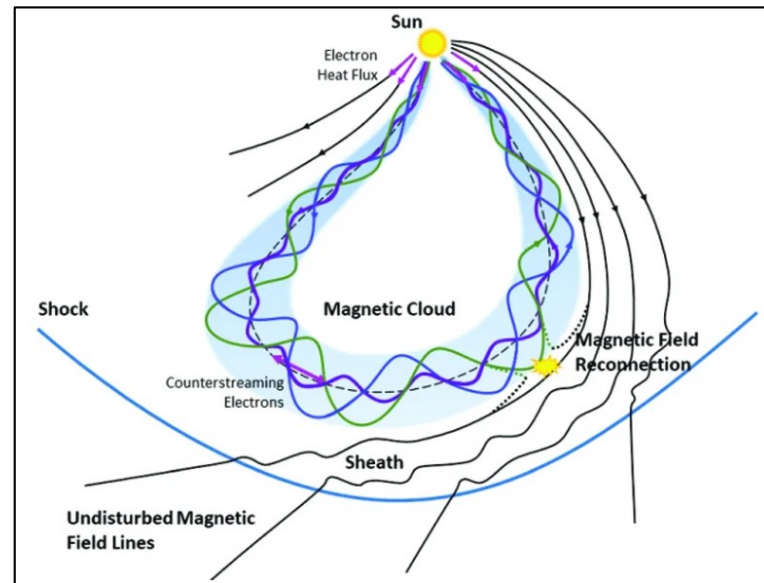
- In a magnetized plasma, the presence of velocity shearing naturally leads to helical structures (“flux ropes”) via the Lorentz force. There are various mathematical representations.
- Related solar topics: filaments, filament channels, magnetic clouds



Gold-Hoyle flare model

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} I \int_C \frac{d\mathbf{l} \times \mathbf{r}'}{|\mathbf{r}'|^3}$$

Biot-Savart law (simple)  
relates currents and B



Generic coronal mass ejection  
(CME) structure

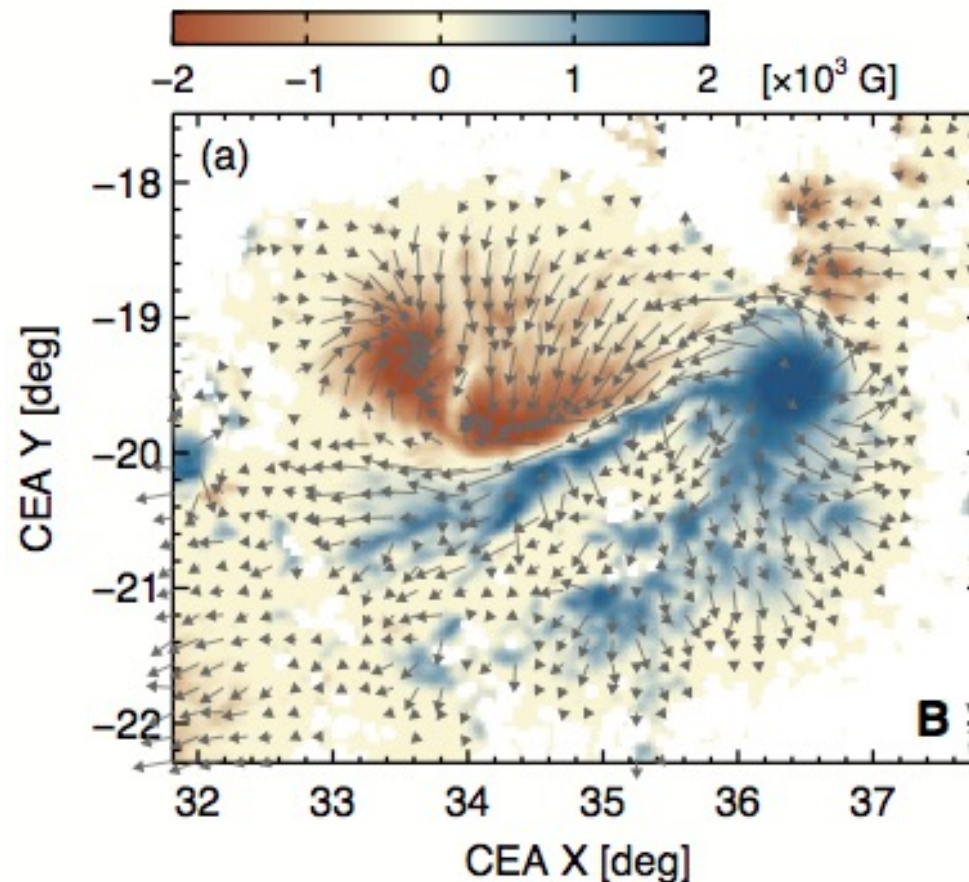
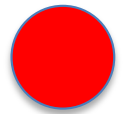
# Maxwell and the EM field

*What is it about the field that likes helices?*

- Maxwell's EM equations unified earlier work (Faraday, Ampere, *etc.*)
- The synthesis allowed Maxwell to make a laboratory-based estimate of  $c$
- The implication of this was the existence of a new object: the *field*
- The EM field requires the *curl* operator  $\nabla$
- Next historical step: Einstein and GR?

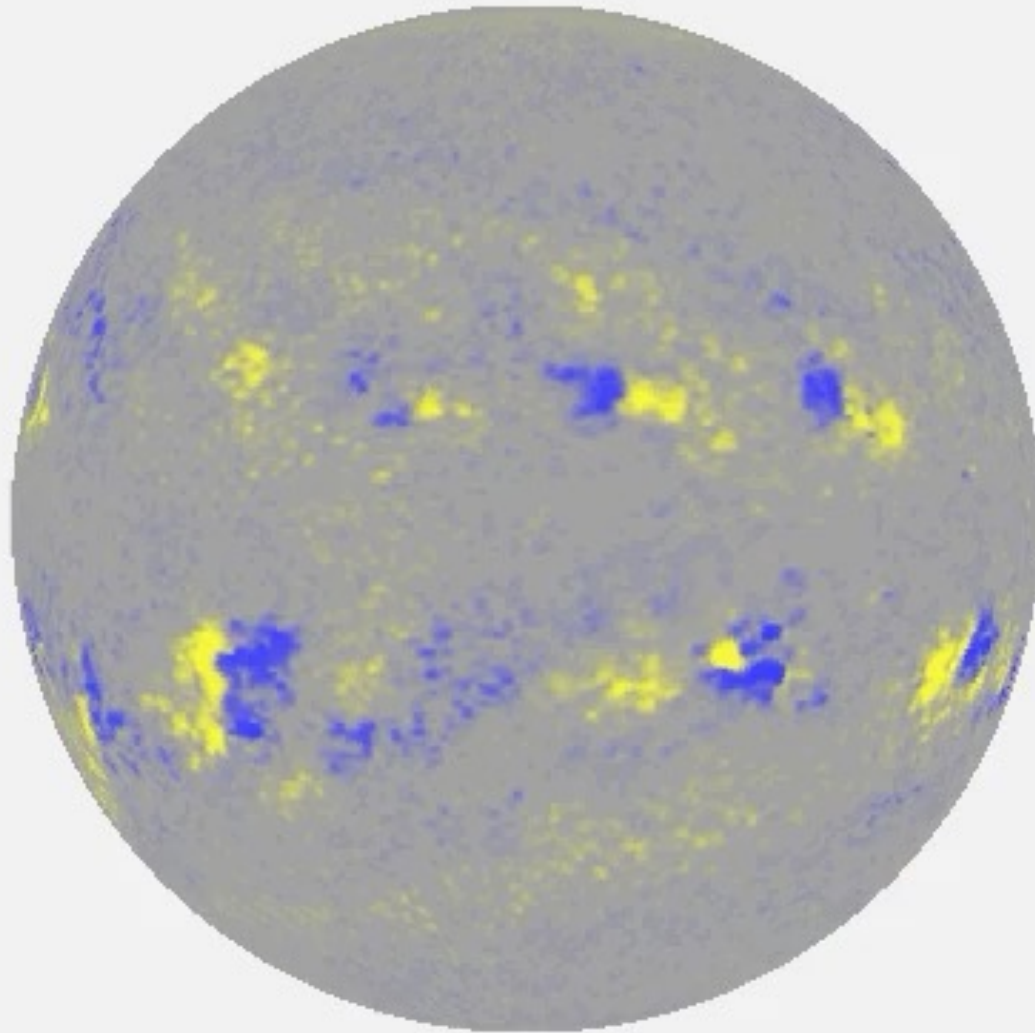


# Vector field at photosphere

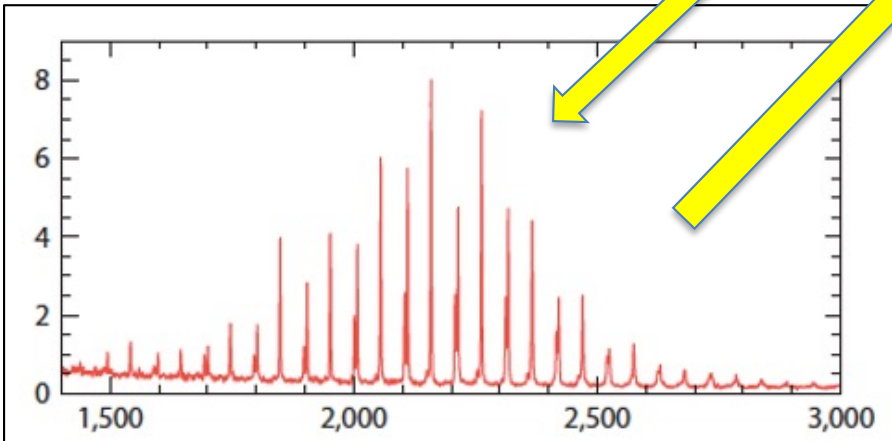
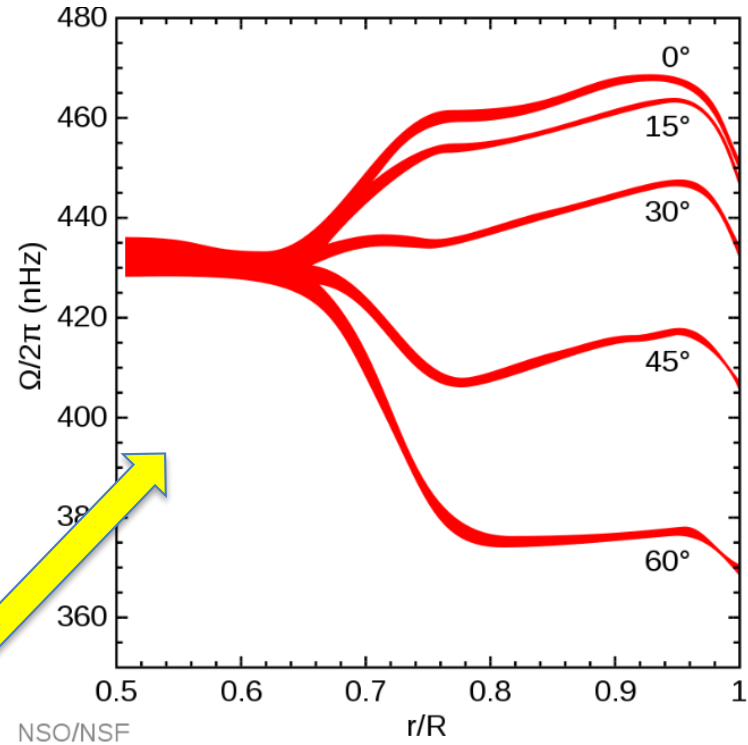
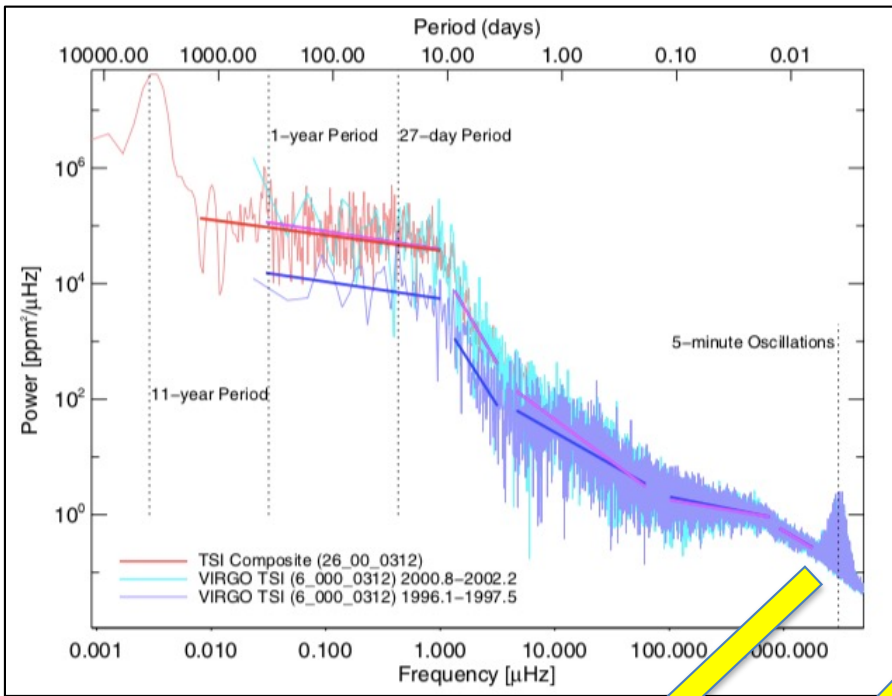
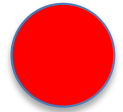


- With both parallel and perpendicular components, modulo a sign, one can infer the vector **B**.
- Note the "Polarity Inversion line" (PIL)
- This is also a "filament channel" with a strong horizontal current

# Routine solar magnetograms



# Helioseismology and the Tachocline

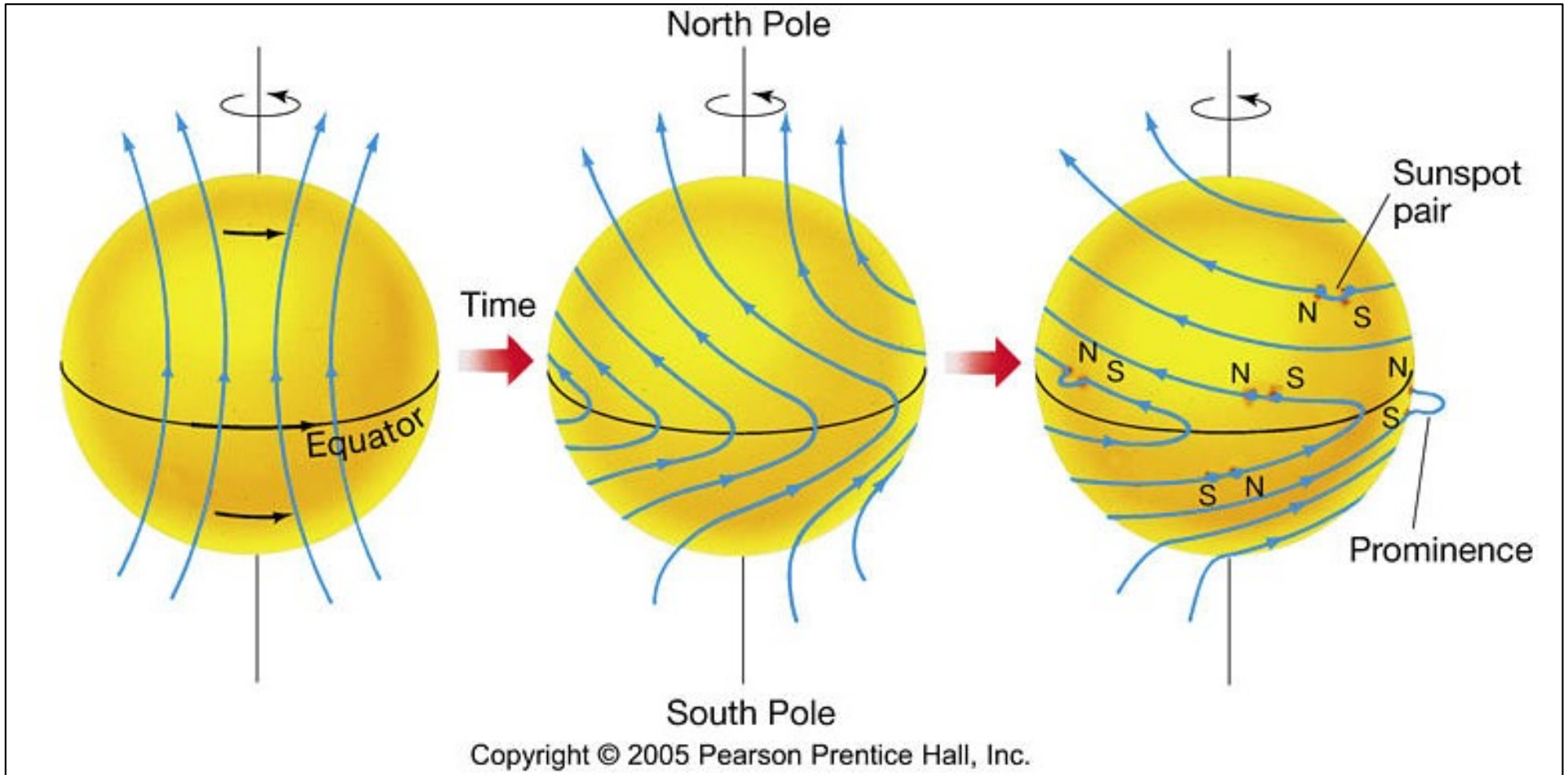


This example from a star, not the Sun!

- The solar resonances (“p-modes” or “5-min oscillations” allow us to locate the “tachocline”.
- \* This may be the site of the solar dynamo and magnetism.



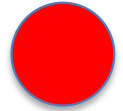
# Global: the Hale cycle



“Poloidal” fields amplify, due to differential rotation, and become “toroidal”

# The problem of coronal magnetism

- We can use the Zeeman effect in the photosphere to determine the  $\mathbf{B}$  (modulo symmetry-breaking) on a “surface,” but the 3D field is required to understand energetics.
- Unfortunately, the corona is optically thin – polarization signals are weak, and the structure is ill-defined because of the 3<sup>rd</sup> dimension.
- Can we just *extrapolate* from the photospheric boundary into the corona? With full Stokes info, we know the vector  $\mathbf{B}$  in the “plane” of the photosphere.



# Extrapolating the photospheric field

i) Field derived from a scalar potential ( $\nabla \times \mathbf{B} = 0$ )

ii) Linear force-free models, LFF ( $\nabla \times \mathbf{B} = \alpha \mathbf{B}$ )

iii) Non-linear force-free models, NLFF:

$$\nabla \times \mathbf{B} = \alpha(x,y) \mathbf{B}$$

.....

iv) MHD models: zero beta, “cold plasma”;  $P_B \gg P_g$

.....

v) MHD models, multi-fluid MHD

vi) Non-MHD models: any dynamics, actually

Note that the solar corona is a near-fully ionized plasma, and so its electrical resistivity is very small (of order  $10^{-7}$  Ohm-m).

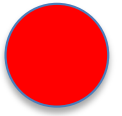
# Potential fields

- Extrapolation would be OK if there were no plasma in the corona, but if that were the case we couldn't see it!
- We could in principle ignore this basic fact and just plow ahead.
- In the absence of currents ( $\nabla \times \mathbf{B} = 0$ )\*, one can uniquely extrapolate the coronal magnetic field via spherical harmonics.
- An example would be to represent the photospheric magnetic field as a set of *magnetic charges* from which a multipole expansion could determine the coronal field.
- In the popular "PFSS" model\*\*, in which the Laplace's equation is solved inside a concentric spherical domain, outside of which the field is assumed to be radial.

\* Thus, solutions of Laplace's equation:  $\nabla^2 \mathbf{B} = 0$

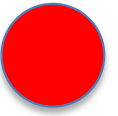
\*\* "Potential Field Source Surface"

# Limitations on mathematical extrapolations



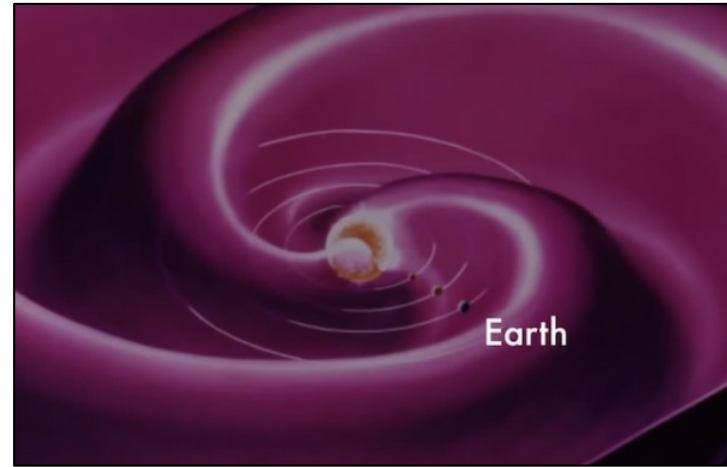
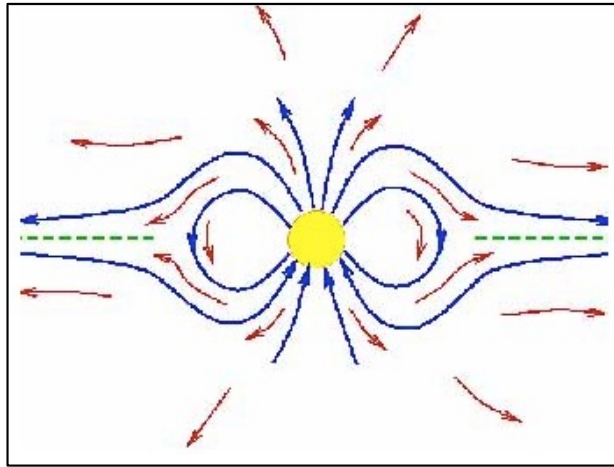
- The chromosphere is not force-free and not static
- Unbalanced currents penetrate the corona
- The base of the corona cannot be an equipotential surface
- The top of the corona (the base of the solar wind) is not low- $\beta$

# Other ways to measure the coronal magnetic field



- **The Hanle effect** results from line depolarization when the collision frequency exceeds the Larmor frequency.
- **Faraday rotation** alters the polarization angle of a linearly polarized background source (e.g., a quasar).
- **Image striations** can be interpreted in terms of field direction: note that would be  $\sim 2/3$  of the problem of measuring  $\mathbf{B}(\mathbf{r})$ .
- **Gyrosynchrotron** radiation has natural circular polarization.
- **Coronal coherent emissions** contain many clues about coronal B.
- **Birefringence effects** polarize even thermal emission from magnetized plasmas.
- **Magnetically induced transitions** a new atomic spectroscopy tool.
- **TeV cosmic-ray shadows** a new technique, largely unused yet.
- ***In situ* measurement** (Parker Solar Probe).

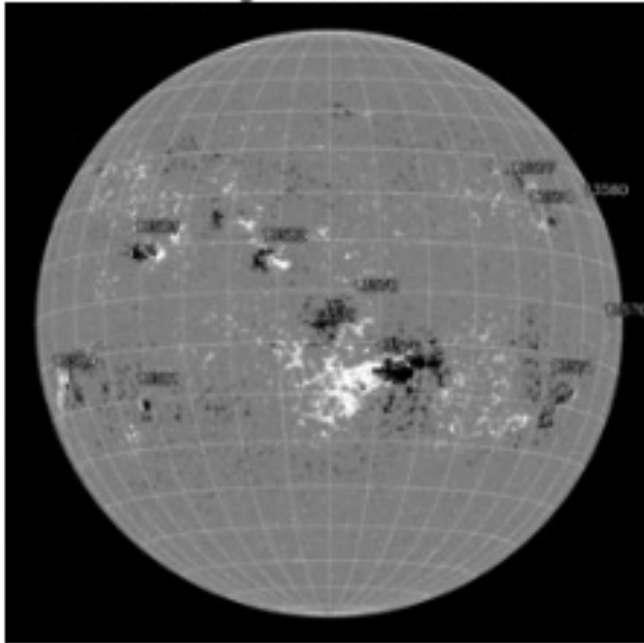
# Heliospheric current sheet



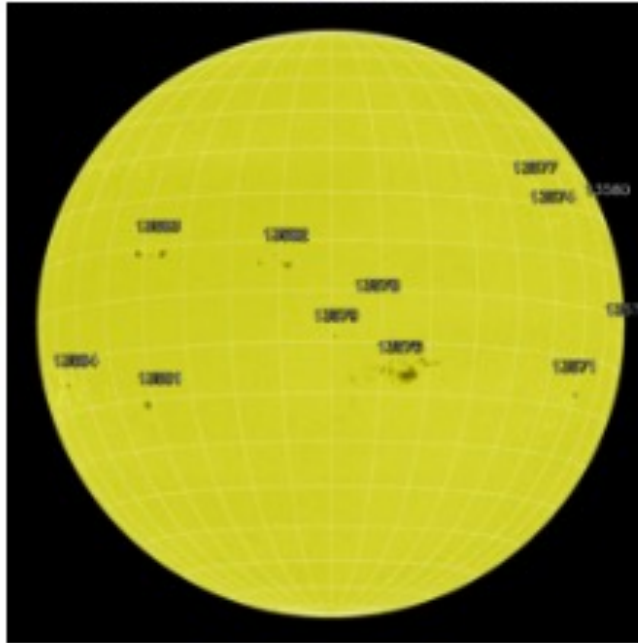
- A solar wind, with radial field, violates the potential-field concept of zero current.
- The rotation of the Sun induces a spiral pattern in the heliosphere due to the solar wind.
- Field asymmetry introduces warps (“sectors” at 1 AU) in the structure.

# The Sun yesterday

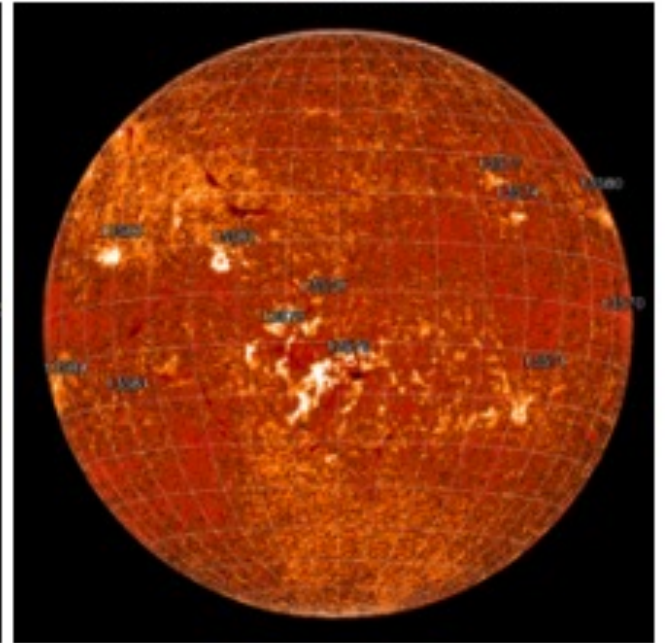
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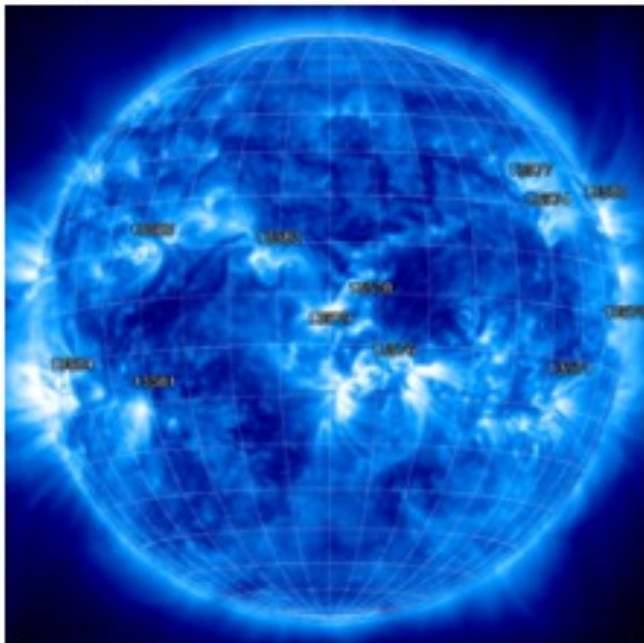
HMI 6173Å 20240211 07:34



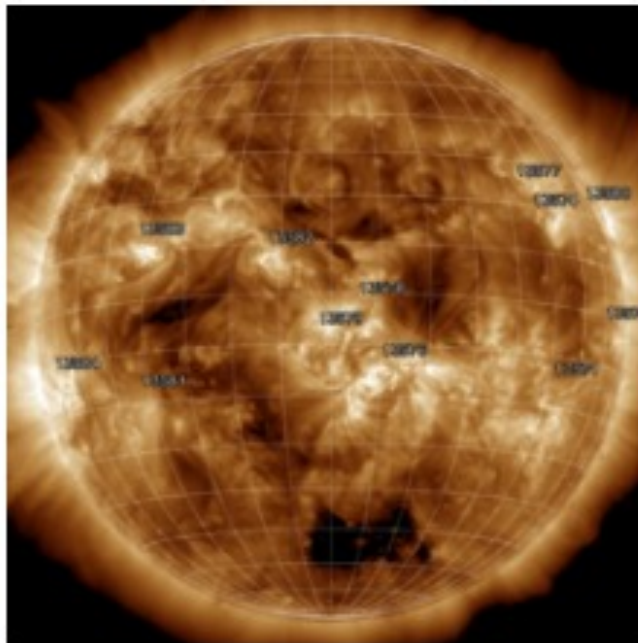
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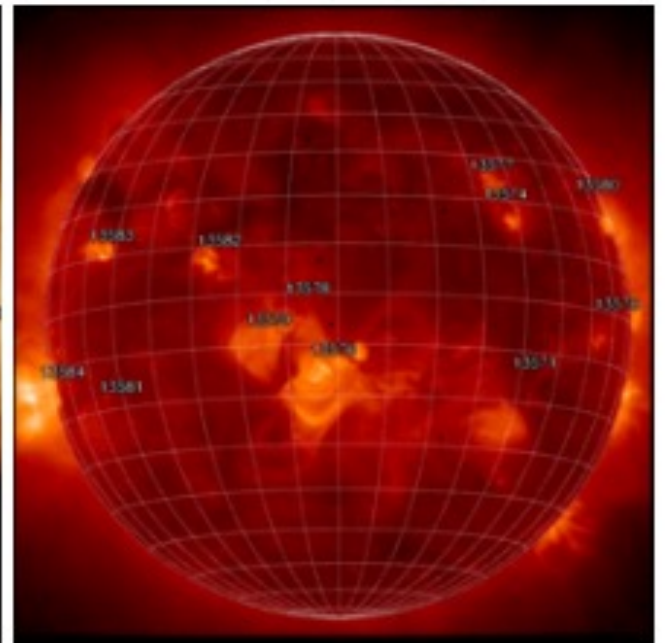
SWAP 174Å 20240211 05:41



AIA 193Å 20240211 08:24



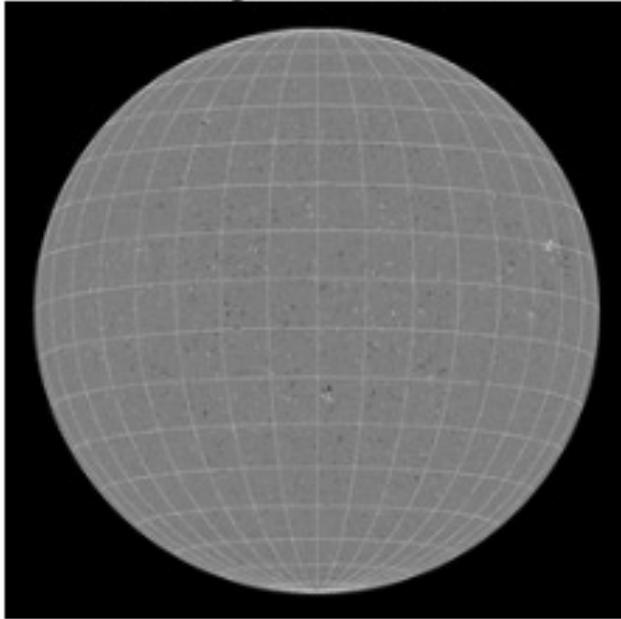
XRT 20240210 06:04



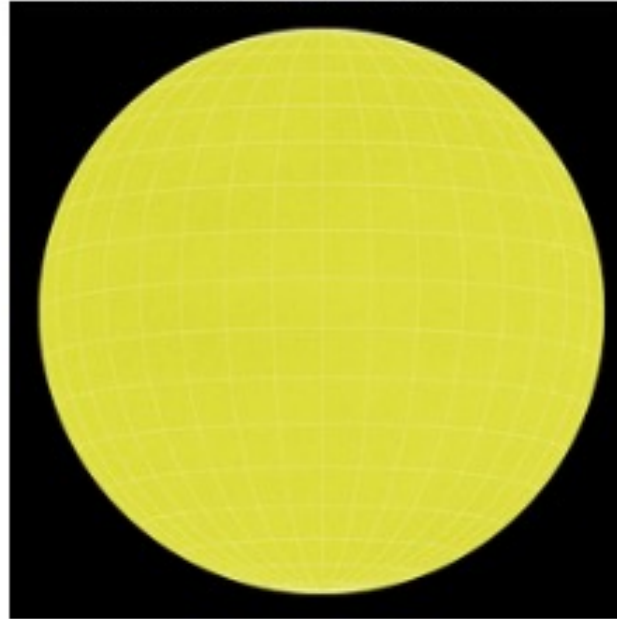


# The Sun five years ago

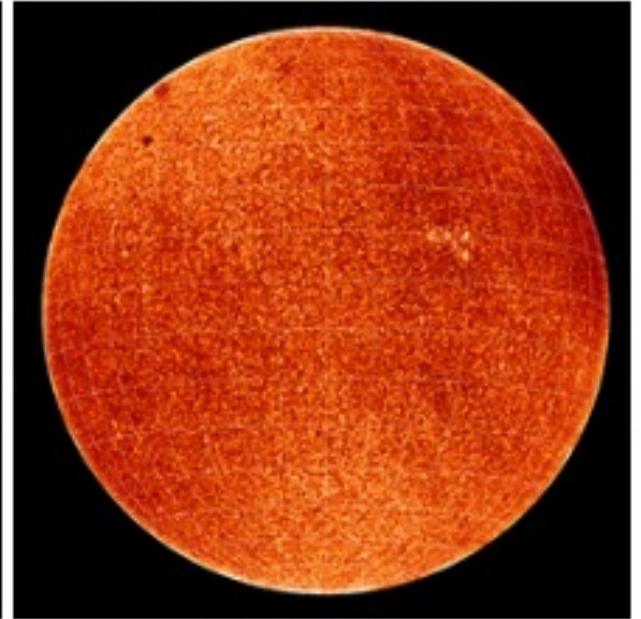
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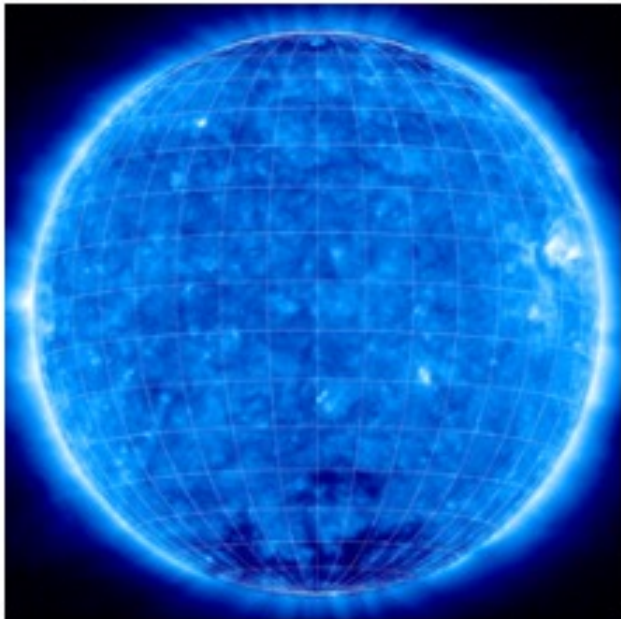
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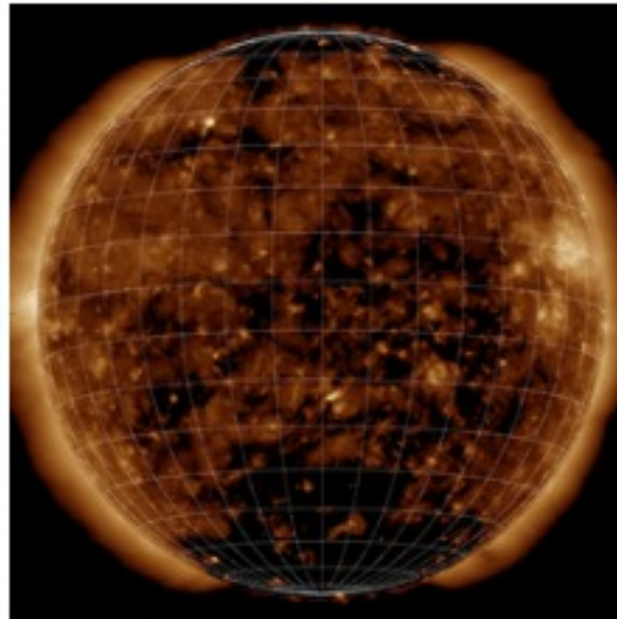
GHN H $\alpha$  20190209 08:47



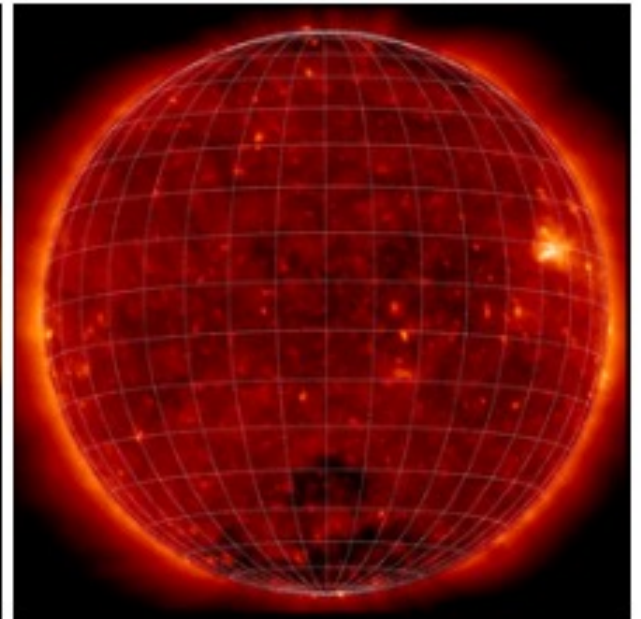
SWAP 174Å 20190211 18:20



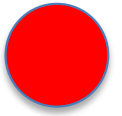
AIA 193Å 20190211 20:24



XRT 20190211 06:01

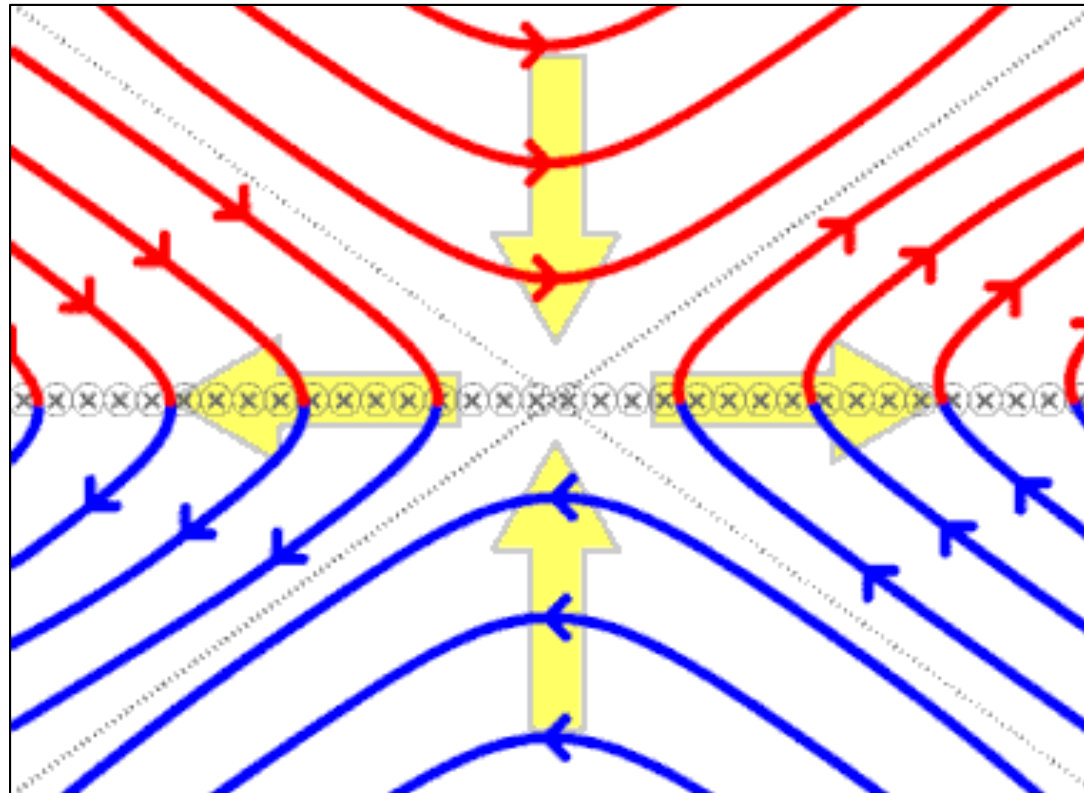


# The stability of the coronal field



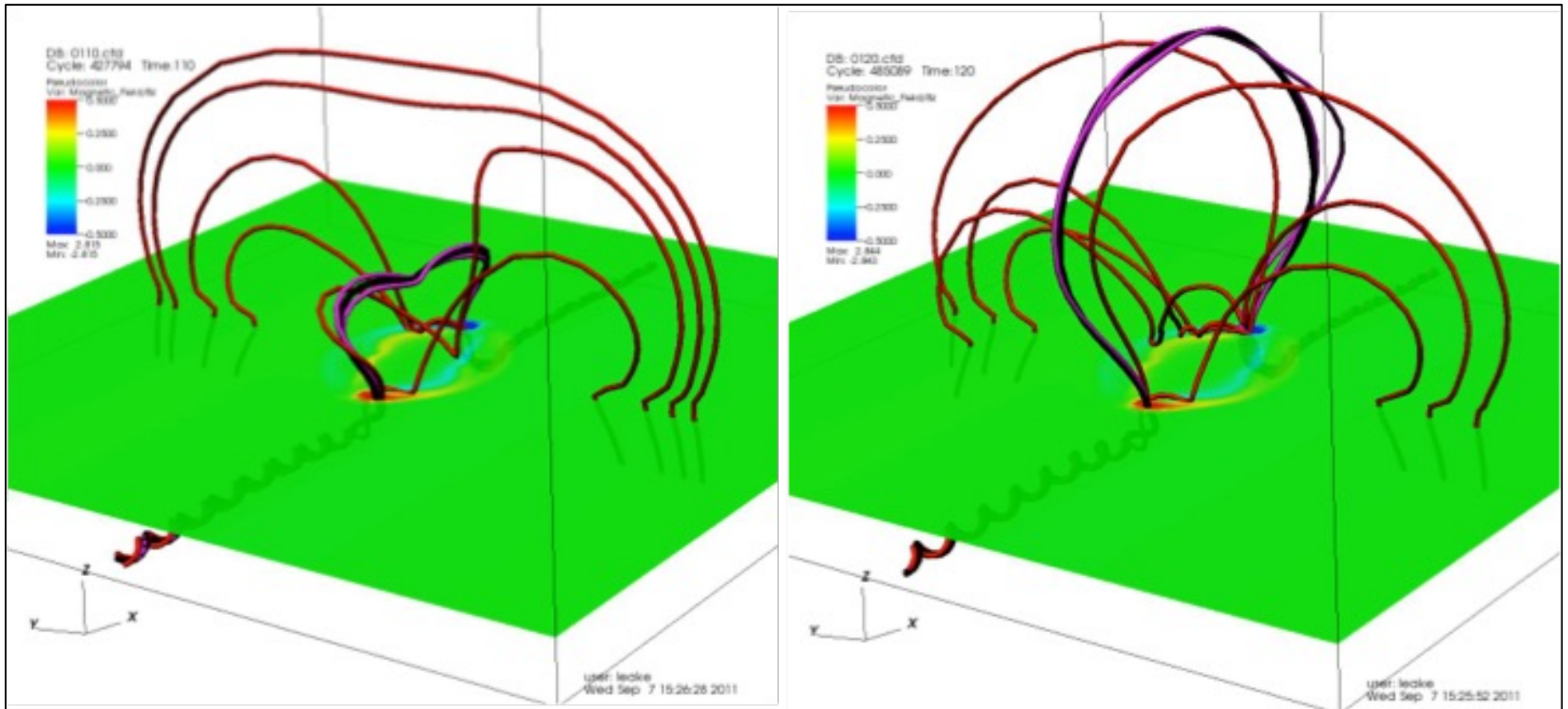
- A solar magnetic transient (flare or CME; coronal mass ejection) derives energy from the magnetic field:  $\mathbf{B}^2/8\pi$  (in CGS units). It must become more potential-like as this happens.
- Large currents ( $10^{12}$  A) may flow in the corona, storing this energy inductively.
- Identifying the plasma instabilities that do this is an open problem now, and of course very important for flare prediction (“space weather”).
  - Ideal MHD instability (no magnetic reconnection), such as the kink instability
  - Resistive instabilities, such as the torus instability or the “tether-cutting” cartoon

# “Magnetic reconnection” (flux transfer)



- This great Wikipedia graphic shows how magnetic field lines are thought to “reconnect,” and this is the key physics of many models of flares.
- Reconnection is often described as a mechanism for energy release, but here the magnetic energy is constant - it is a steady state.









# The “torus instability”

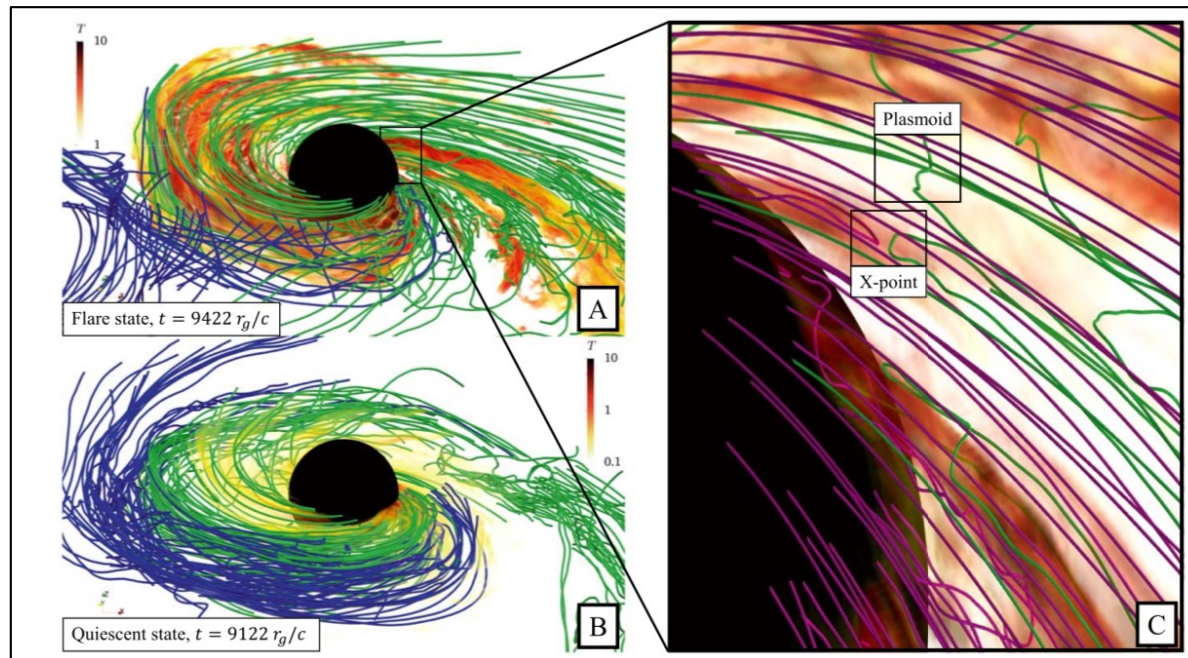


- A *flux rope* forms above a *photospheric inversion line*.
- This structure is held in place by overlying field.
- Torok & Kliem (2005) suggest that if this overlying field decreases in strength with height, the structure may blow.

# Magnetic reconnection elsewhere

## Black Hole Flares: Ejection of Accreted Magnetic Flux through 3D Plasmoid-mediated Reconnection

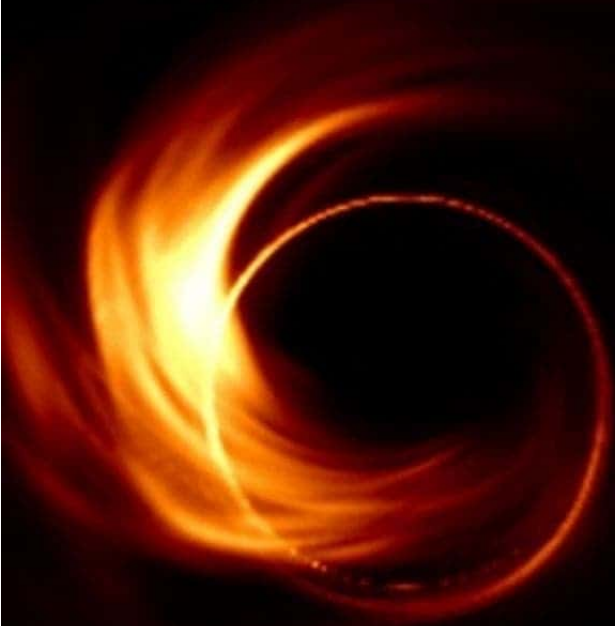
B. Ripperda<sup>1,2,3,11</sup> , M. Liska<sup>4,5,11</sup> , K. Chatterjee<sup>6,7</sup> , G. Musoke<sup>7</sup> , A. A. Philippov<sup>1</sup> , S. B. Markoff<sup>7,8</sup> ,  
A. Tchekhovskoy<sup>9</sup> , and Z. Younsi<sup>10</sup> 



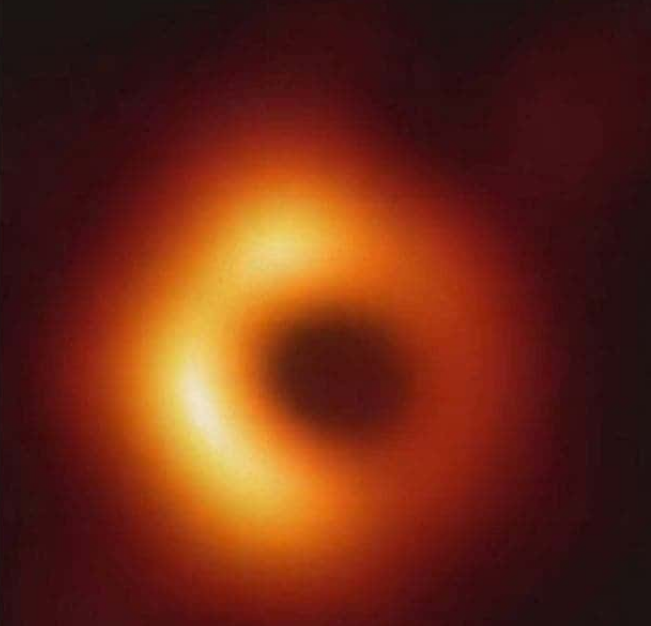
- Flaring activity occurs near black-hole horizons.
- One can explain it in the same way “standard models” explain solar flares, ie “plasmoid-driven” reconnection.

# Imaging a black hole

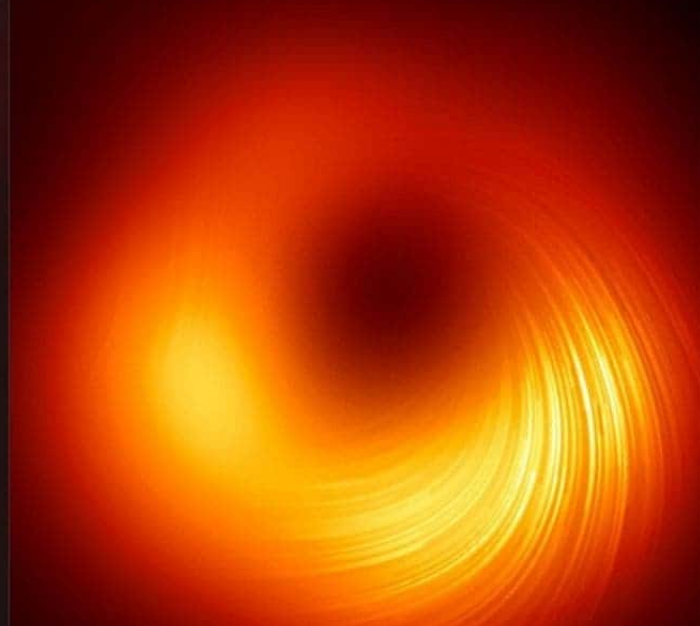
**BLACK HOLES: HOW IT STARTED. HOW IT'S GOING**



**THEORETICAL  
IMAGE OF M87'S  
BLACK HOLE**



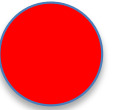
**FIRST DIRECT  
IMAGE OF M87'S  
BLACK HOLE**



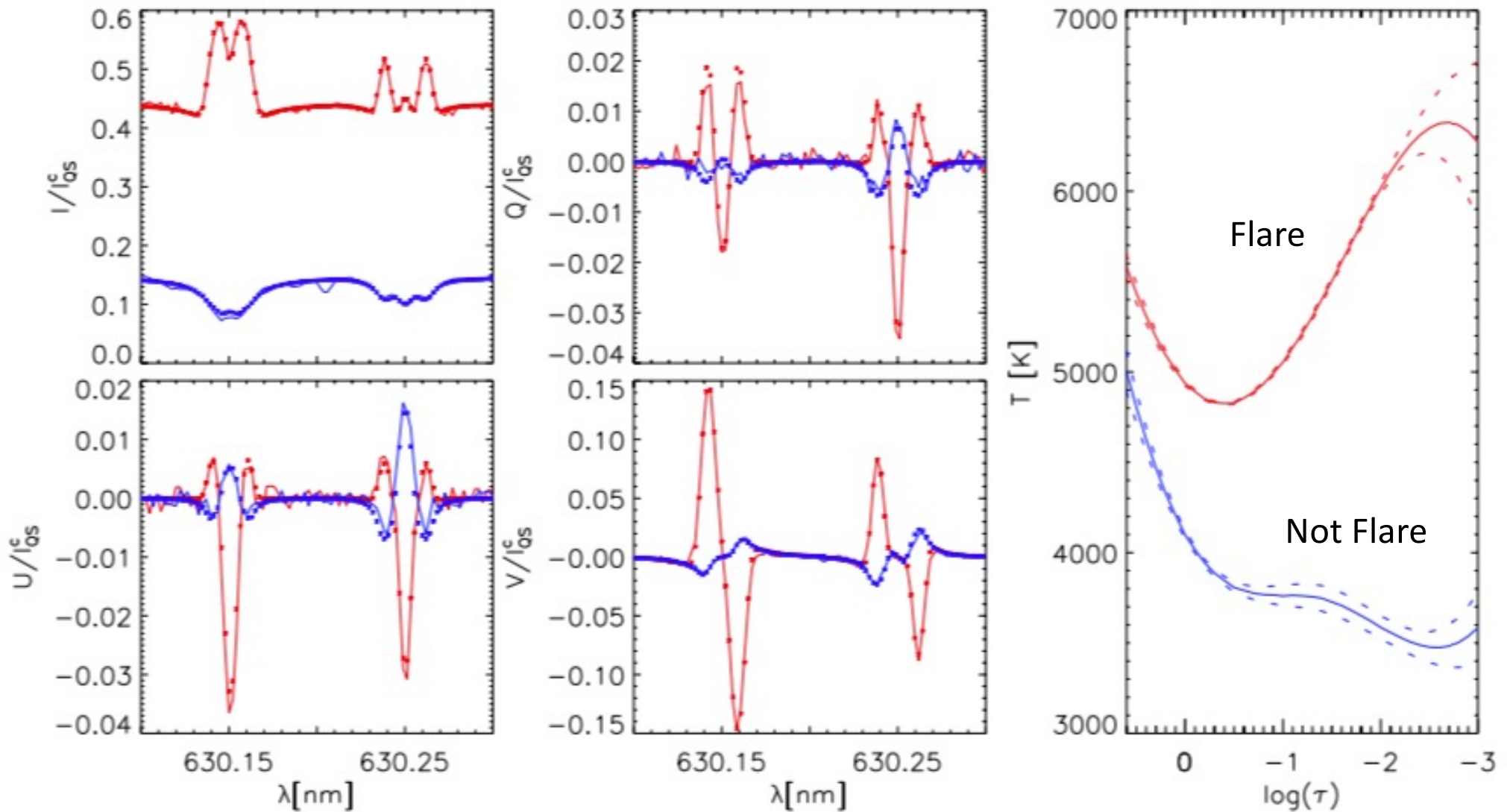
**SECOND DIRECT  
IMAGE OF M87'S  
BLACK HOLE**

# An example research paper

- Jurčák *et al.* 2018, “Heating of the solar photosphere during a white-light flare”
  - The behavior of an actual photospheric spectral line
  - “Line inversions” to obtain quasi-3D structure
  - What the Stokes parameters can tell you
  - How flares might work
- The “SIR” Stoke Inversion code translates line profiles, “forward fitting”  $T$  to a coarse 1D map at 5 points over a four-decade range of optical depths: many, many assumptions!

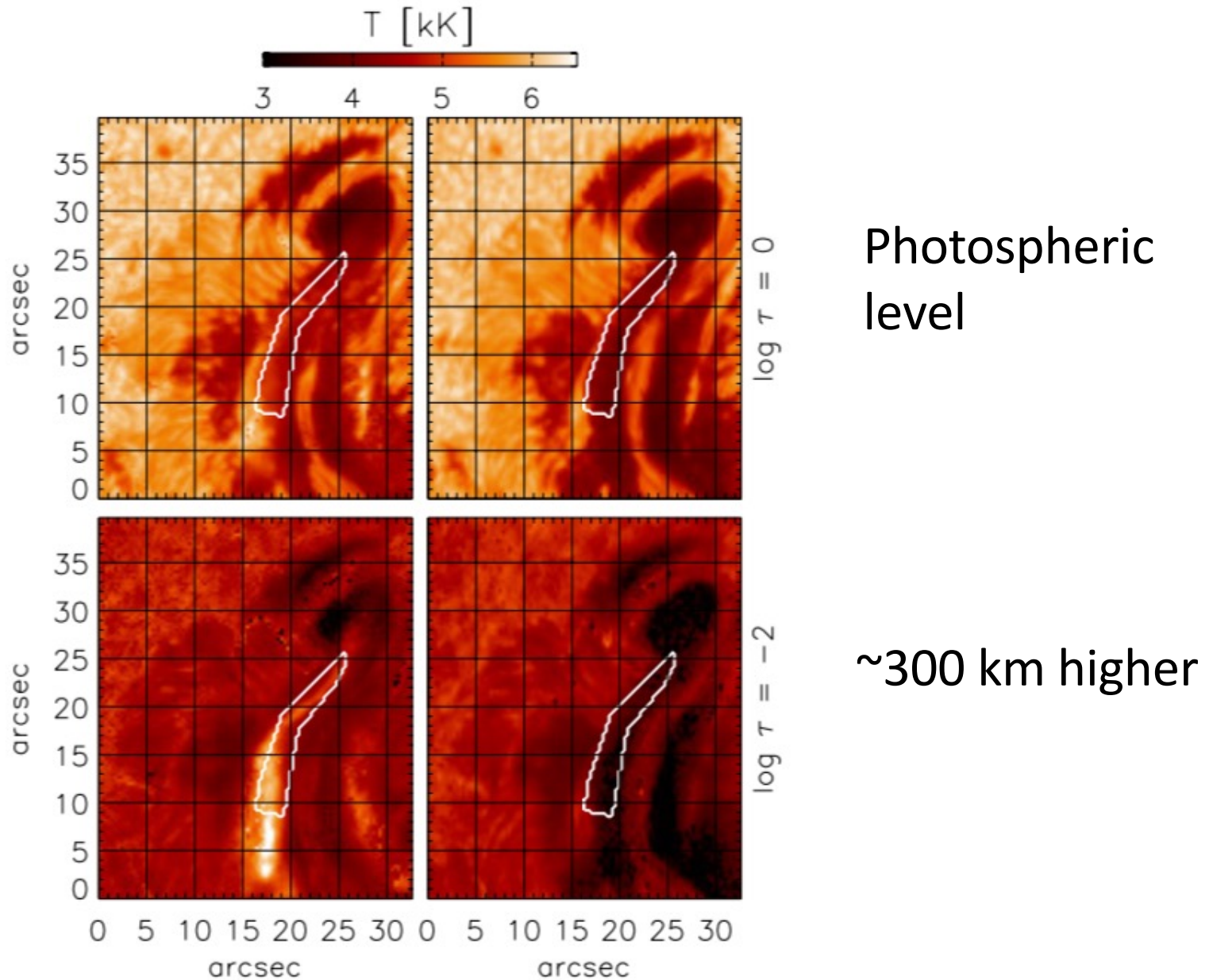


# *Hinode* line profiles and inversion for T



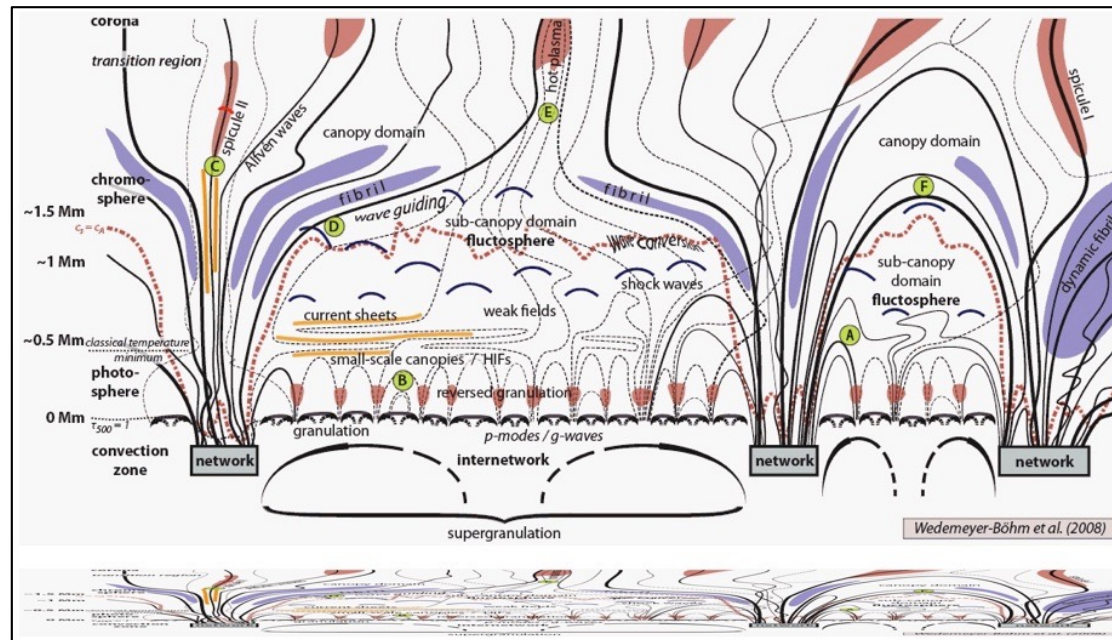


# *Hinode* line profiles and inversion



# Conclusions

- Magnetism makes the solar atmosphere interesting from the point of view of plasma physics.
- ALMA (now), DKIST (“first light” already), PSP and Orbiter (in orbit) offer wonderful new observational opportunities.





# Useful homework questions based on this material

- What are the forces acting on plasma in a coronal magnetic loop, rough quantitative estimate?
- For what value of  $B$  can a flux tube escape from the tachocline and emerge in one cycle?
- What is the electrical potential at the photosphere?
- At what stereoscopic angle can Stokes (I,V) do better than Stokes (I,Q,U,V)?

*These are questions to be answered by rough approximation, and the necessary facts should be here on the slides. I will post discussion of these items at*

*<http://www.ssl.berkeley.edu/~hudson/presentations/supa.220207/>*